



Statistical QCD

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IOP: 1994-1999.

The Institute of Mathematical Sciences

Hadronic world

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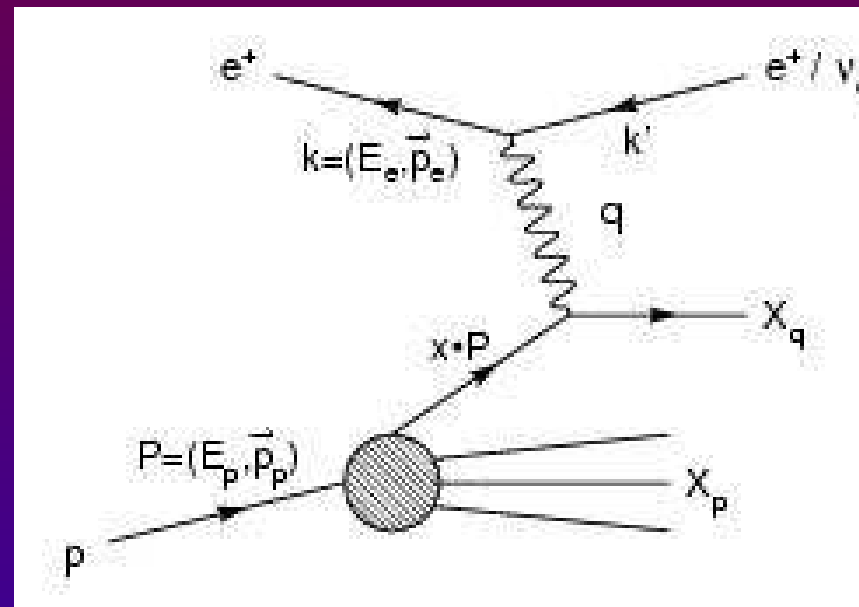
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If hadrons are fundamental/elementary? What is the theory of hadron world?

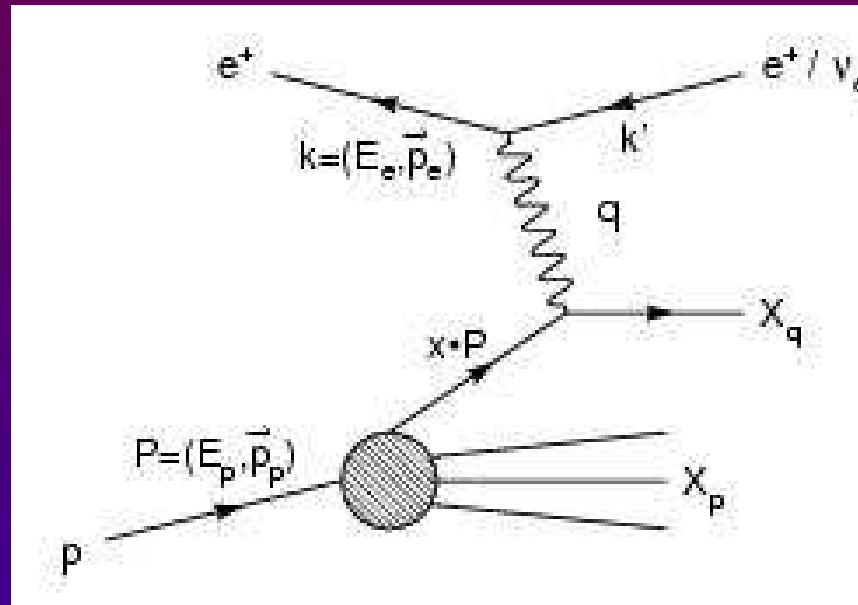
Structure of Hadrons

High energy electron proton scattering showed hadrons are composites of quarks and gluons



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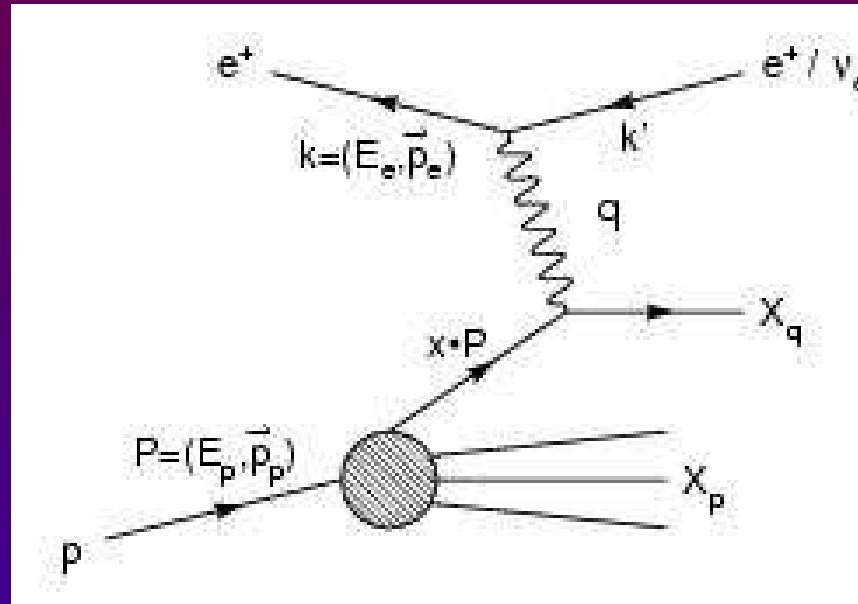
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By Friedman, Kendall, Taylor (1967-1973) Noble 1990.

Structure of Hadrons

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Theory of quarks and gluons must answer questions of hadronic world

Theory of Quarks and Gluons (QCD)

Quantum electrodynamics (QED)

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Particles with electric charge
Electron, positron, etc..

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Quarks, anti-quarks, gluons

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Photon – charge neutral

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Gluons – colored (8)

QCD

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}^a F^{a\mu\nu} + \sum_{f=u,d,s} \bar{\psi}_\alpha^f (i\gamma^\mu D_\mu - m_f)^{\alpha\beta} \psi_\beta^f$$

$$F_{\mu\nu}^a = \partial_\mu G_\nu^c - \partial_\nu G_\mu^c + gf^{cba} G_\mu^b G_\nu^a$$

$$D_\mu = \frac{\partial}{\partial x_\mu} - igT^a G_\mu^a$$

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This observation is known as **color confinement**

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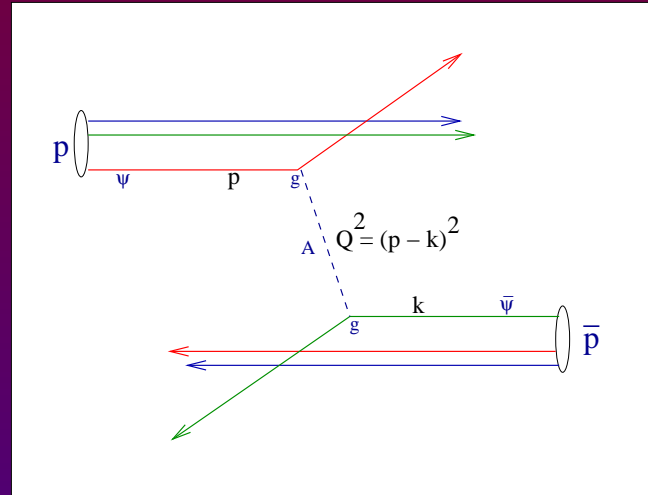
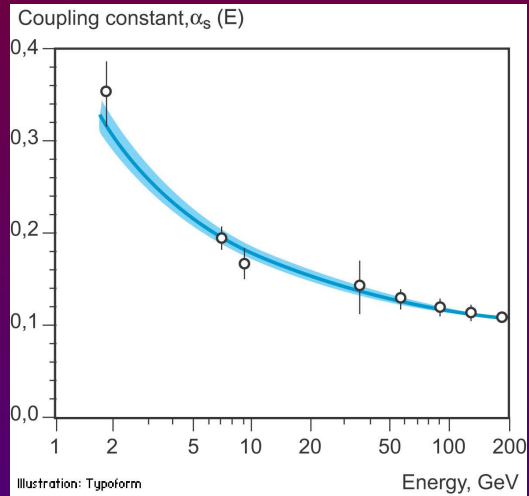
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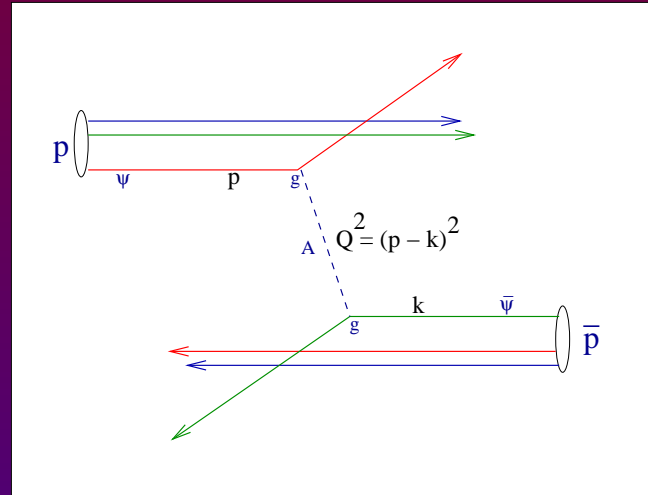
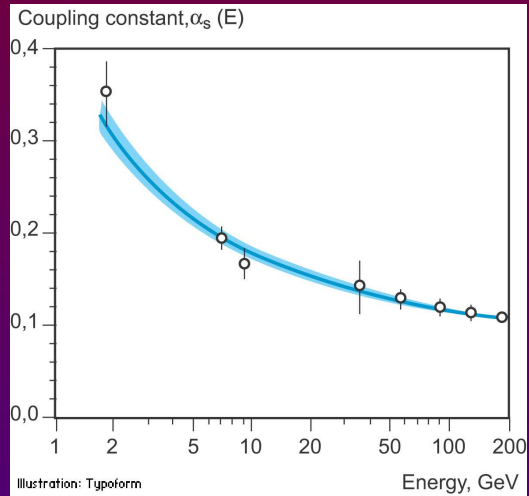
This observation is known as **color confinement**

Analytical proof of color confinement is still an open problem.

Asymptotic Freedom

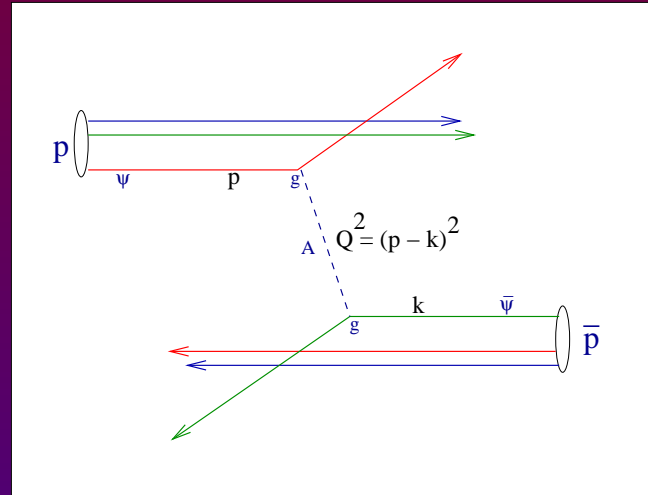
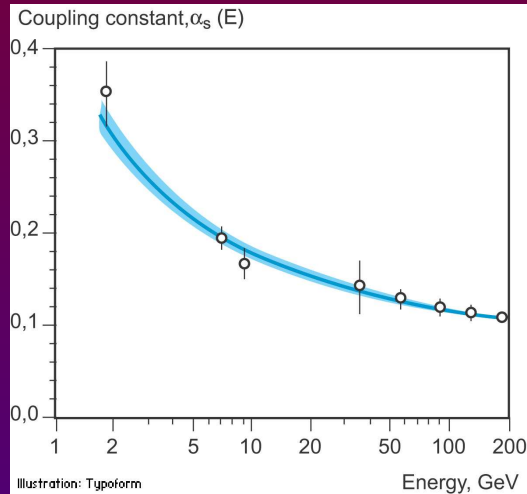


Asymptotic Freedom



$$g \sim \frac{1}{\log(Q^2/\lambda^2)}$$

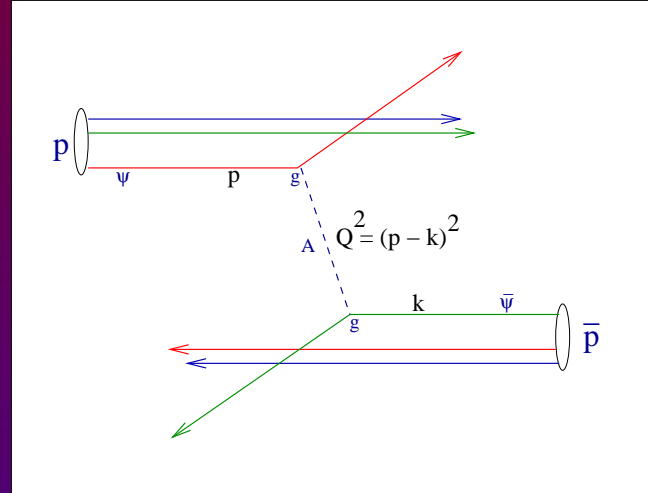
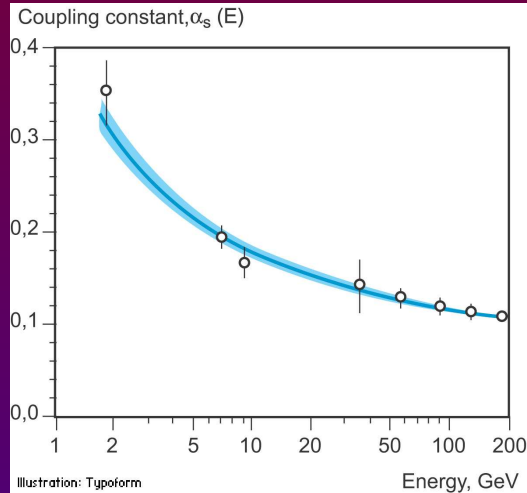
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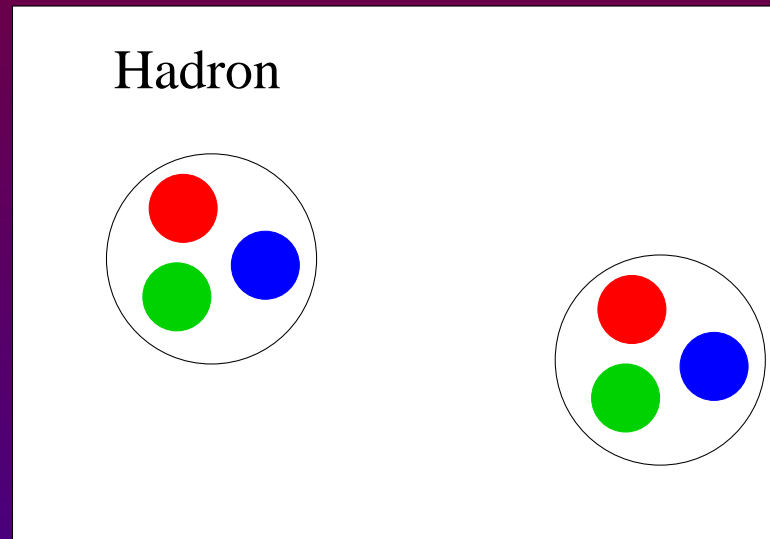


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Gross, Pultzer, Wilczek....2004 Nobel Prize
 g is small for $Q^2 \gg \lambda^2$, analytic (perturbative)

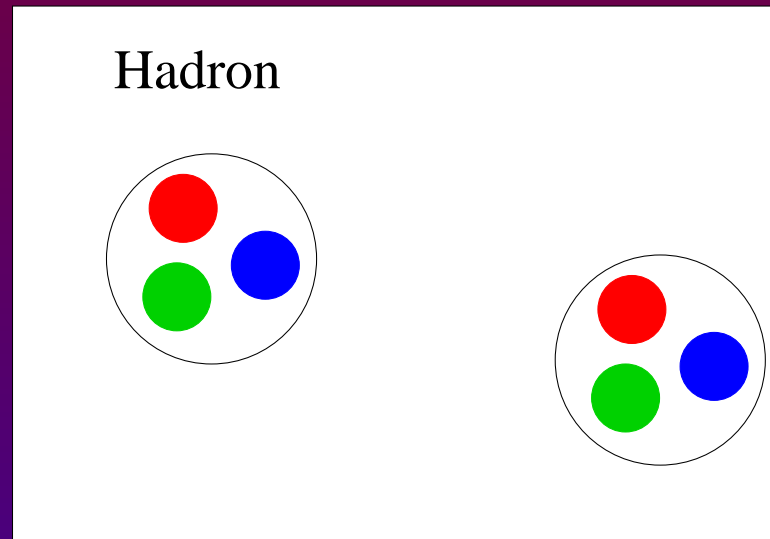
Few higher order corrections

Confinement



Hadrons are of size ~ 1.0 fm, so need ~ 1.0 fm³ space to exist

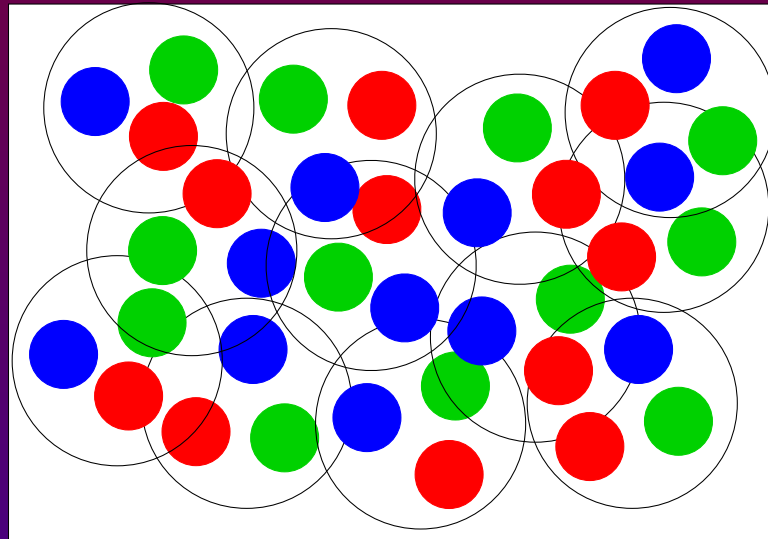
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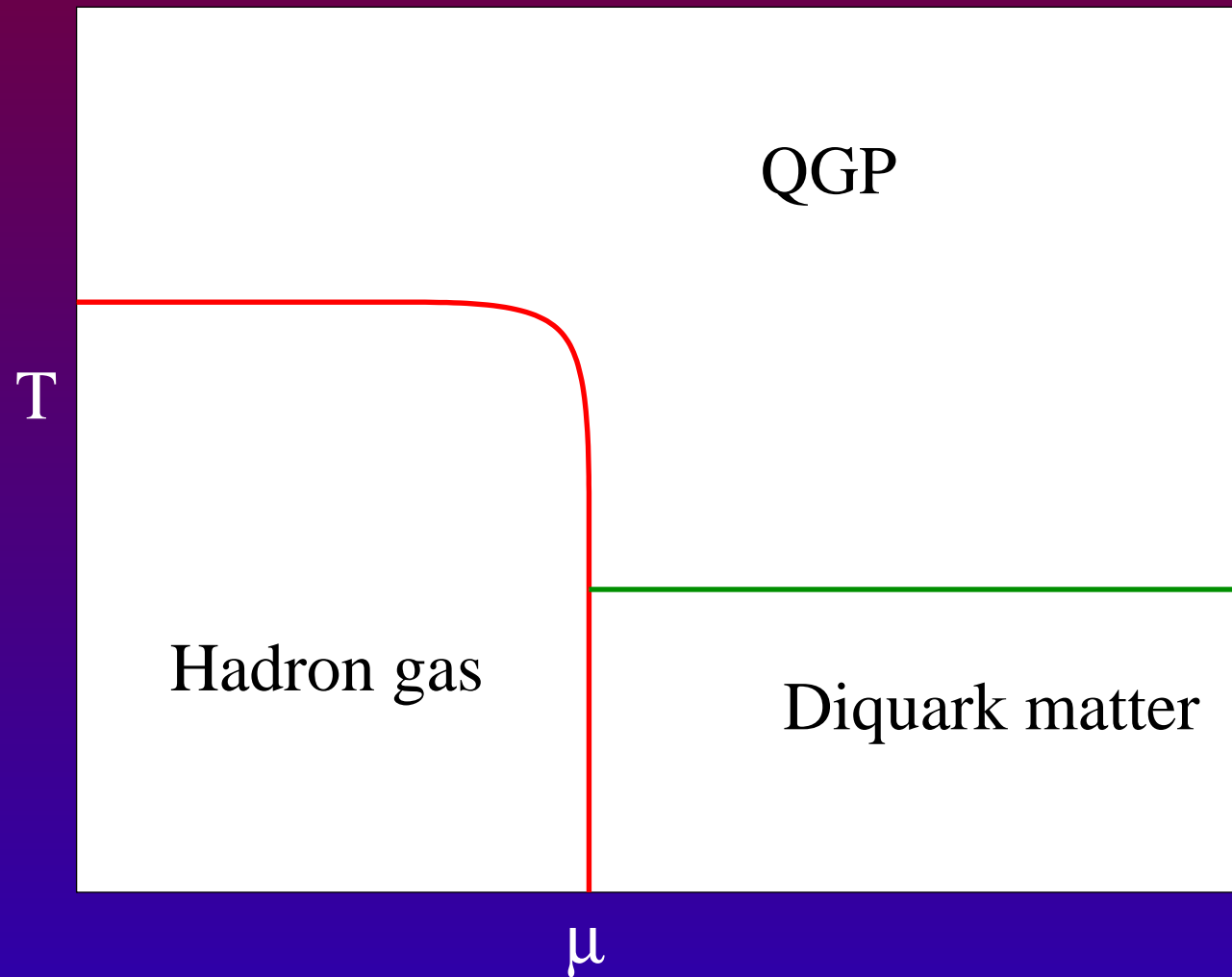
What happens when many hadrons are put in a small volume?

Deconfinement



Now quarks can propagate distances $\gg 1$ fm
Hadrons have melted away.....and the system is “Quark
gluon plasma”

Hadronic matter at finite T and μ



QCD thermodynamics at finite T, μ

$$\begin{aligned}\mathcal{Z} &= \text{Tr}[e^{-H/T}] \\ \text{free energy} &= -\frac{T}{V} \ln \mathcal{Z} \\ \text{energy density} &= \frac{T^2}{V} \frac{\partial \ln \mathcal{Z}}{\partial T} \\ \text{pressure} &= \frac{T}{V} \frac{\partial \ln \mathcal{Z}}{\partial V}\end{aligned}$$

QCD at finite T

$$\mathcal{Z} \propto \text{Tr}[e^{-H/T}] \longrightarrow \mathcal{Z}_{QCD} \propto \int D[G, \bar{\psi}, \psi] e^{-S}$$

$$S = \int_0^{\frac{1}{T}} dt \int d^3x \left[\frac{1}{4} F_{\mu\nu}^a F^{a\mu\nu} + \sum_f \bar{\psi}_\alpha^f (\gamma_E^\mu D_\mu + m_f)^{\alpha\beta} \psi_\beta^f \right]$$

$$\mathcal{Z} \propto \int DG \det(M) e^{-S_{gluon}} = \int DG e^{-S_{QCD}}$$

$S_{QCD} = S_{gluon} - \log(\det M)$. Changing temperature amounts to changing the extent of imaginary time direction.

QCD at finite T

$$\langle O(G) \rangle = \frac{\int DG O(G) e^{-S_{QCD}}}{\int DG e^{-S_{QCD}}}$$

Non-perturbative method, **lattice QCD** is the only alternative

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Suppose the system makes a random walk in $A(x)$ -space.....

$$G_1(x) \rightarrow G_2(x) \rightarrow G_3(x) \rightarrow G_4(x) \dots \dots \dots$$

$$\langle O(G)_{ap} \rangle = \frac{\sum_{i=1}^N O(G_i)}{N}$$

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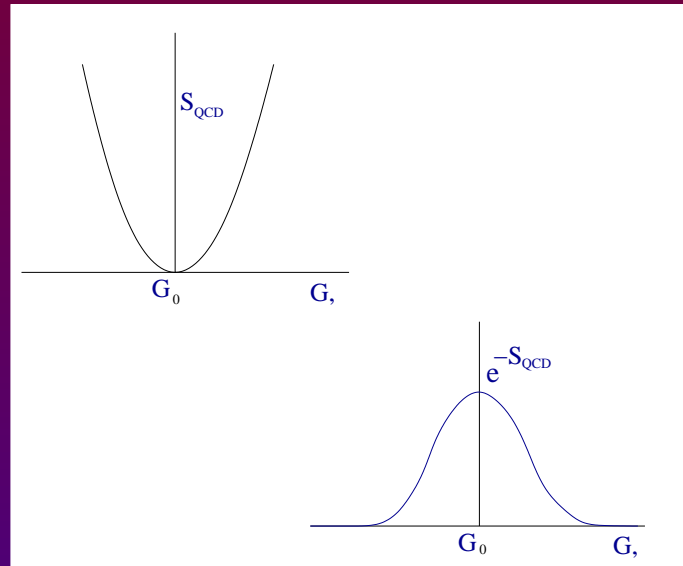
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$$\langle O(A)_{ap} \rangle \neq \langle O(A) \rangle$$

Lattice QCD



If in the random walk the system spends more “time” near G_0 if histogram $H(G)$follows $e^{-S_{\text{QCD}}}$ then

$$\langle O(G) \rangle \equiv \frac{\int DG O(G) e^{-S_{\text{QCD}}}}{\int DG e^{-S_{\text{QCD}}}} = \frac{\sum_{i=1}^N O(G_i)}{N}$$

Lattice QCD

- How to generate N number of G 's with probability $e^{-S_{QCD}}$?

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Accept G' with probability $\exp(-\Delta S)$

$$G \rightarrow G' \rightarrow G'' \rightarrow G''' \dots\dots$$

$$G(x, y, z, t) \rightarrow G'(x, y, z, t)$$

$$G(x_i, y_i, z_i, t_i) \rightarrow G'(x_i, y_i, z_i, t_i)$$

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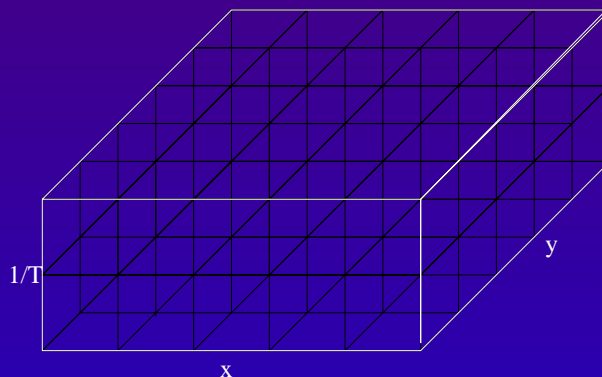
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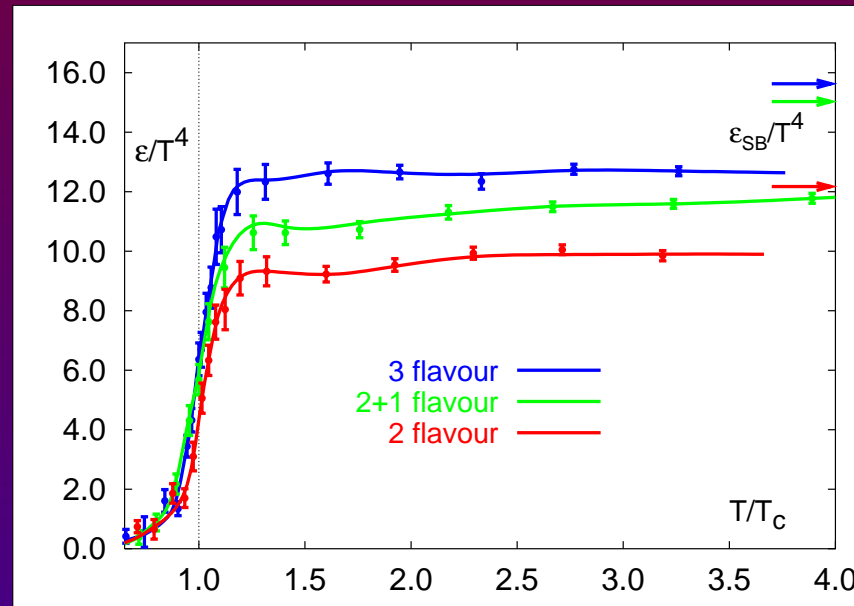
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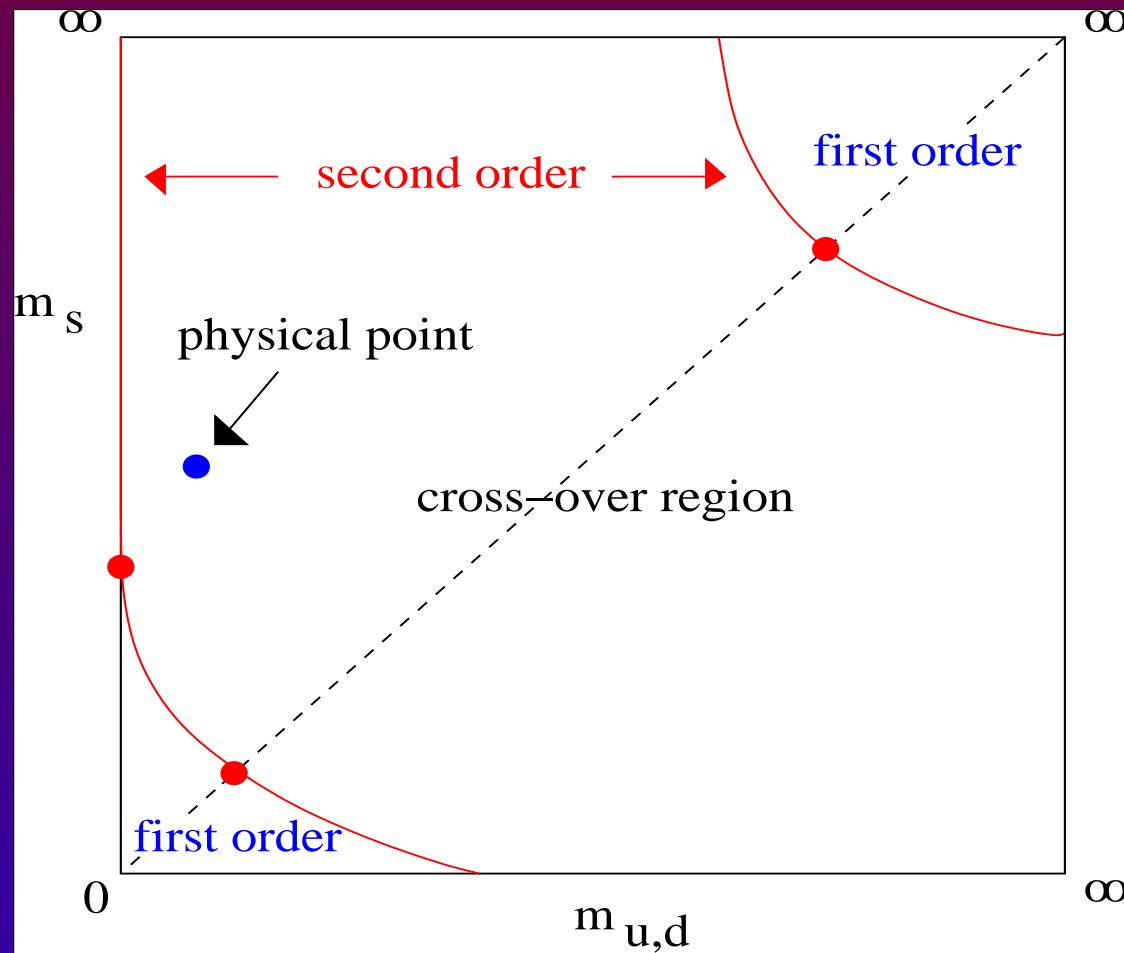
QCD at finite T



A sharp increase in the number of degrees of freedom
 \Rightarrow Melting of hadrons \Rightarrow Color deconfinement

This is the only theoretical evidence of deconfinement,
success of non-perturbative methods

QCD Phase Diagram ($\mu = 0$)



QCD Phase Diagram

