

Strong subadditivity in quantum information



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Outline

- What is strong subadditivity?
- SSA and the Holevo bound
- SSA and multipoint quantum dense coding
- SSA and resonating valence bond states

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von Neumann entropy

- arb quantum state ρ
- von Neumann entropy of ρ , denoted $S(\rho)$, given by

$$S(\rho) = -\text{tr} (\rho \log \rho)$$

Subadditivity

- $S(\rho_{AB}) \leq S_A + S_B$

Strong subadditivity

- $S_{ABC} + S_A \leq S_{AB} + S_{AC}$

Strong subadditivity

- $S_{ABC} + S_A \leq S_{AB} + S_{AC}$

OR

- $S_B + S_C \leq S_{AB} + S_{AC}$

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Lieb & Ruskai, J Math Phys 1973

Strong subadditivity

- $S_{ABC} + S_A \leq S_{AB} + S_{AC}$

OR

- $S_B + S_C \leq S_{AB} + S_{AC}$

- Both versions are equivalent & imply subadditivity.

Strong subadditivity

- $S_{ABC} + S_A \leq S_{AB} + S_{AC}$

OR

- $S_B + S_C \leq S_{AB}$


$$r_{ABC} = r_A \otimes r_{BC}$$

- Both versions r equivalent & imply subadditivity.

Strong subadditivity

- $S_{ABC} + S_A \leq S_{AC}$

OR

For version 2 to subadditivity, first go to version 1.

- $S_B + S_C \leq S_{AB} + S_{AC}$

- Both versions are equivalent & imply subadditivity.

Outline

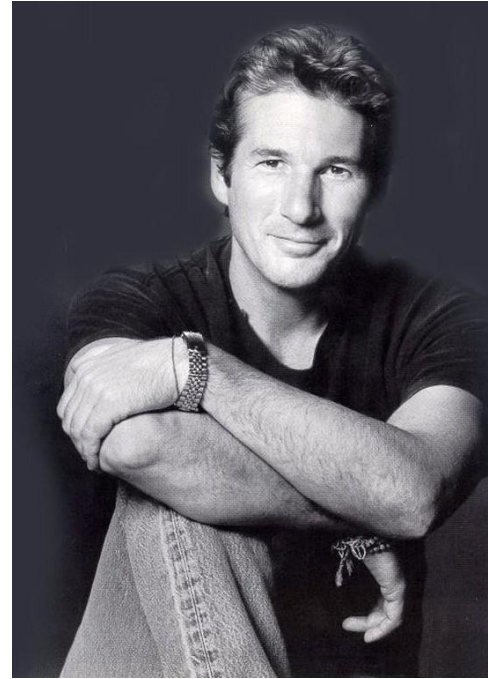
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- SSA and resonating valence bond states

Encoding classical information in quantum states

Encoding classical information in quantum states

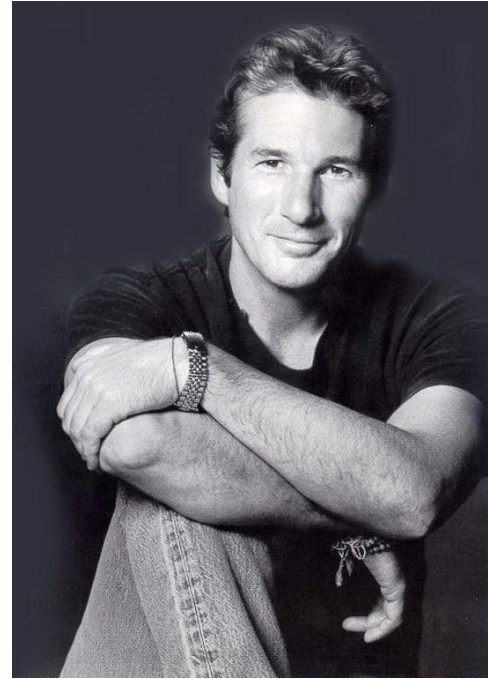


Encoding classical information in quantum states

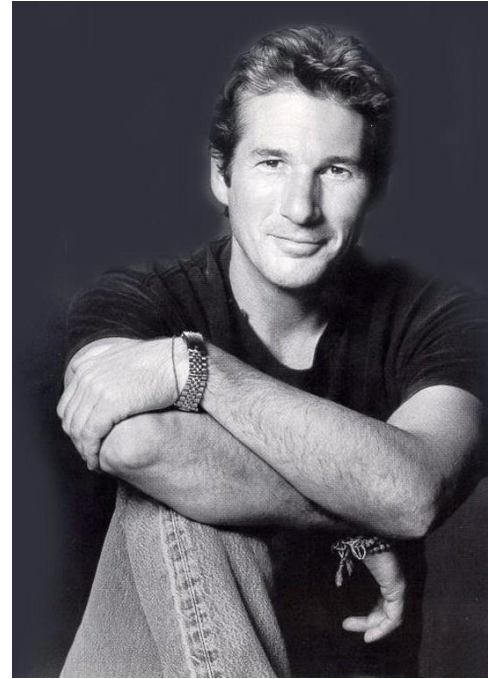




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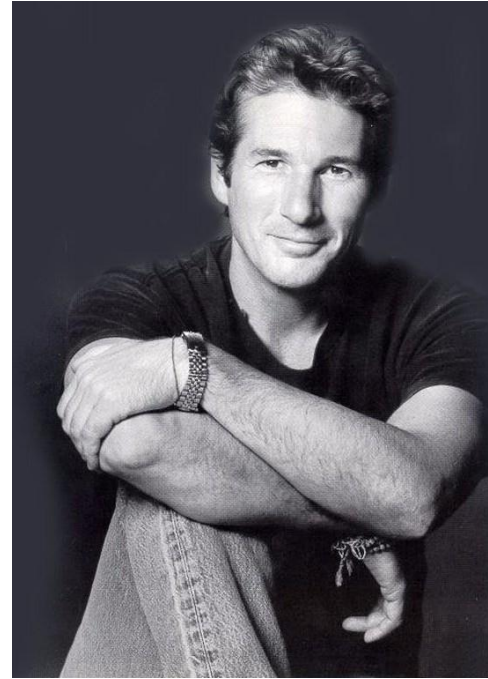


Encoding classical information in quantum states



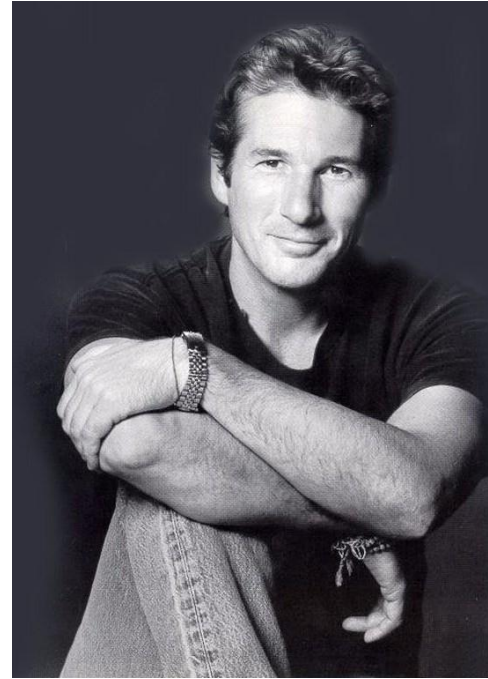
- Alice (sender) encodes the classical variable x , that occurs with probability p_x , in the quantum state ρ_x . I.e., she produces the ensemble $\mathcal{E} = \{p_x, \rho_x\}$.

Encoding classical information in quantum states



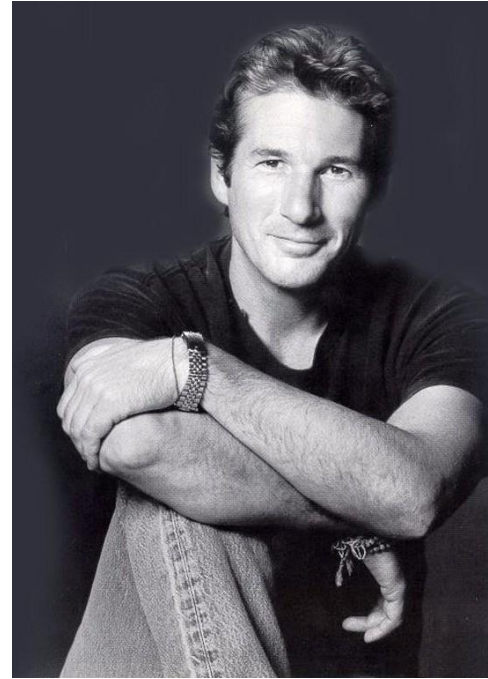
- Alice (sender) encodes the classical variable x , that occurs with probability p_x , in the quantum state ρ_x . I.e., she produces the ensemble $\mathcal{E} = \{p_x, \rho_x\}$.
- She then sends it to Bob (receiver).

Encoding classical information in quantum states



- Task of Bob: Gather information about x .

Encoding classical information in quantum states



- Task of Bob: Gather information about x .
- Accessible information = Maximal classical information that can be extracted by Bob from $\mathcal{E} = \{\rho_x, \rho_x\}$.

Holevo theorem (1973): Upper bound on accessible information

- Initial Ensemble: $\mathcal{E} = \{p_x, \rho_x\}$
- Accessible information $\leq \chi$ ← Holevo quantity
- $\chi = S(\rho) - \sum p_x S(\rho_x)$; $\rho = \sum p_x \rho_x$.
- Accessible information $\leq S(\rho) \leq \log_2 d$

Lower bound: Jozsa, Robb, & Wootters, Phys. Rev. A '94

Outline of the proof of Holevo bound

- Usually involved.

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- But, becomes uncluttered with use of SSA.

Schumacher, Westmoreland, Wootters, PRL 1996

Outline of the proof of Holevo bound

- Usually involved.
- But, becomes uncluttered with use of SSA.
- Essential step is to prove that $\chi_{AB} \geq \chi_A$, and this is where SSA is used.

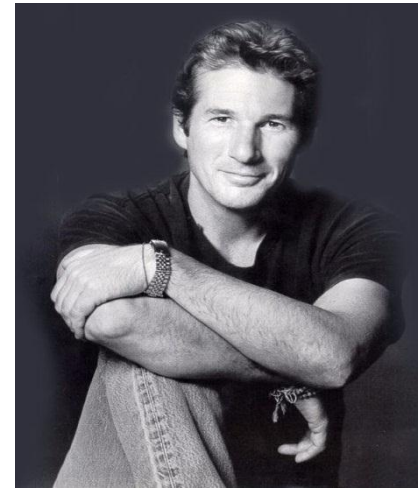
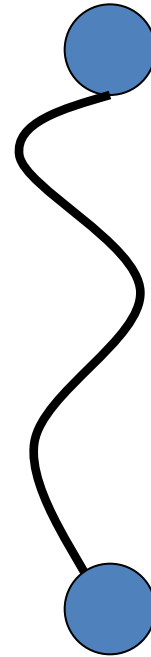
Schumacher, Westmoreland, Wootters, PRL 1996

LOCC counterpart of Holevo theorem:

Upper bound on
locally accessible info



$$\mathcal{E} = \{p_x, \rho_x\}$$



Entanglement correction to Holevo theorem

- Accessible information $\leq \log_2 d$
- Locally accessible information $\leq \log_2 d - E_{av}$

Universally true: Holds for arbitrary bipartite ensembles.

SSA plays an imp role, once again.

Badziag, Horodecki, Sen(De), & Sen, Phys. Rev. Lett. '03

Sen(De), Sen, & Lewenstein, Phys. Rev. A '06

Outline

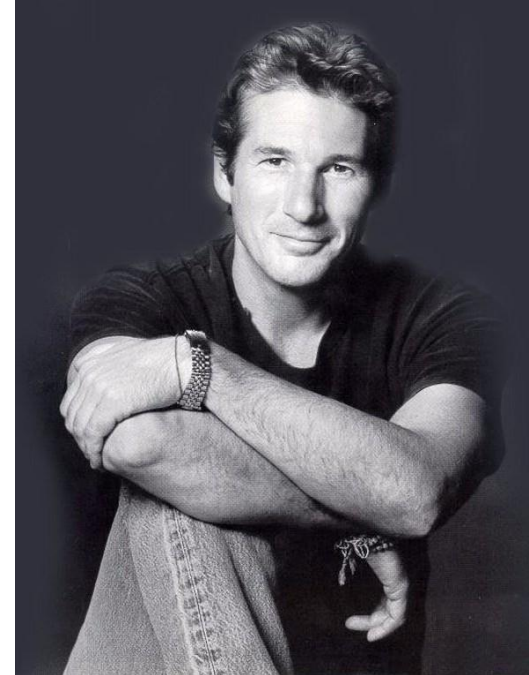
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Quantum dense coding



Alice

Sender

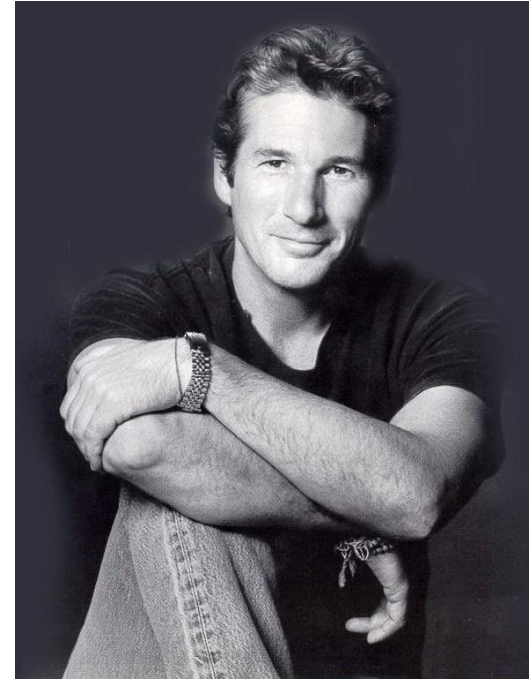


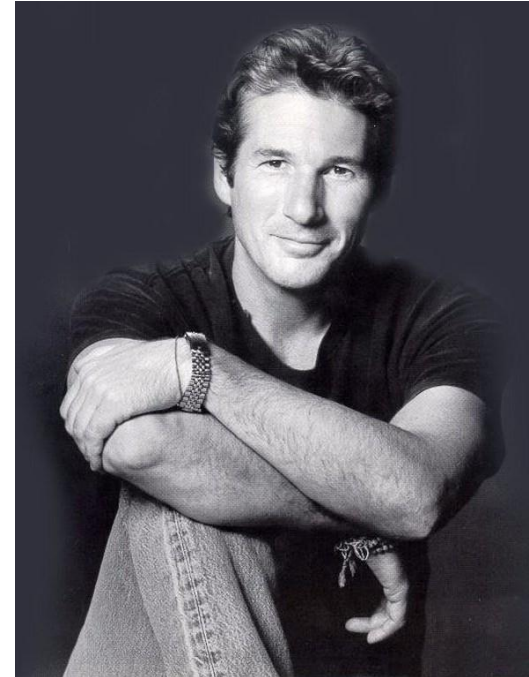
Bob

Receiver

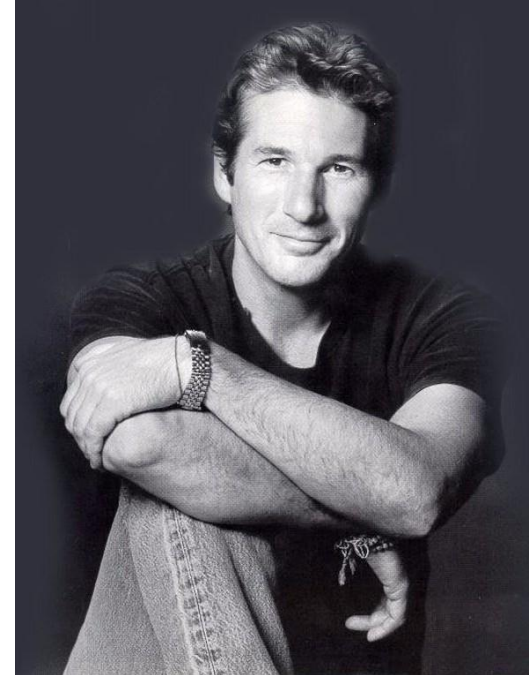


Alice is in Bhubaneswar.
Bob is in Delhi.





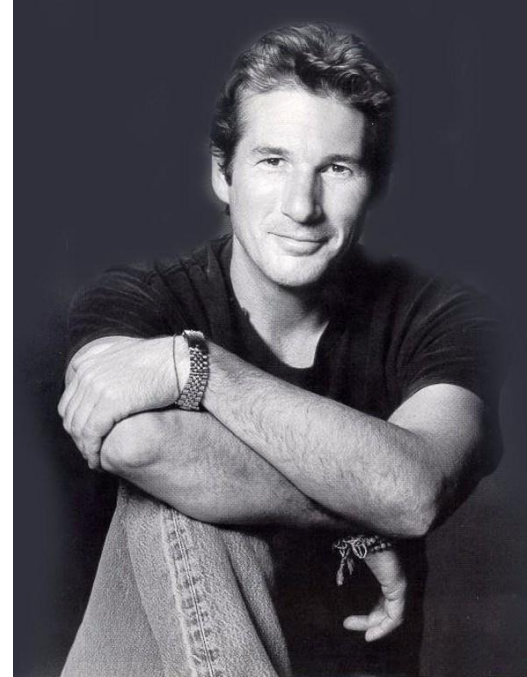
Alice wants to send info about weather in Bhubaneswar to Bob.



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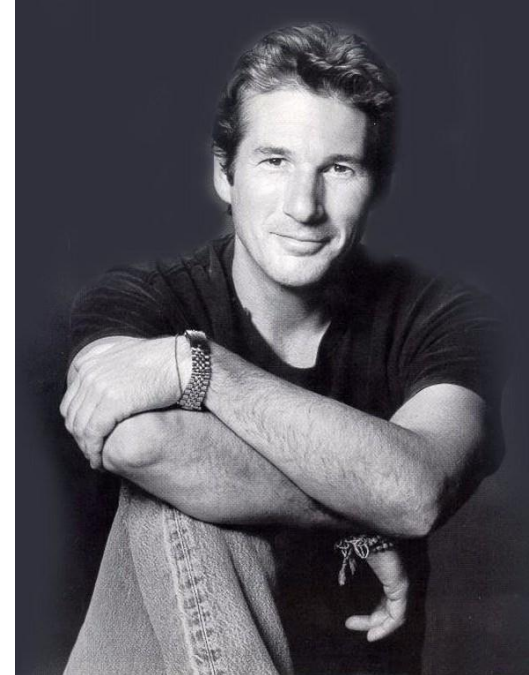
Sunny or not

Windy or not



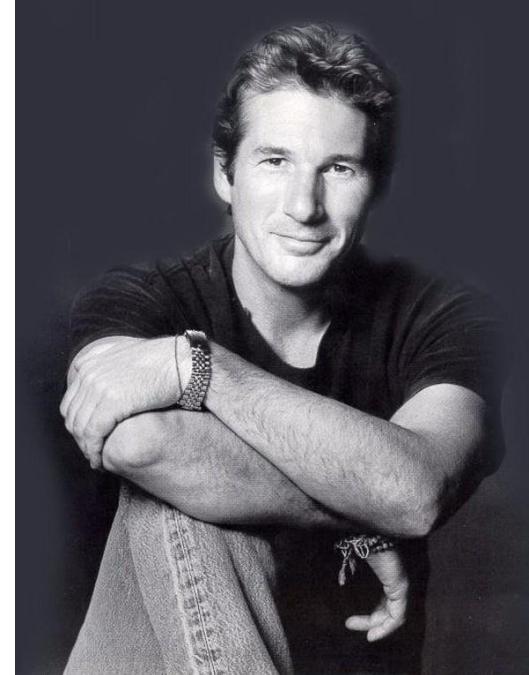
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2 bits



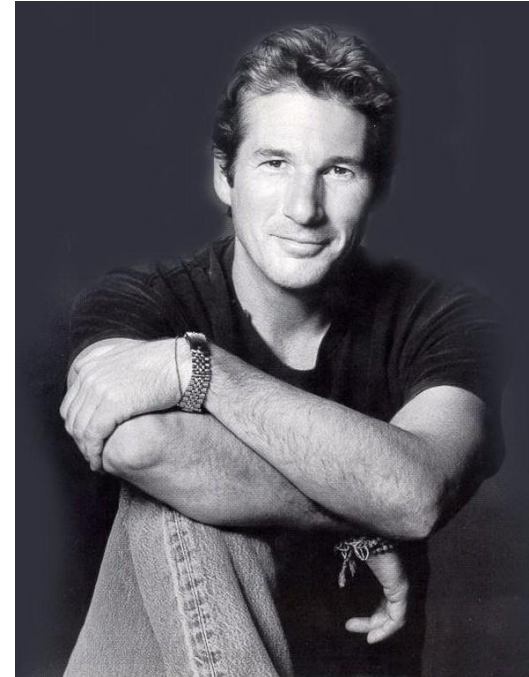
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Can be sent by using ...

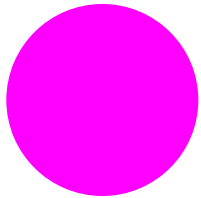


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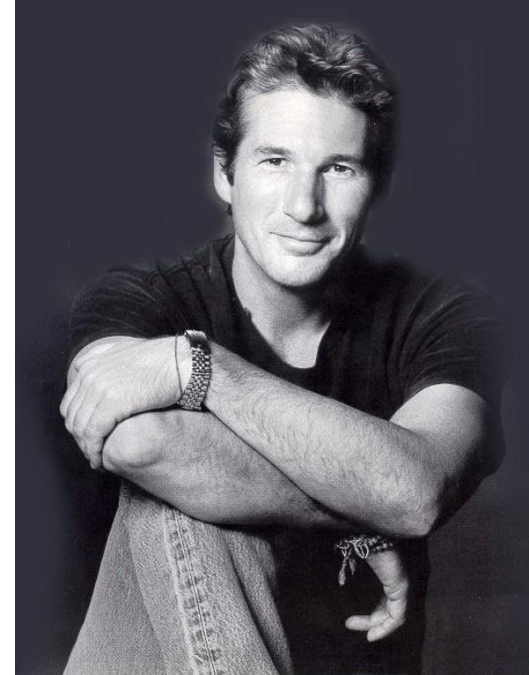
Four balls of different colors.



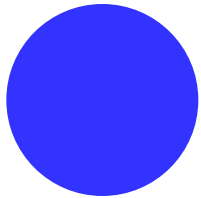
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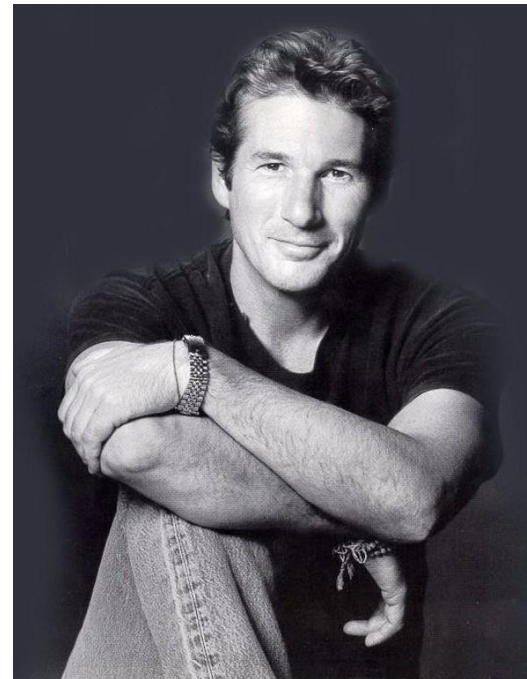
sunny and not windy



Alice wants to send info about weather in Bhubaneswar to Bob.

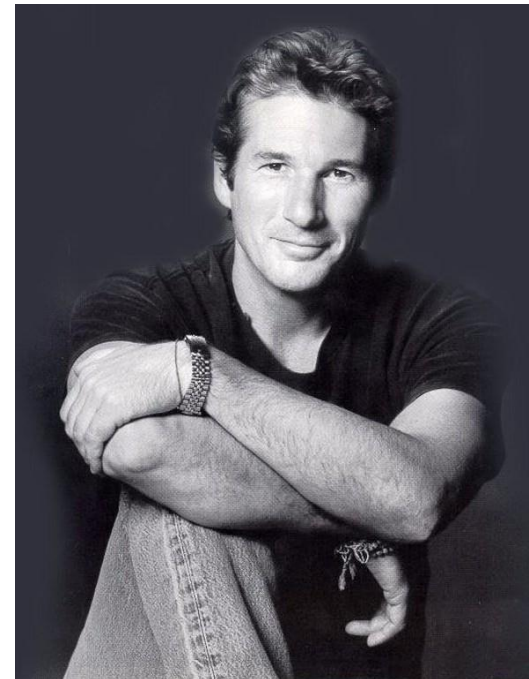


sunny and windy



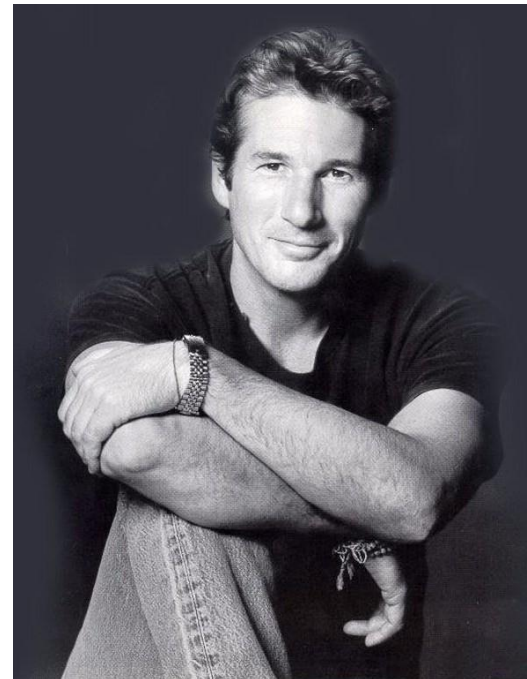
Alice wants to send info about weather in Bhubaneswar to Bob.

2 bits require 4 dim.



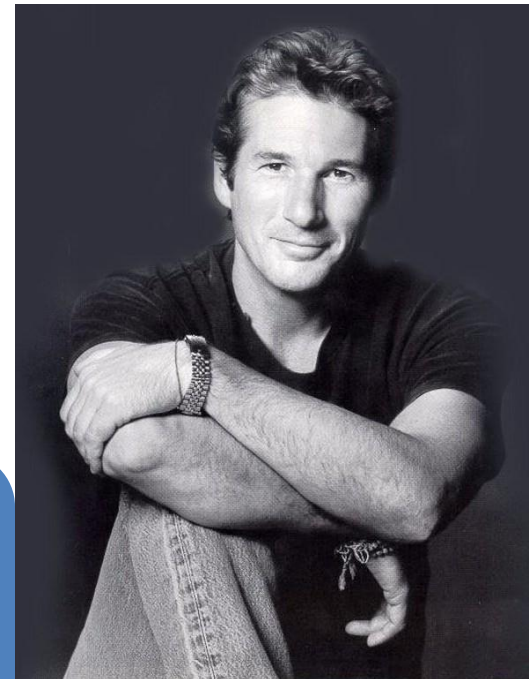
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Using shared entanglement between Alice & Bob, ...



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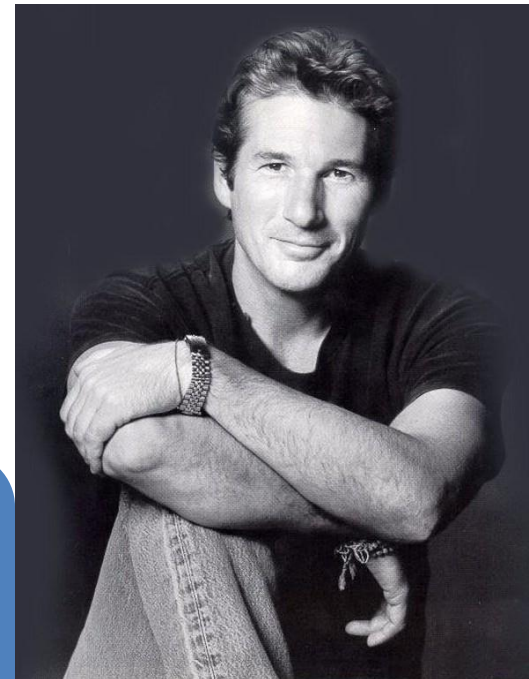
Using shared entanglement between Alice & Bob, ...



$$|\uparrow\rangle\otimes|\downarrow\rangle - |\downarrow\rangle\otimes|\uparrow\rangle$$

Alice wants to send info about weather in Bhubaneswar to Bob.

Using shared entanglement between Alice & Bob, ...

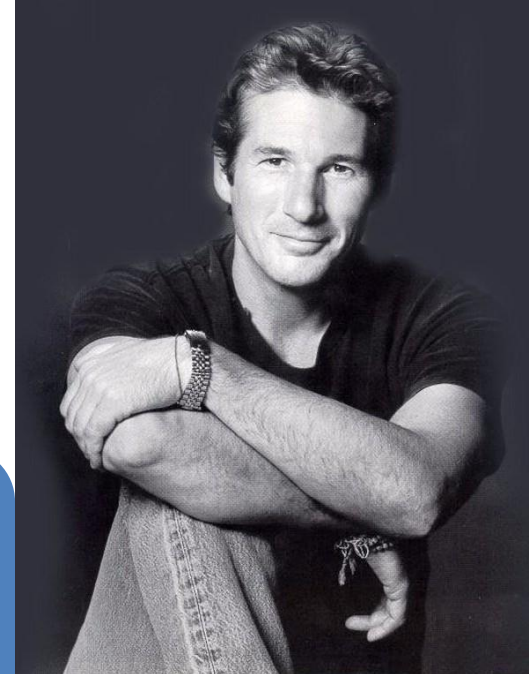


$|\uparrow\rangle \otimes |\downarrow\rangle - |\downarrow\rangle \otimes |\uparrow\rangle$

Alice wants to send info about weather in Bhubaneswar to Bob.

2 bits require 2 dim.

Quantum dense coding

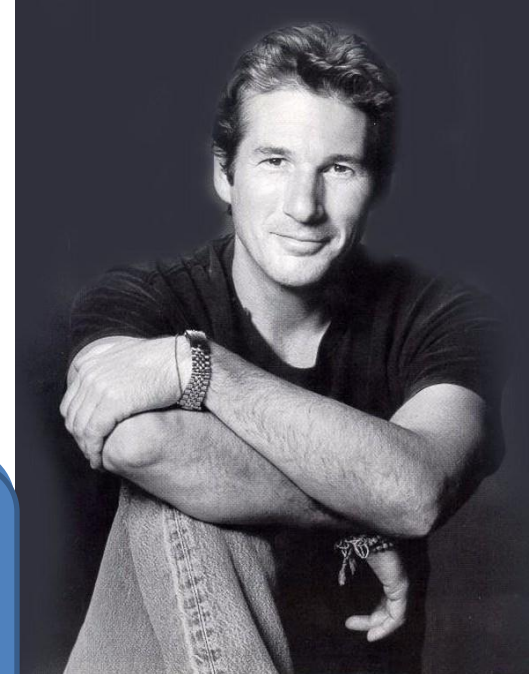


$$|\uparrow\rangle \otimes |\downarrow\rangle - |\downarrow\rangle \otimes |\uparrow\rangle$$

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2 bits require **2 dim.**

Quantum dense coding

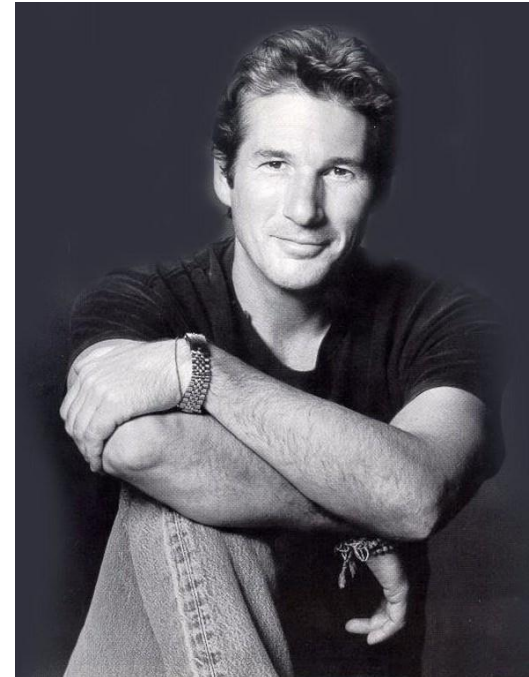
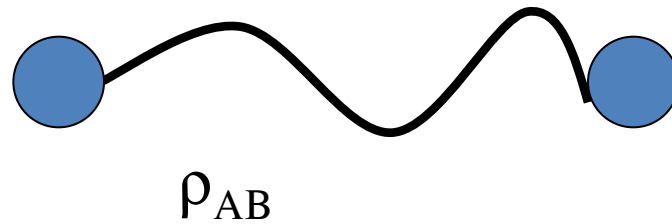


Bennett & Wiesner
1992

Alice wants to send info about weather in Bhubaneswar to Bob.

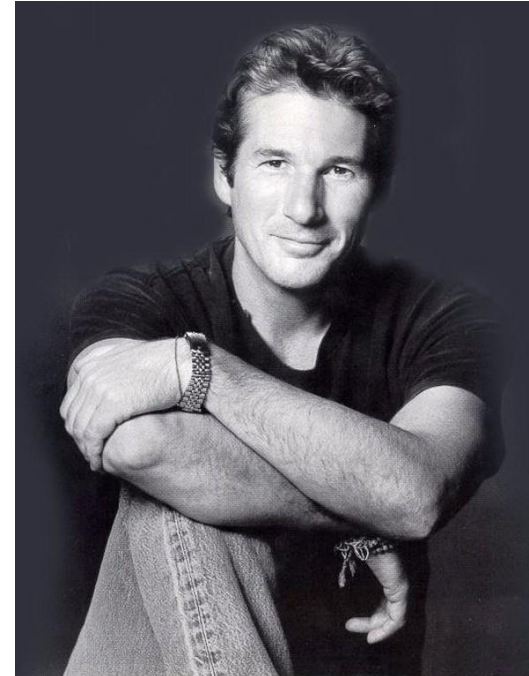
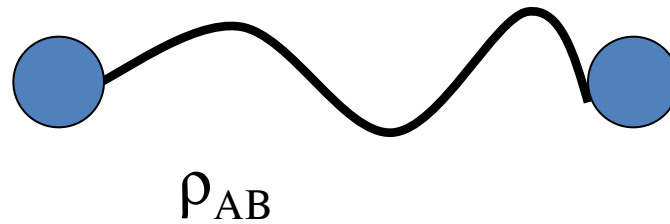
2 bits require **2 dim.**

Capacity of quantum dense coding



- Capacity of quantum dense coding = amount of classical info that can be sent via a given shared state.
- The Holevo theorem can be used to find this capacity.

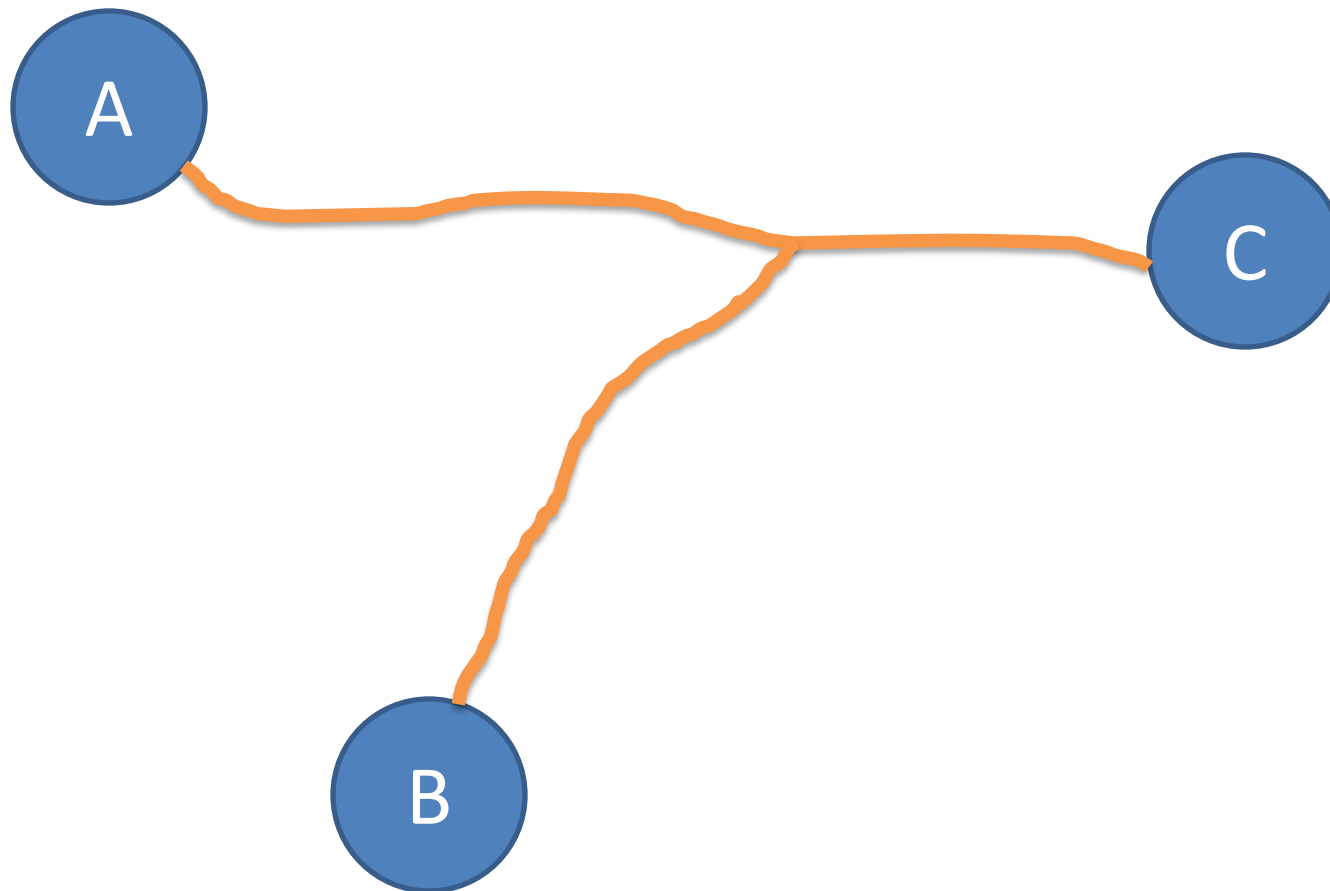
Capacity of quantum dense coding



- Capacity of quantum dense coding =

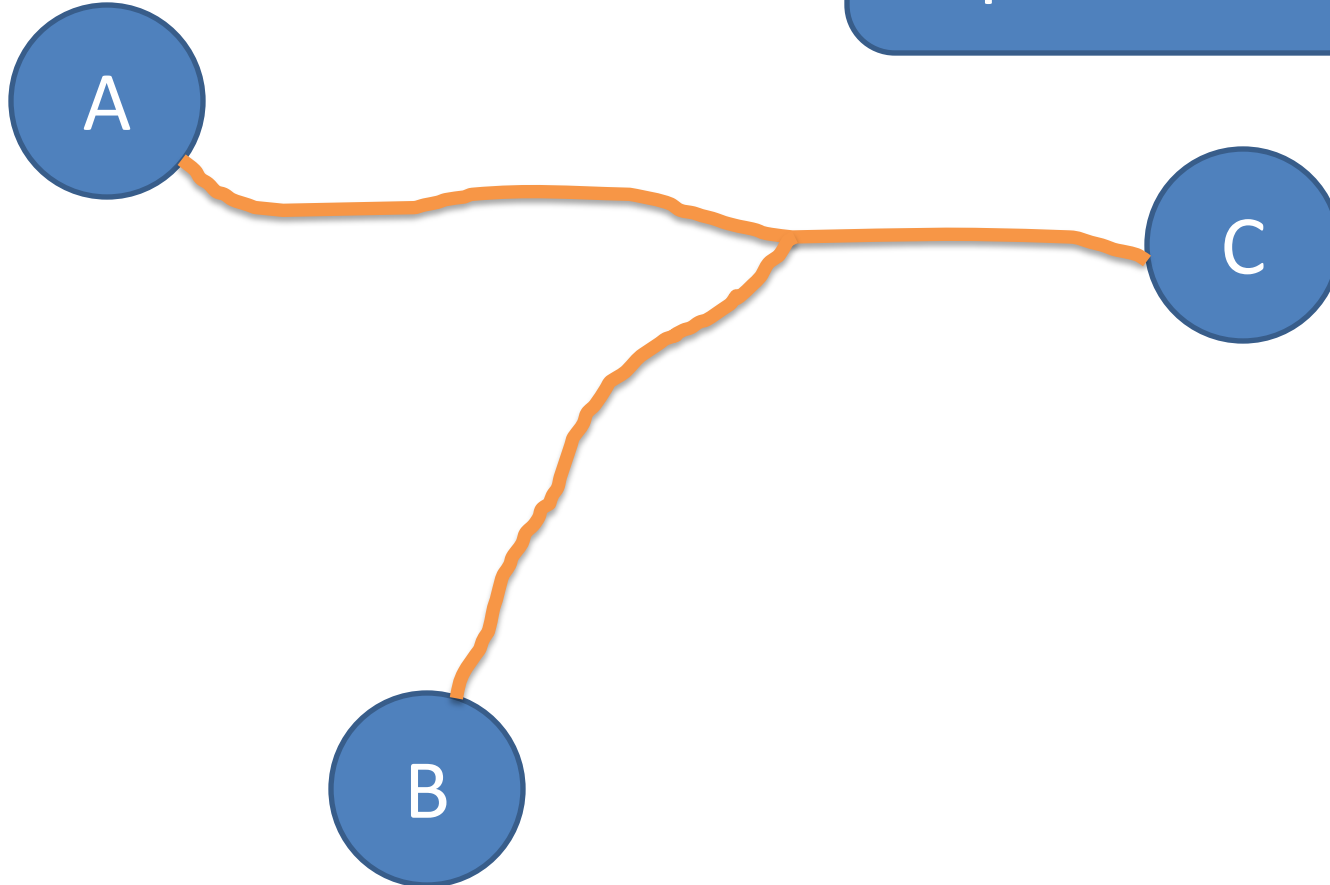
$$\log d + S_B - S_{AB}$$

Distributed Q Dense Coding



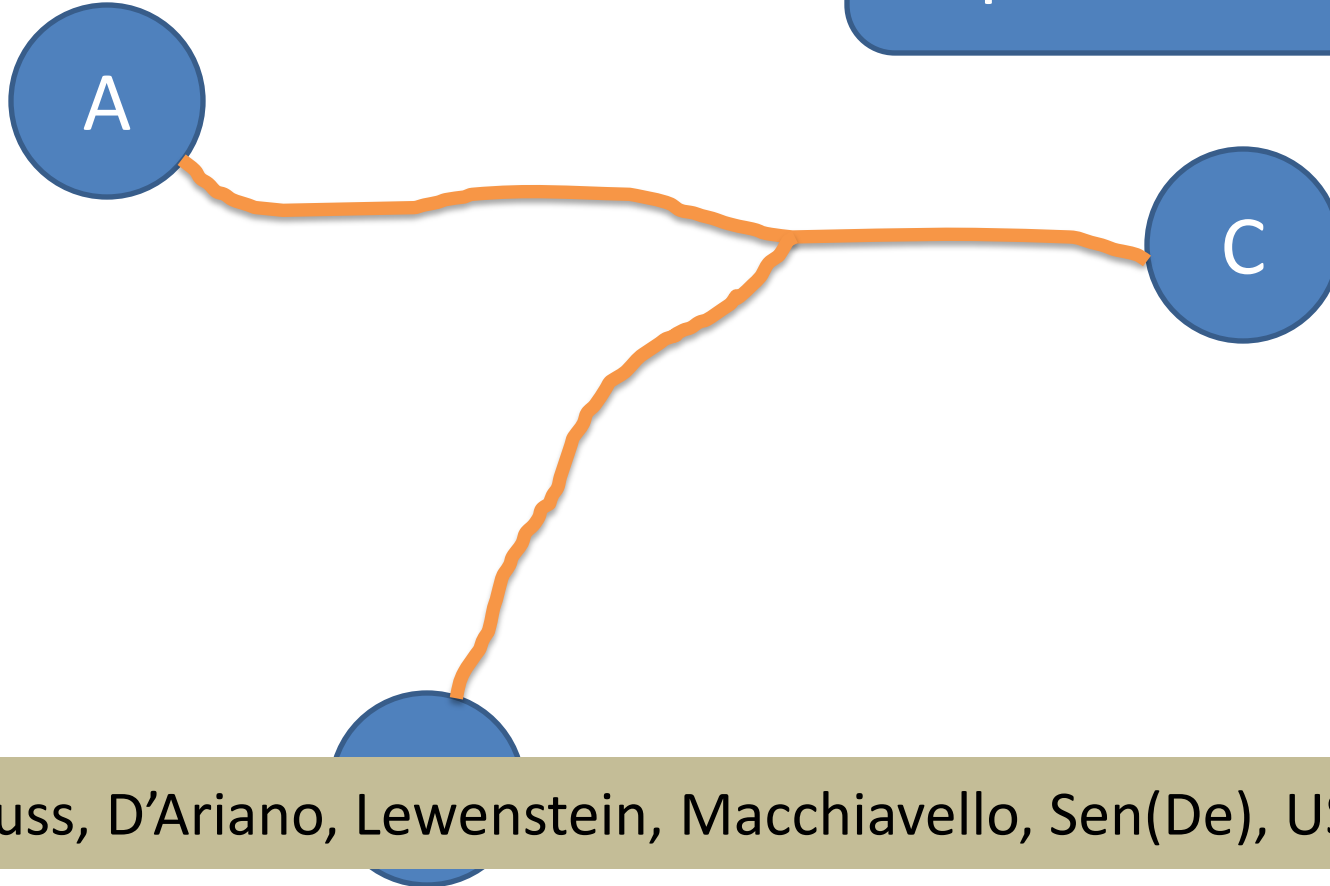
Distributed Q Dense Coding

Towards a
quantum internet



Distributed Q Dense Coding

Towards a
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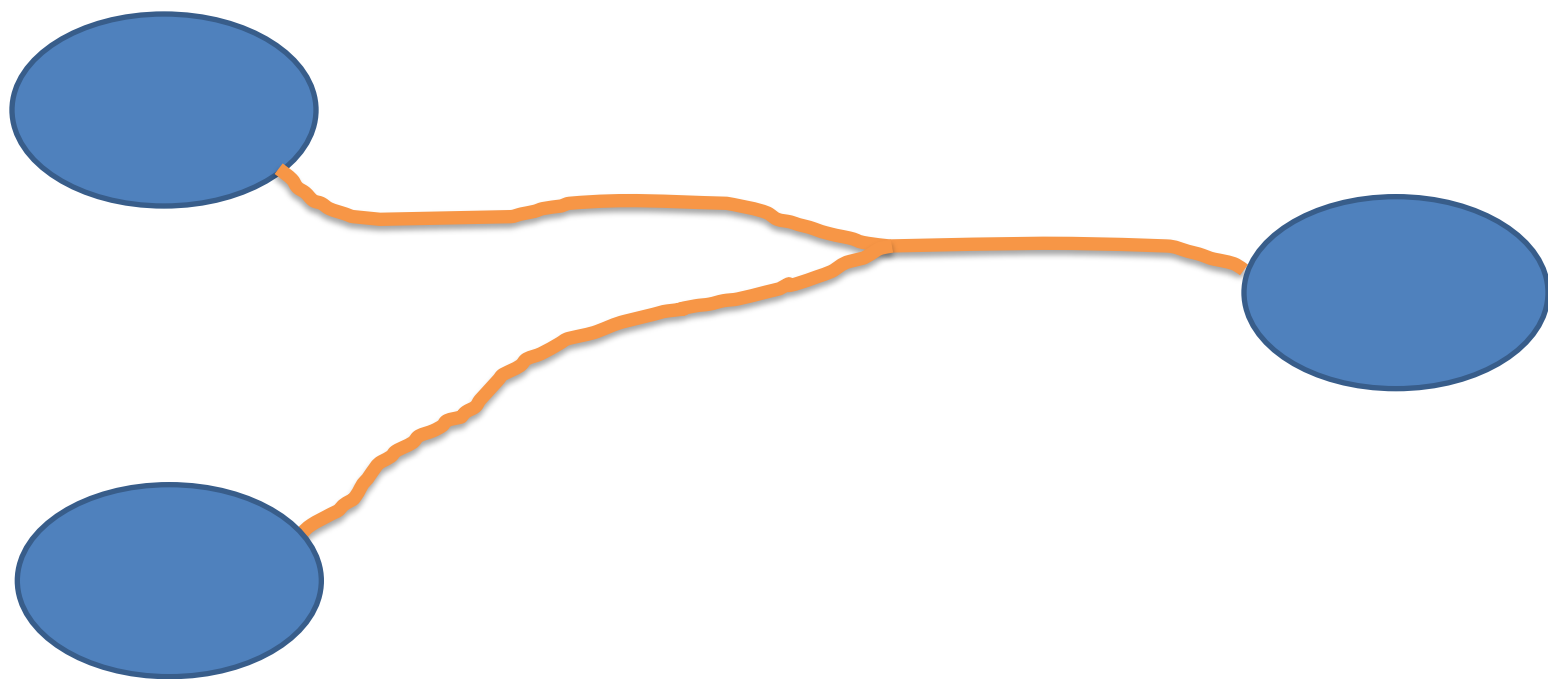


Bruss, D'Ariano, Lewenstein, Macchiavello, Sen(De), US, PRL 2004

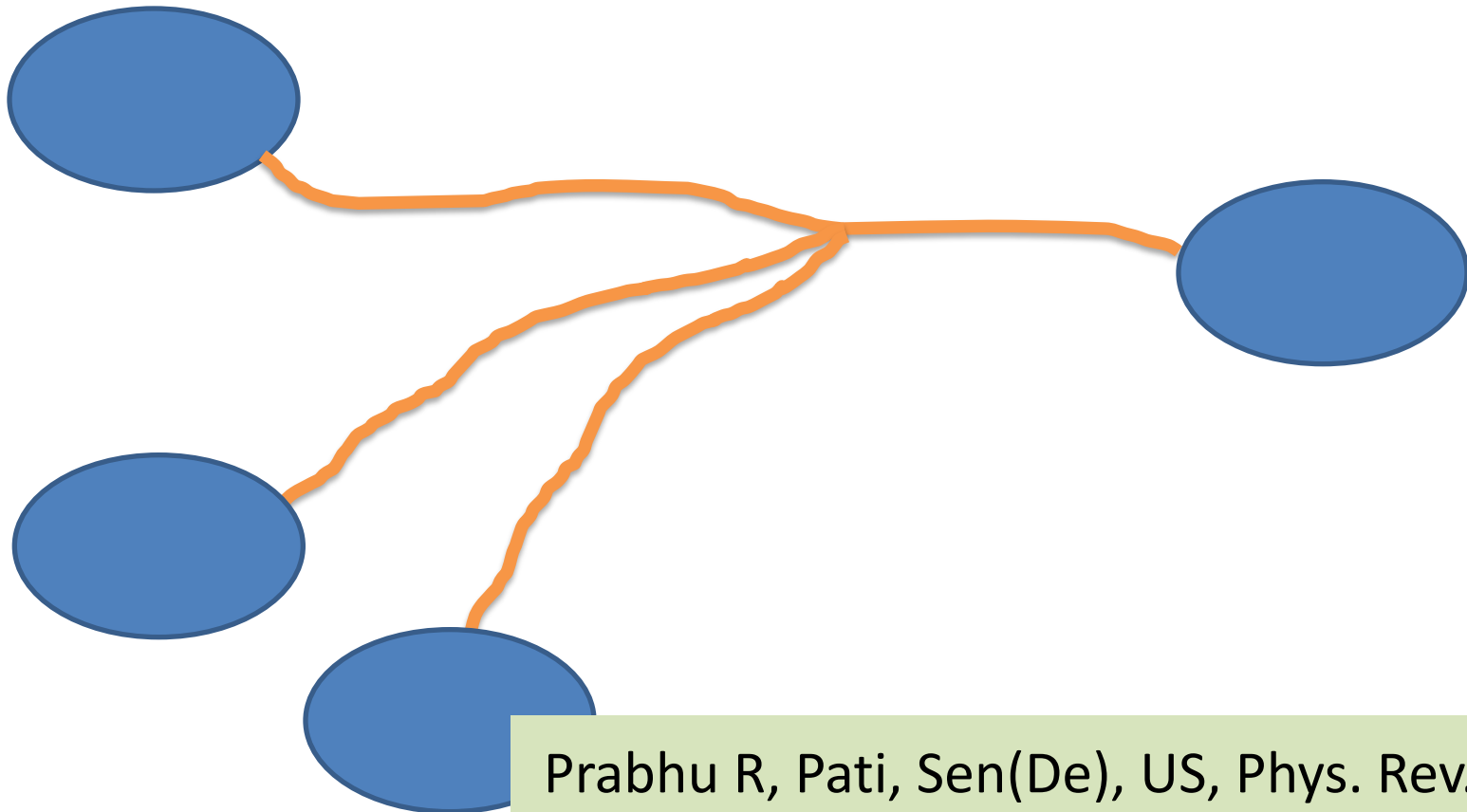
Exclusion pple for Q Dense Coding

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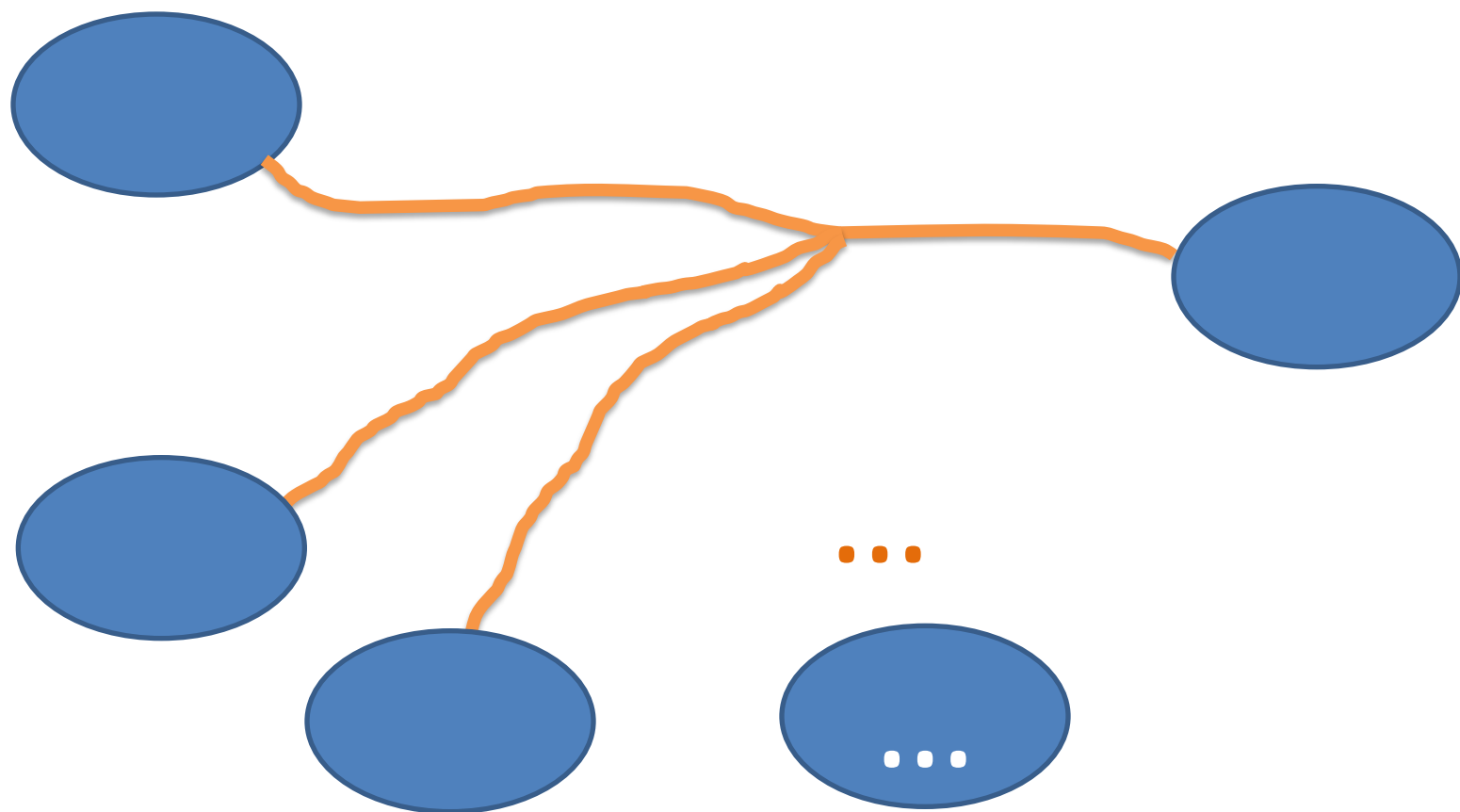


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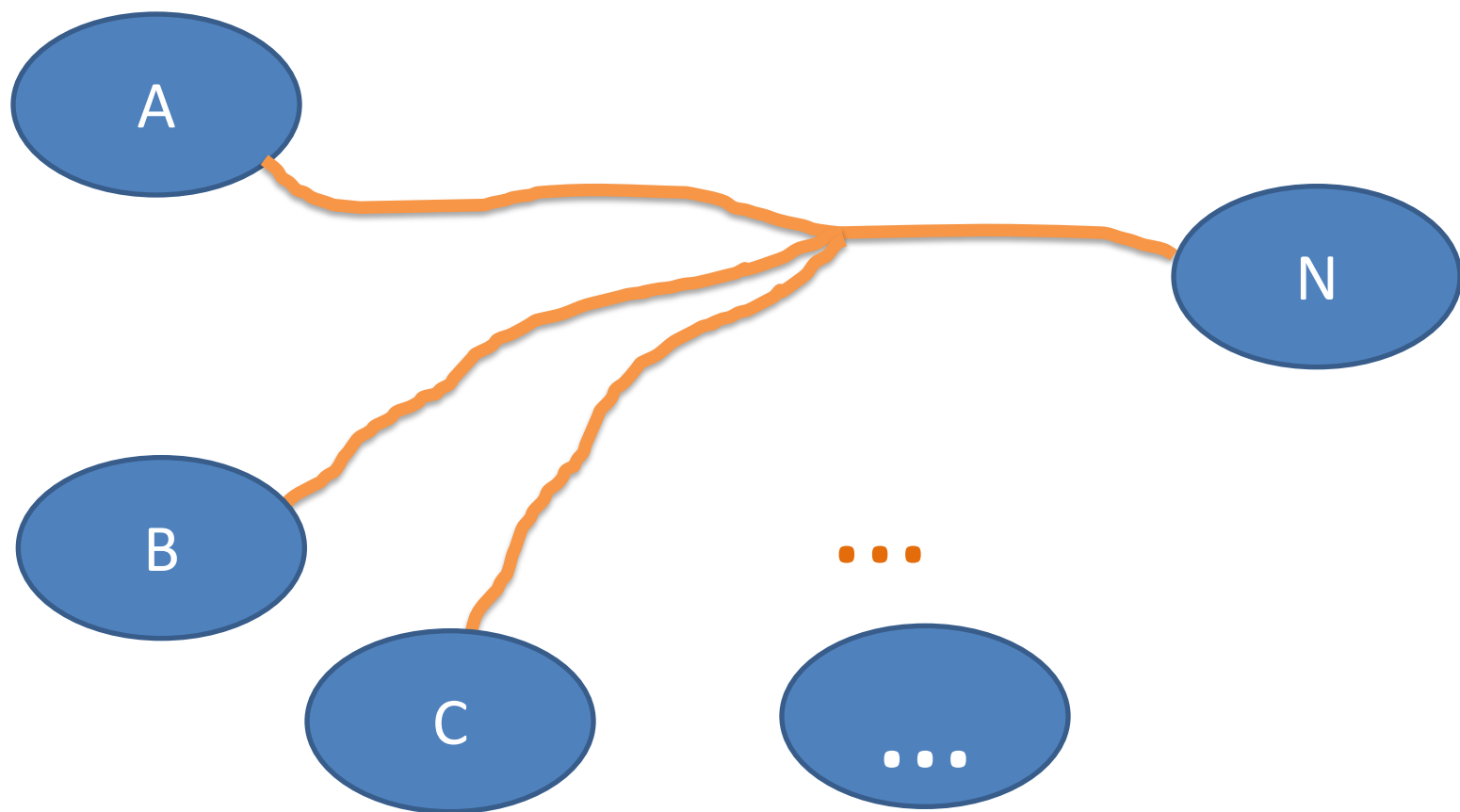


Prabhu R, Pati, Sen(De), US, Phys. Rev. A 2013

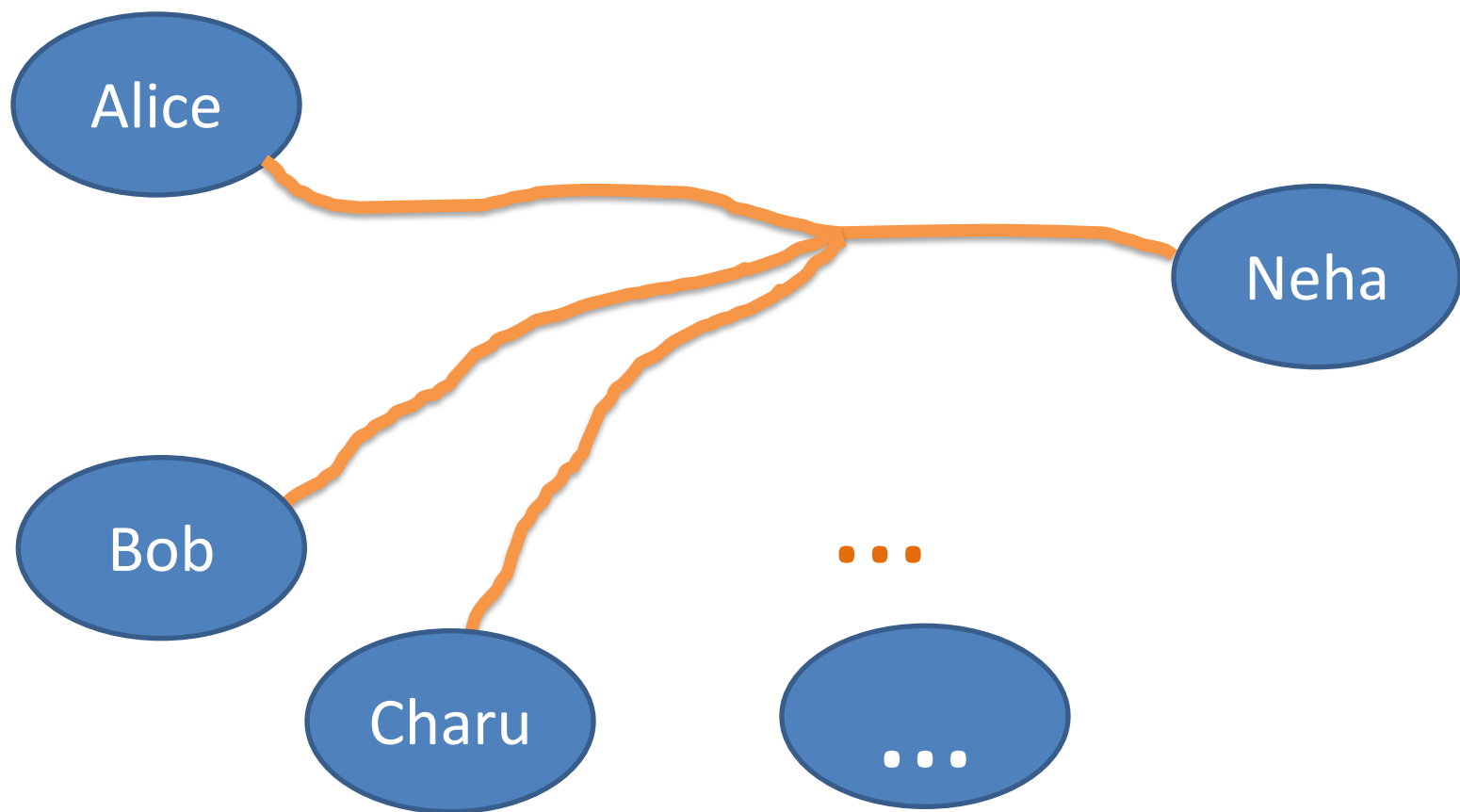
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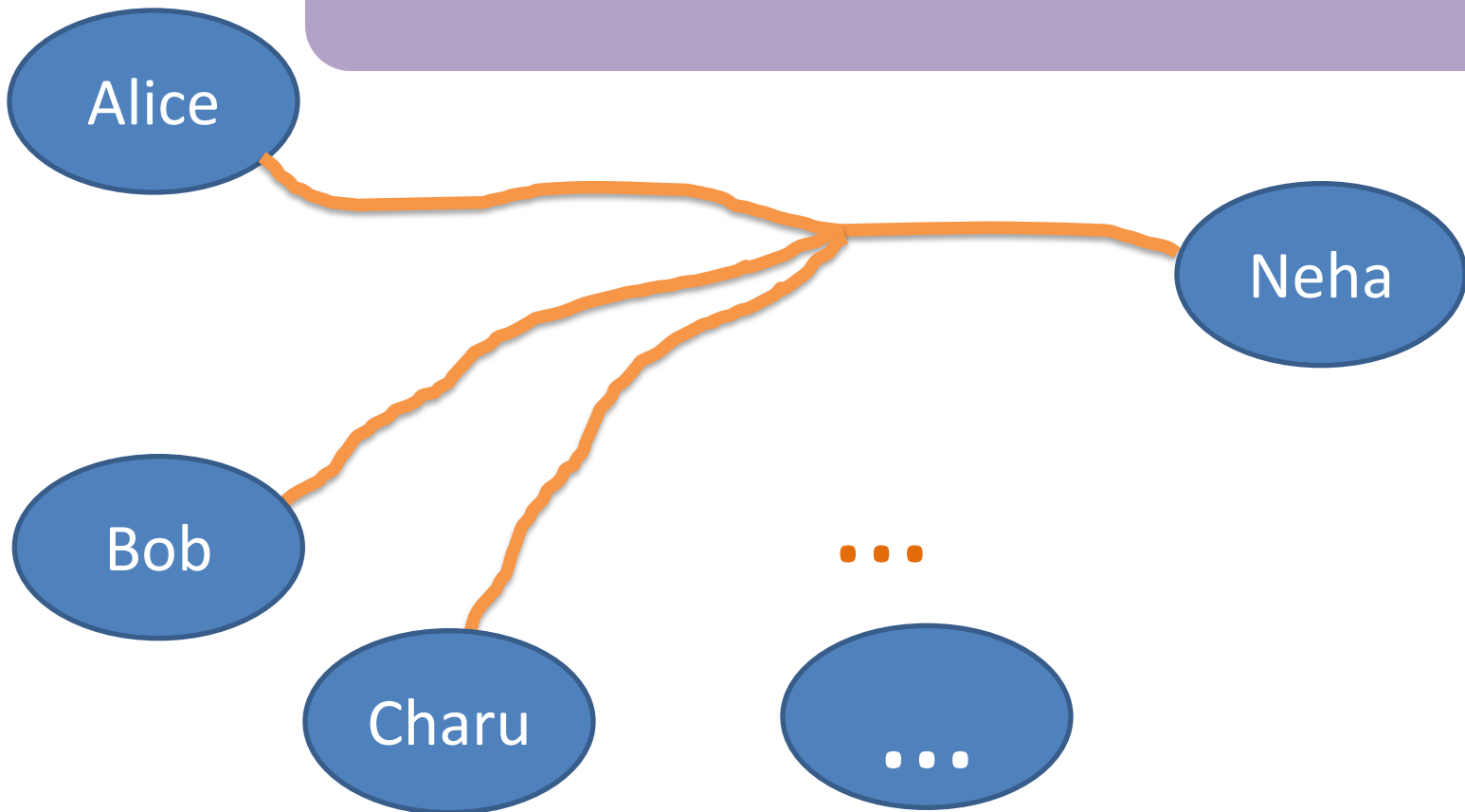


Exclusion pple for Q Dense Coding



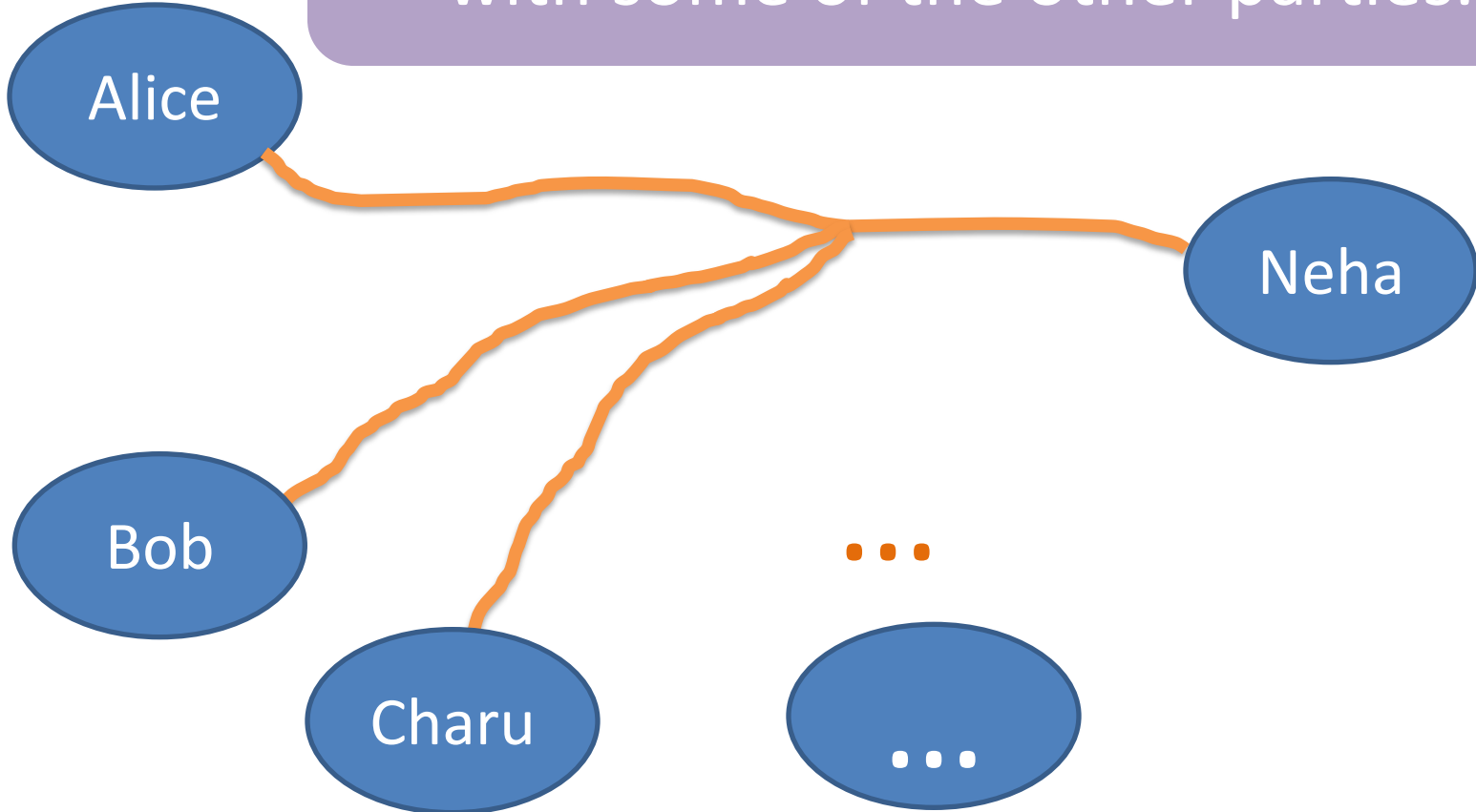
Exclusion pple for Q Dense Coding

N-party quantum state shared.



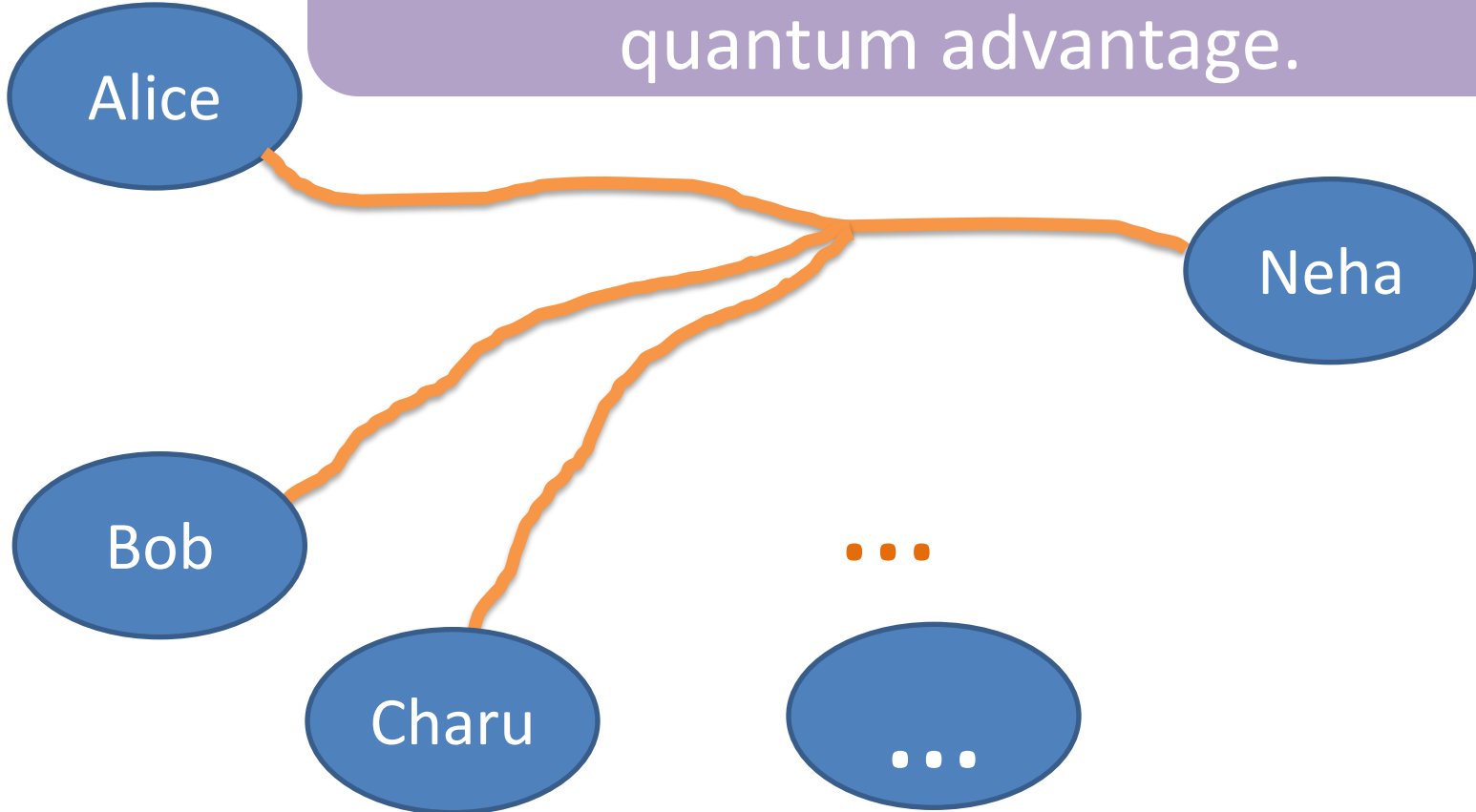
Exclusion pple for Q Dense Coding

Alice wishes to perform dense coding with some of the other parties.



Exclusion pple for Q Dense Coding

For every shared multiparty q state, **at most one** channel from Alice has a quantum advantage.



Exclusion principle for Q Dense Coding

For every shared multiparty q state, **at most one** channel from Alice has a quantum advantage.

Alice

Only two options possible:

C C C C

OR

Q C C C

Charu

...

Exclusion principle for Q Dense Coding

Note that
albeit of a
shown to b

This monogamy is stricter than that of quantum correlations. W state have qc in AB **and** AC for whatever qc u may choose!

Only two options possible:

C C C C

OR

Q C C C

Charu

Outline of proof for 3-party states

- C_{AB} is quantum if $S_B - S_{AB} \geq 0$.
- This is becoz $C_{AB} = \log d + S_B - S_{AB}$,
and $\log d$ can be achieved by classical means.
- And, C_{AC} is quantum if $S_C - S_{AC} \geq 0$.
- But, SSA dictates that sum of the LHSs ≤ 0 .
- So, either both ≤ 0 or only one ≥ 0 .

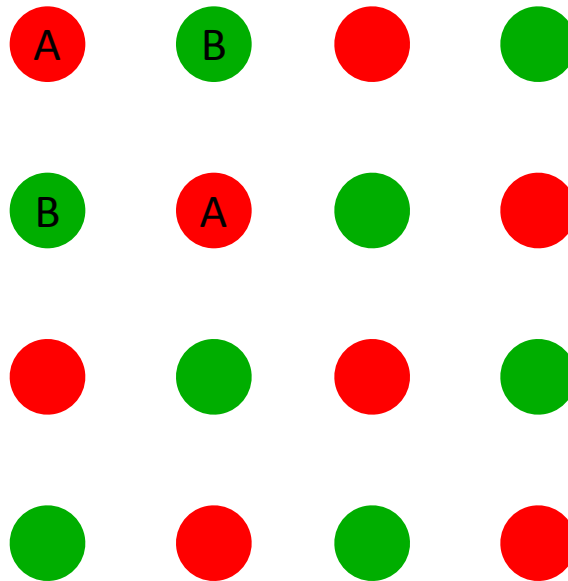
- For larger # of parties, suppose that two (or more channels) r quantum.
- Eg. let AB and AC be quantum for $\rho_{ABC \dots N}$.
- Then ρ_{ABC} is such that AB and AC r quantum.
- This contradicts the 3-party result.

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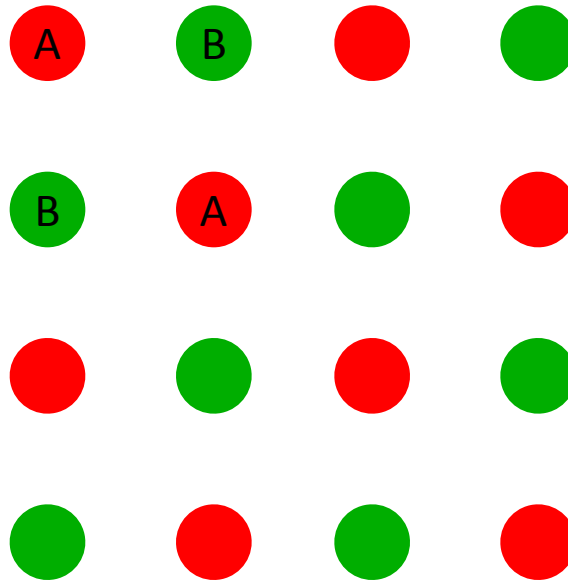
Resonating-Valence-Bond states

- $|\psi\rangle = \sum h(i_1, \dots, i_N; j_1, \dots, j_N) |(i_1, j_1)\rangle \dots |(i_N, j_N)\rangle$



Resonating-Valence-Bond states

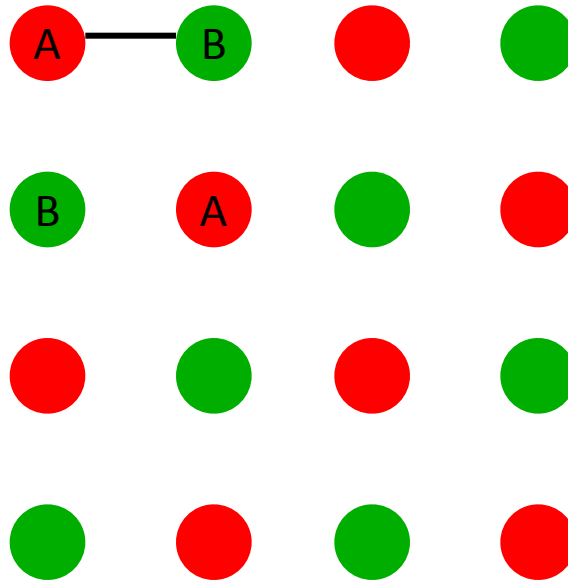
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- $|(i, j)\rangle = (|up\rangle |down\rangle - |down\rangle |up\rangle)/2$

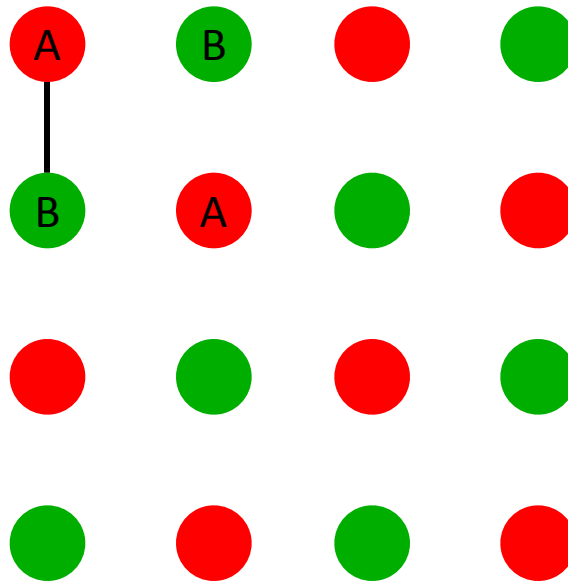
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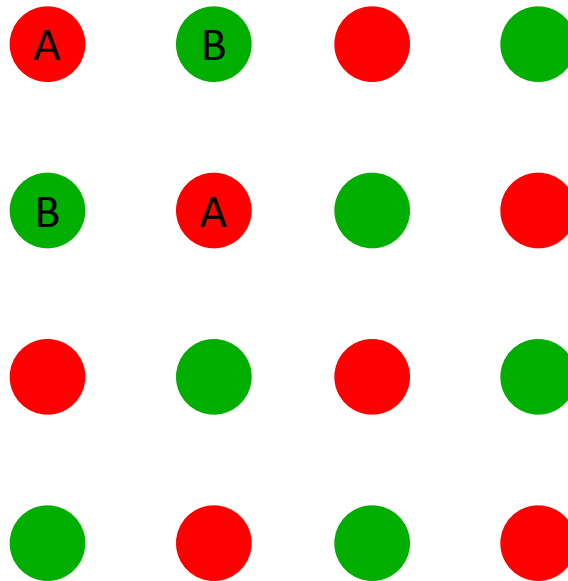
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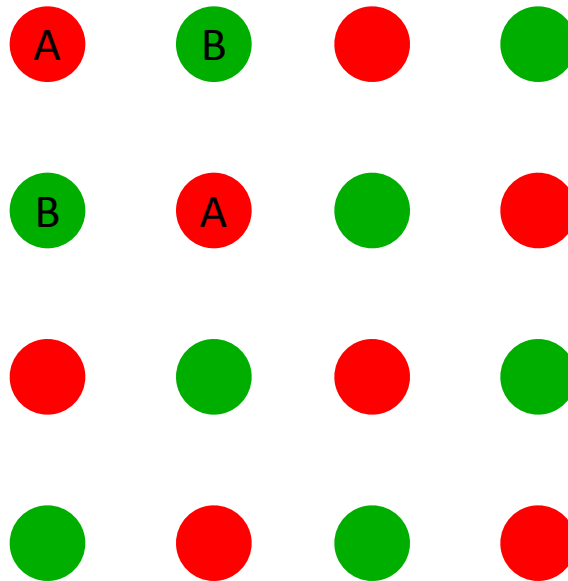
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- h is assumed to be isotropic over the lattice.

Resonating-Valence-Bond states

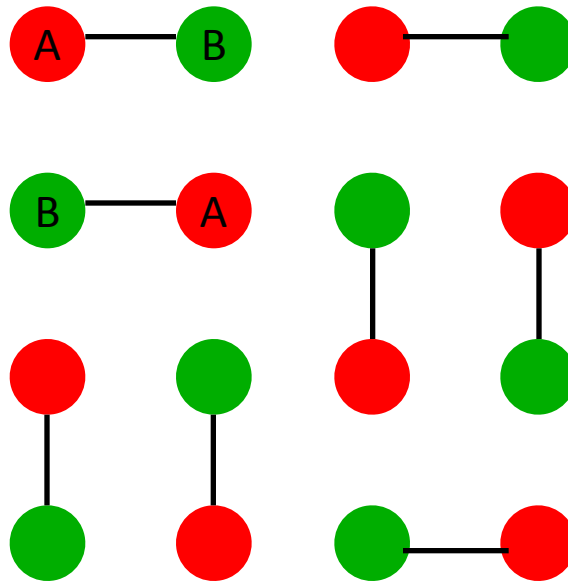
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- Coverings of only nearest neighbor dimers.

Resonating-Valence-Bond states

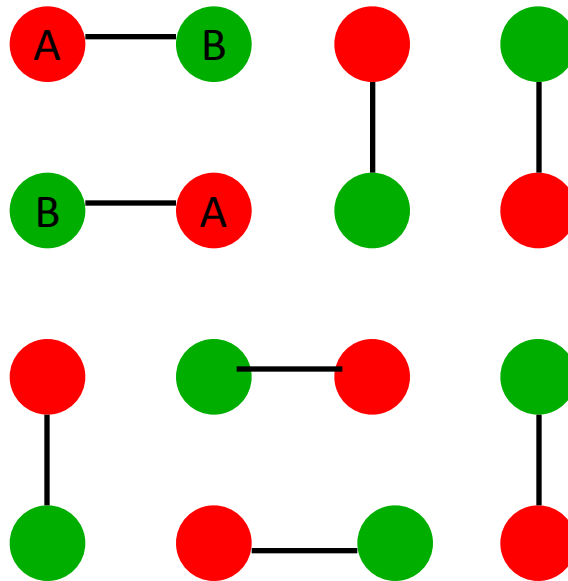
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Resonating-Valence-Bond states

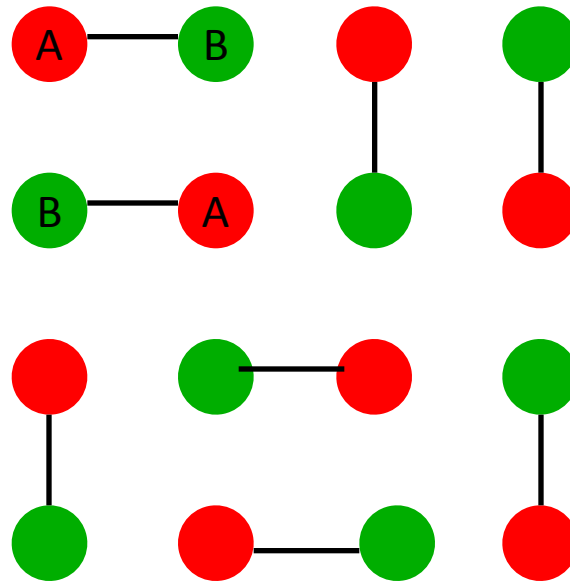
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Resonating-Valence-Bond states

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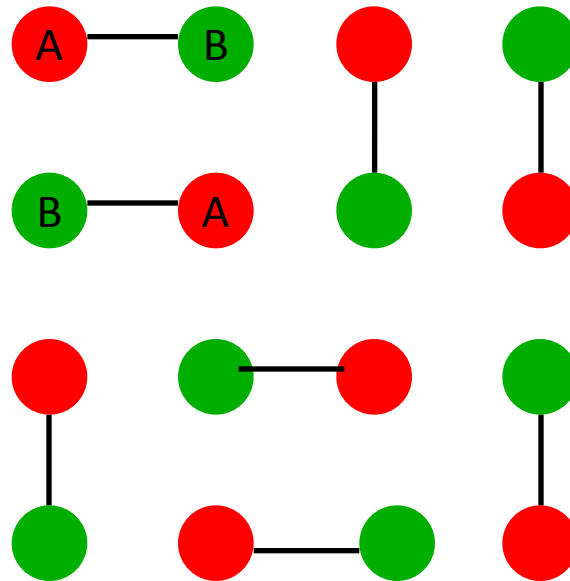


+ other nn coverings

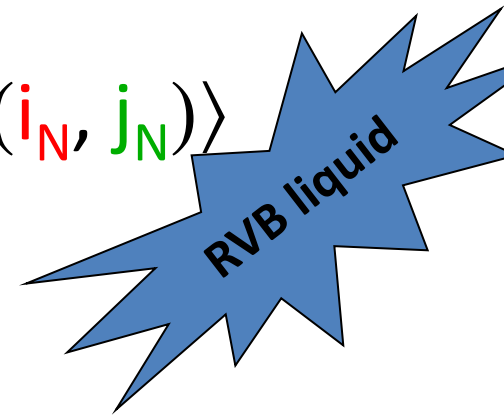
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Resonating-Valence-Bond states

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+ other nn coverings



- Coverings of only nearest neighbor dimers.

Genuine multiparty entanglement

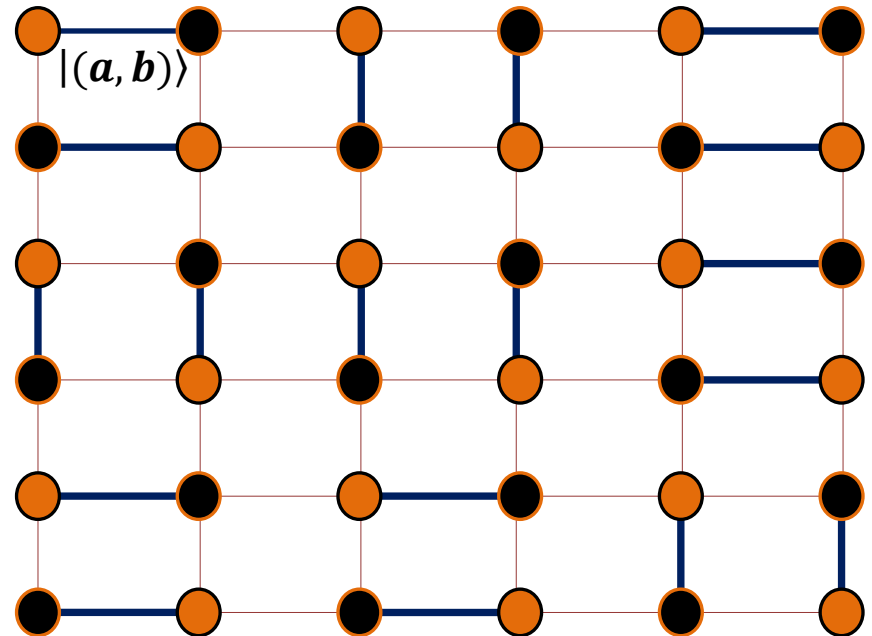
A multiparty pure quantum state is said to be **genuinely multiparty entangled** if

it is entangled across **every** bipartition.

ENTANGLEMENT PROPERTIES:

MULTIPARTITE ENTANGLEMENT IN ISOTROPIC RVB

CONSIDER AN **ODD: REST**
BIPARTITE CUT



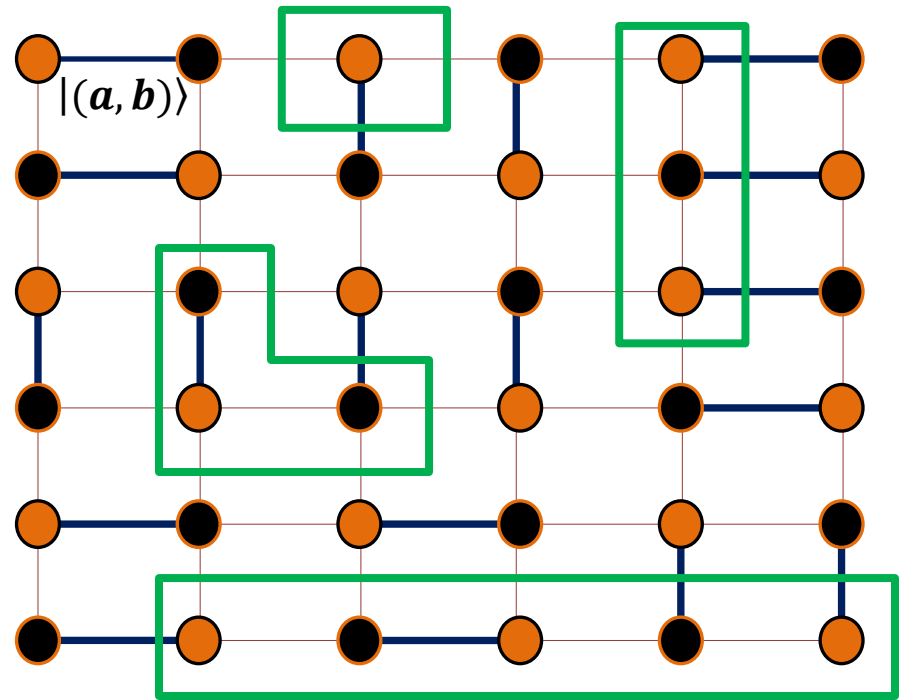
Chandran, Kaszlikowski, Sen(De), US, Vedral, PRL'07
Dhar, Sen(De), US, PRL'13

ENTANGLEMENT PROPERTIES:

MULTIPARTITE ENTANGLEMENT IN ISOTROPIC RVB

CONSIDER AN **ODD: REST**
BIPARTITE CUT

$$\rho_{odd} = \text{Tr}_{rest}(|\psi\rangle\langle\psi|_{RVB})$$



Chandran, Kaszlikowski, Sen(De), US, Vedral, PRL'07

Dhar, Sen(De), US, PRL'13

ENTANGLEMENT PROPERTIES:

MULTIPARTITE ENTANGLEMENT IN ISOTROPIC RVB

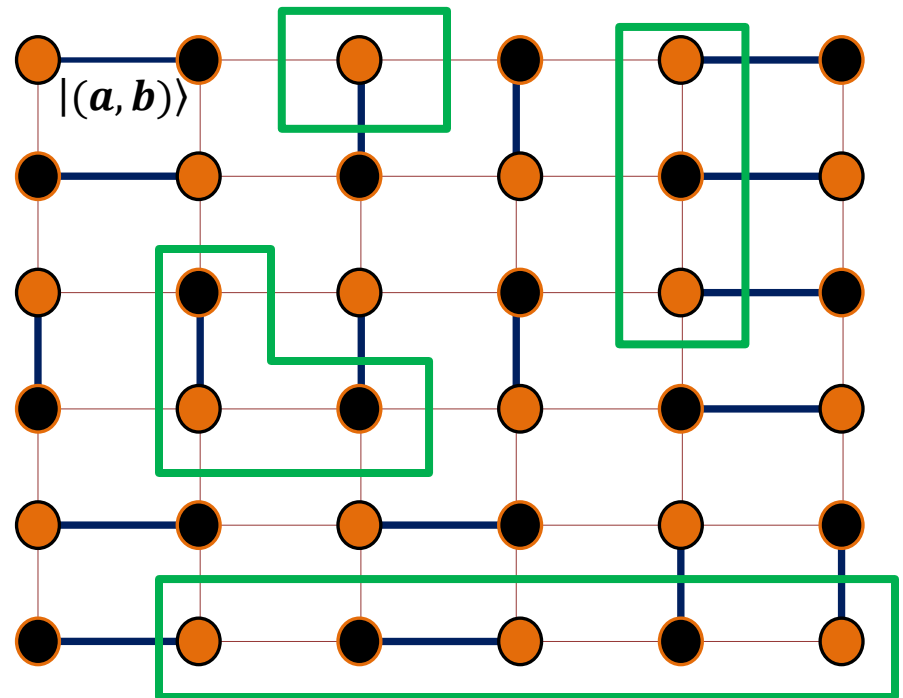
CONSIDER AN **ODD: REST**
BIPARTITE CUT

$$\rho_{odd} = \text{Tr}_{rest}(|\psi\rangle\langle\psi|_{RVB})$$

There is no pure ρ_{odd} that is rotationally invariant.

(Since RVB is rotationally invariant, ρ_{odd} has to be so.)

ρ_{odd} is mixed. The bipartition ODD: rest is entangled.

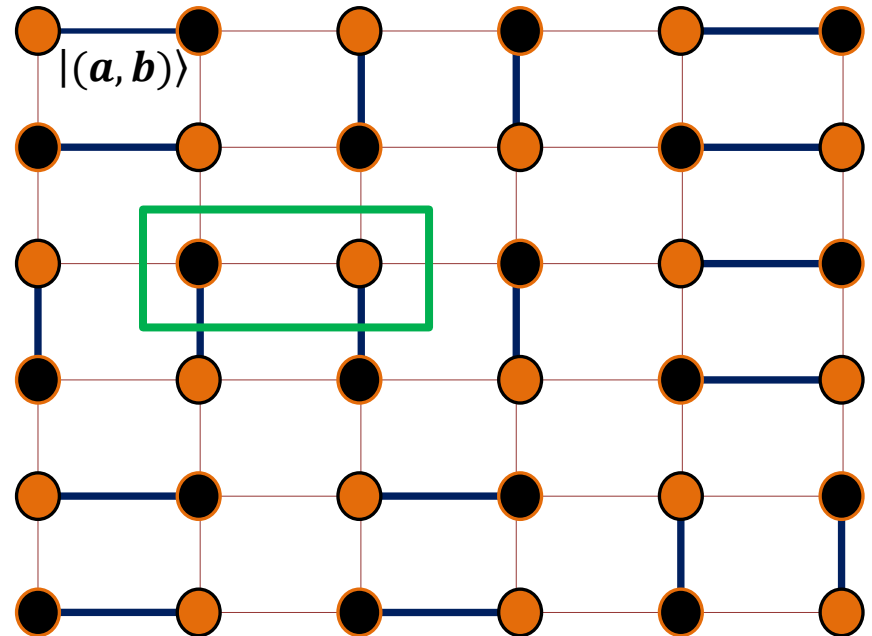


Chandran, Kaszlikowski, Sen(De), US, Vedral, PRL'07
Dhar, Sen(De), US, PRL'13

ENTANGLEMENT PROPERTIES:

MULTIPARTITE ENTANGLEMENT IN ISOTROPIC RVB

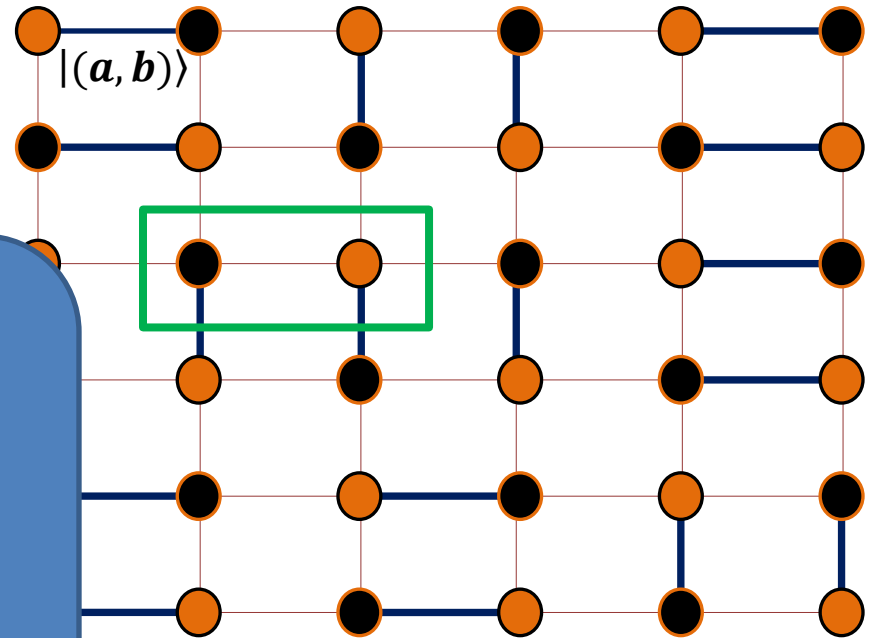
CONSIDER A **EVEN: REST**
BIPARTITE CUT



ENTANGLEMENT PROPERTIES:

MULTIPARTITE ENTANGLEMENT IN ISOTROPIC RVB

CONSIDER A **EVEN: REST**
BIPARTITE CUT



Can be shown to be
entangled by using
“strong subadditivity”

ENTANGLEMENT PROPERTIES:

MULTIPARTITE ENTANGLEMENT IN ISOTROPIC RVB

FOR INFINITE
INFINITE:IN
BIPARTITION

ASSUME EACH
(GREEN) IS

WE CAN ALWAYS
STRIPS THAT

WE CAN AGAIN
SINGLE SITE THAT IS
REQUIRED TO BE PURE

Similar proofs can show that
**INFINITE : INFINITE
BIPARTITIONS ARE ALSO
ENTANGLED**

In conclusion, ...

- Strong subadditivity is very useful.
- Has a host of important applications in QIC.

Thank you!



Planck & Co
Quantum
SMBOS (2013)
Physics D
12.0.188 (198)
 $E = \frac{h\nu}{\lambda}$
 $E = \frac{hc}{\lambda}$
 $\lambda = \frac{hc}{E}$
...

QIC Group @ HRI, 2013

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References r incomplete!