Aspects of Uncertainty and Complementarity

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Uncertainty and Complementarity





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Uncertainty and Complementarity



- 2 Complementarity and Entanglement
- Complementarity and Uncertainty
- 4 Conclusions



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Bohr's Complementarity Principle



Niels Bohr in 1928

In describing the results of quantum mechanical experiments, certain physical concepts are complementary. If two concepts are complementary, an experiment that clearly illustrates one concept will obscure the other complementary one....

--"The Quantum Postulate and the Recent Development of Atomic Theory," Supplement to Nature, April 14, 1928, p.580







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Which slit did the electron pass through?

Getting the "Welcher-Weg" (which-way) information





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Which slit did the electron pass through?

Getting the "Welcher-Weg" (which-way) information



No Interference!



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Complementarity in the 2-slit Experiment

Complementarity = Wave-particle duality?



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Complementarity = Wave-particle duality?

 In the two-slit experiment: the "which-way" information vs existence of interference pattern.



Complementarity = Wave-particle duality?

- In the two-slit experiment: the "which-way" information vs existence of interference pattern.
- They can NEVER be observed at the same time, in the same experiment.



Uncertainty and Complementarity



Heisenberg: Measurement of particle position results in uncontrollable disturbance in its momentum, washing out the interference pattern.



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IPQI 2014, 28/2 12

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Uncertainty and Complementarity



Heisenberg: *Measurement of particle position results in uncontrollable disturbance in its momentum, washing out the interference pattern.* Way out?



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Uncertainty and Complementarity



Heisenberg: *Measurement of particle position results in uncontrollable disturbance in its momentum, washing out the interference pattern.* Way out?

Determine which-way without disturbing the particle?







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Uncertainty and Complementarity

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... Einstein thought he had found a counterexample to the uncertainty principle. "It was quite a shock for Bohr he did not see the solution at once. During the whole evening he was extremely unhappy, going from one to the other and trying to persuade them that it couldn't be true, that it would be the end of physics if Einstein were right; but he couldn't produce any refutation.





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ROSENFELD (1968) Fundamental Problems in Elementary Particle Physics

Proceedings of the Fourteenth Solvay Conference, Interscience, New York, p. 232.



Replace the static source slit



Figures after Bo

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IPQI 2014, 28/2 14 / 38

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Replace the static source slit



by a movable slit

Figures after Bo

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Replace the static source slit



by a movable slit Recoil of slit \implies which-way information without disturbing the particle



Figures after B





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• Particle going through upper/lower slit has momentum $\pm p_0$





- Particle going through upper/lower slit has momentum $\pm p_0$
- Momentum conservation \implies recoil $\mp p_0$ of slit



- Particle going through upper/lower slit has momentum ±p0
- Momentum conservation \implies recoil $\mp p_0$ of slit
- Momentum of slit → which-way information



$$\Delta p_x = 2p\sin(\theta/2)$$



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$$\Delta p_x = 2p\sin(\theta/2) \approx p\theta = \frac{h}{\lambda}\theta$$



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$$\Delta p_x = 2p\sin(\theta/2) \approx p\theta = \frac{h}{\lambda}\theta = \frac{h}{\lambda}\frac{d}{L}$$



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IPQI 2014, 28/2 16 / 38



$$\Delta p_x = 2p\sin(\theta/2) \approx p\theta = \frac{h}{\lambda}\theta = \frac{h}{\lambda}\frac{d}{L}$$

This is the limit on accuracy of measuring recoil momentum.



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 For particles passing through Slit A and those through slit B:

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• Min uncertainty in position of source slit: $\Delta x = \frac{\hbar}{2\Delta p_x} = \frac{\lambda L}{4\pi d}$.





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- This is the uncertainty in position of a fringe.





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- Fringe separation $\delta x = \frac{\lambda L}{d}$.



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- This is the uncertainty in position of a fringe.
- Fringe separation $\delta x = \frac{\lambda L}{d}$.
- Interference pattern is lost!

• Complementarity enforced by Uncertainty Principle?



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- Complementarity enforced by Uncertainty Principle?
- Getting which-way information will necessarily disturb the state of the particle.

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- Complementarity enforced by Uncertainty Principle?
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- Disturbance will be enough to wash out interference.



- Complementarity enforced by Uncertainty Principle?
- Getting which-way information will necessarily disturb the state of the particle.
- Disturbance will be enough to wash out interference.
- This viewed as a restatement of Uncertainty Principle

Realization of Recoiling-Slit Experiment

PHYSICAL REVIEW A 75, 062105 (2007)

Trapped-ion realization of Einstein's recoiling-slit experiment

Robert S. Utter and James M. Feagin*

Department of Physics, California State University-Fullerton, Fullerton, California 92834, USA (Received 10 July 2006; revised manuscript received 9 October 2006; published 13 June 2007)

We analyze photon scattering by a harmonically trapped ion using two-port interferometry of the scattered photon and coherent-state measurement of the ion's external recoil motion. We examine how the coherent-state measurement could be used to mimick both momentum and position ion measurements and thus a modern realization of Wootters and Zurek's pioneering analysis of Einstein's historic recoiling-slit gedanken experi-





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IPQI 2014, 28/2 18 / 38

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Realization of Recoiling-Slit Experiment

Letters to Nature > Abstract

Letters to Nature



Nature 411, 166-170 (10 May 2001) | doi:10.1038/35075517; Received 22 December 2000; Accepted 7 March 2001

A complementarity experiment with an interferometer at the quantum-classical boundary

P. Bertet, S. Osnaghi, A. Rauschenbeutel, G. Nogues, A. Auffeves, M. Brune, J. M. Raimond & S. Haroche Physics Nobel 2012

 Laboratoire Kastler Brossel, Département de Physique, Ecole Normale Supérieure, 24 rue Lhomond, F-75231, Paris Cedex 05, France





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Quantum measurement

According to von Neumann

A quantum measurement consists of two processes:



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$$\sum_{i=1}^{n} c_{i} |d_{i}\rangle |\psi_{i}\rangle \xrightarrow[Process 2]{} |d_{k}\rangle |\psi_{k}\rangle$$

"The Measurement Problem".



Using von Neumann's process 1

Two orthogonal states of the particle depending on the path: slit 1: $|\psi_1\rangle$ slit 2: $|\psi_2\rangle$ Two orthogonal momentum states of the recoiling slit: $|p_1\rangle$ and $|p_2\rangle$.



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(a) Final state of particle+slit: necessary entanglement :

$$|\Psi\rangle = |\psi_1\rangle|p_1\rangle + |\psi_2\rangle|p_2\rangle$$



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and is enough to rule out interference!



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Without which-way information

Amplitude for finding the particle at point *x* on the screen is

$$\Psi(x) = \psi_1(x) + \psi_2(x).$$



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Amplitude for finding the particle at point *x* on the screen is

$$\Psi(x)=\psi_1(x)+\psi_2(x).$$

Probability (intensity):

 $|\Psi(x)|^2 = |\psi_1(x)|^2 + |\psi_2(x)|^2 + \frac{\psi_1^*(x)\psi_2(x) + \psi_2^*(x)\psi_1(x)}{\psi_1(x)}.$



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interference

WITH which-way information

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Interference vanishes if which-way information is obtained. •



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Can this argument be made more quantitative?


Suppose our detector distinguishes the two paths inaccurately.



Radhika Vathsan (BITS Goa)

Uncertainty and Complementarity

IPQI 2014, 28/2 24 / 38

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Suppose our detector distinguishes the two paths inaccurately.

This means "which-way" states $\langle d_1 | d_2 \rangle \neq 0$.



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This means "which-way" states $\langle d_1 | d_2 \rangle \neq 0$.

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Amplitude that the paths are perfectly distinguished



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Is there a relationship between them to capture complementarity?



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Uncertainty and Complementarity

IPQI 2014, 28/2 25 / 38

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So D^2 is the probability of correctly distinguishing the two paths.

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Uncertainty and Complementarity

IPQI 2014, 28/2 25 / 38

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t = 0: particle emerges from the double-slit with amplitude

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$$\Psi(x,0) = A\left(|d_1\rangle e^{-\frac{(x-d/2)^2}{4\epsilon^2}} + |d_2\rangle e^{-\frac{(x+d/2)^2}{4\epsilon^2}}\right),\,$$



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After time *t*, traveling a distance *L*, amplitude for particle to arrive at *x* on screen:

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where $A_t = \frac{1}{\sqrt{2}} [\sqrt{2\pi} (\epsilon + i\hbar t/2m\epsilon)]^{-1/2}$



Gaussian Wave-packet Model

Probability of finding particle at point *x* on the screen



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Probability of finding particle at point *x* on the screen

$$\begin{aligned} |\Psi(x,t)|^2 &= 2|A_t|^2 e^{-\frac{x^2+d^2/4}{2\sigma_t^2}}\cosh(xd/2\sigma_t^2) \\ &\times \left(1+|\langle d_1|d_2\rangle|\frac{\cos\left(\frac{xd\lambda L/2\pi}{4\epsilon^4+(\lambda L/2\pi)^2}+\theta\right)}{\cosh(xd/2\sigma_t^2)}\right) \end{aligned}$$



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$$\mathcal{V} \equiv \frac{I_{max} - I_{min}}{I_{max} + I_{min}}$$



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$$\mathcal{V}^2 + \mathcal{D}^2 \le 1.$$

Englert-Greenberger-Yasin duality relation



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Englert-Greenberger-Yasin duality relation A *quantitative* statement of complementarity



Origin of Complementarity?

Quantum correlations?



D.M. Greenberger, A. Yasin, Phys. Lett. A 128, 391 (1988), "Simultaneous wave and particle knowledge in a neutron interferometer",



B-G. Englert, *Phys. Rev. Lett.* **77**, 2154 (1996), "Fringe visibility and which-way information: an inequality"



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H. Wiseman, Phys. Lett. A 311, 285 (2003). "Directly observing momentum transfer in twin-slit which-way experiments"

Does the particle really receive a "momentum kick"?



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Uncertainty and duality

• "Which-way" states of the recoiling slit: $|d_1
angle$ and $|d_2
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(normalized, not necessarily orthogonal)



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Eigenstates of some observable \hat{P} with eigenvalues ± 1 .



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= 1 - \Delta P^2

Uncertainty and Duality

Correlation of detector states with particle states:

 $\Psi(x) = \psi_1(x)|p_1\rangle + \psi_2(x)|p_2\rangle.$



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Uncertainty and Complementarity

IPQI 2014, 28/2 32 / 38
Correlation of detector states with particle states:

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• Consider a basis change:

$$|p_1\rangle + |p_2\rangle \rightarrow \psi_1(x) + \psi_2(x)$$

 $|p_1\rangle - |p_2\rangle \rightarrow \psi_1(x) - \psi_2(x)$



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angle - |p_2
angle
ightarrow \psi_1(x) - \psi_2(x)$$

• \implies \exists another observable \hat{Q} with eigenvalues ± 1 and corresponding eigenstates

$$\begin{aligned} |q_1\rangle &= (|p_1\rangle + |p_2\rangle)/\sqrt{2} \\ |q_2\rangle &= (|p_1\rangle - |p_2\rangle)/\sqrt{2} \end{aligned}$$

Correlation of detector states with particle states:

$$\Psi(x) = \psi_1(x)|p_1\rangle + \psi_2(x)|p_2\rangle.$$

• Consider a basis change:

$$|p_1
angle + |p_2
angle
ightarrow \psi_1(x) + \psi_2(x) |p_1
angle - |p_2
angle
ightarrow \psi_1(x) - \psi_2(x)$$

• \implies \exists another observable \hat{Q} with eigenvalues ± 1 and corresponding eigenstates

$$\begin{array}{ll} |q_1\rangle &=& (|p_1\rangle + |p_2\rangle)/\sqrt{2} \\ |q_2\rangle &=& (|p_1\rangle - |p_2\rangle)/\sqrt{2} \end{array}$$

• The particle states can be correlated with these states:

$$\Psi(x) = \frac{c_1}{\sqrt{2}} [\psi_1(x) + \psi_2(x)] |q_1\rangle + \frac{c_2}{\sqrt{2}} [\psi_1(x) - \psi_2(x)] |q_2\rangle$$



Correlate the detected particles on the screen with the measured eigenstate of \hat{Q} ($c_1 = c_2$ case)



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Correlate the detected particles on the screen with the measured eigenstate of \hat{Q} ($c_1 = c_2$ case)



Two complementary interference patterns corresponding to $|q_1\rangle$ and $|q_2\rangle$.



Correlate the detected particles on the screen with the measured eigenstate of \hat{Q} ($c_1 = c_2$ case)



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Correlate the detected particles on the screen with the measured eigenstate of \hat{Q} ($c_1 = c_2$ case)



Two complementary interference patterns corresponding to $|q_1\rangle$ and $|q_2\rangle$.



For any c_1 , c_2 ,

$$|\Psi(x)|^{2} = \frac{|\psi_{1}(x)|^{2} + |\psi_{2}(x)|^{2}}{2} + \frac{|c_{1}|^{2} - |c_{2}|^{2}}{2} \left[\psi_{1}^{*}(x)\psi_{2}(x) + \psi_{2}^{*}(x)\psi_{1}(x)\right].$$



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For any c_1, c_2 ,

$$|\Psi(x)|^2 = rac{|\psi_1(x)|^2 + |\psi_2(x)|^2}{2} + rac{|c_1|^2 - |c_2|^2}{2} \left[\psi_1^*(x)\psi_2(x) + \psi_2^*(x)\psi_1(x)
ight].$$

Fringe visibility: $V^2 \le (|c_1|^2 - |c_2|^2)^2$.



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Fringe visibility: $V^2 \le (|c_1|^2 - |c_2|^2)^2$.

The uncertainty in \hat{Q} , in this *entangled* state:

$$\Delta Q^2 = 1 - (|c_1|^2 - |c_2|^2)^2.$$



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$$\Delta Q^2 = 1 - (|c_1|^2 - |c_2|^2)^2.$$

Thus

$$\mathcal{V}^2 \leq 1 - \Delta Q^2.$$



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Fringe visibility: $V^2 \le (|c_1|^2 - |c_2|^2)^2$.

The uncertainty in \hat{Q} , in this *entangled* state:

$$\Delta Q^2 = 1 - (|c_1|^2 - |c_2|^2)^2.$$

Thus

$$\mathcal{V}^2 \le 1 - \Delta Q^2.$$

Combining with the earlier result $D^2 = 1 - \Delta P^2$, we get

$$\mathcal{D}^2 + \mathcal{V}^2 \le 2 - [\Delta P^2 + \Delta Q^2].$$

The Sum Uncertainty Relation

Sum uncertainty relation for angular momenta ¹

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$$\Delta L_x^2 + \Delta L_y^2 + \Delta L_z^2 \ge \ell$$



¹Hoffmann, Takeuchi, *Phys. Rev. A* 68, 032103 (2003).

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$$\Delta L_x^2 + \Delta L_y^2 + \Delta L_z^2 \ge \ell$$

Implication for Pauli spin matrices

$$\Delta \sigma_x^2 + \Delta \sigma_y^2 + \Delta \sigma_z^2 \ge 2, \qquad \qquad \Delta \sigma_x^2 + \Delta \sigma_y^2 \ge 1.$$



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$$\Delta L_x^2 + \Delta L_y^2 + \Delta L_z^2 \ge \ell$$

Implication for Pauli spin matrices

$$\Delta \sigma_x^2 + \Delta \sigma_y^2 + \Delta \sigma_z^2 \ge 2, \qquad \Delta \sigma_x^2 + \Delta \sigma_y^2 \ge 1.$$

In our case, $\hat{P} = \hat{\sigma}_z$, $\hat{Q} = \hat{\sigma}_x$. So, $\Delta P^2 + \Delta Q^2 \ge 1$.

¹Hoffmann, Takeuchi, *Phys. Rev. A* 68, 032103 (2003).

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$$\Delta L_x^2 + \Delta L_y^2 + \Delta L_z^2 \ge \ell$$

Implication for Pauli spin matrices

$$\begin{split} & \Delta \sigma_x^2 + \Delta \sigma_y^2 + \Delta \sigma_z^2 \geq 2, \qquad \Delta \sigma_x^2 + \Delta \sigma_y^2 \geq 1. \\ & \text{In our case, } \hat{\pmb{P}} = \hat{\pmb{\sigma}}_z, \ \hat{\pmb{Q}} = \hat{\pmb{\sigma}}_x. \text{ So, } \boxed{\Delta P^2 + \Delta Q^2 \geq 1}. \\ & \text{Using this on} \\ & \mathcal{D}^2 + \mathcal{V}^2 \leq 2 - [\Delta P^2 + \Delta Q^2]. \end{split}$$
we get

$$\mathcal{D}^2 + \mathcal{V}^2 \le 1.$$

¹Hoffmann, Takeuchi, *Phys. Rev. A* 68, 032103 (2003).

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$$\Delta L_x^2 + \Delta L_y^2 + \Delta L_z^2 \ge \ell$$

Implication for Pauli spin matrices

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In our case, $\hat{P} = \hat{\sigma}_z$, $\hat{Q} = \hat{\sigma}_x$. So, $\Delta P^2 + \Delta Q^2 \ge 1$. Using this on

$$\mathcal{D}^2 + \mathcal{V}^2 \le 2 - [\Delta P^2 + \Delta Q^2].$$

we get

$$\mathcal{D}^2 + \mathcal{V}^2 \le 1$$

 The duality relation also emerges from the sum uncertainty relation.

 ¹Hoffmann, Takeuchi, Phys. Rev. A 68, 032103 (2003).

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- For any two orthogonal states of the recoiling slit (say) $|\xi_1\rangle$ and $|\xi_2\rangle$, one can *always* find operators \hat{P} and \hat{Q} whose uncertainties enforce complementarity.
 - $\hat{\pmb{P}} = |\xi_1\rangle\langle\xi_1| |\xi_2\rangle\langle\xi_2|$ $\hat{\pmb{Q}} = |\xi_1\rangle\langle\xi_2| + |\xi_2\rangle\langle\xi_1|$



 $\hat{\pmb{P}} = |\xi_1\rangle\langle\xi_1| - |\xi_2\rangle\langle\xi_2|$ $\hat{\pmb{Q}} = |\xi_1\rangle\langle\xi_2| + |\xi_2\rangle\langle\xi_1|$

• Englert-Greenberger-Yasin duality relation emerges from correlations and also from the sum uncertainty relation.



 $\hat{\boldsymbol{P}} = |\xi_1\rangle\langle\xi_1| - |\xi_2\rangle\langle\xi_2|$ $\hat{\boldsymbol{Q}} = |\xi_1\rangle\langle\xi_2| + |\xi_2\rangle\langle\xi_1|$

- Englert-Greenberger-Yasin duality relation emerges from correlations and also from the sum uncertainty relation.
- Complementarity enforced by correlations and the uncertainty relations are two sides of a coin (provided the observables are correctly identified).



 $\hat{\boldsymbol{P}} = |\xi_1\rangle\langle\xi_1| - |\xi_2\rangle\langle\xi_2|$ $\hat{\boldsymbol{Q}} = |\xi_1\rangle\langle\xi_2| + |\xi_2\rangle\langle\xi_1|$

- Englert-Greenberger-Yasin duality relation emerges from correlations and also from the sum uncertainty relation.
- Complementarity enforced by correlations and the uncertainty relations are two sides of a coin (provided the observables are correctly identified).
- Momentum back-action of the recoiling slit on the particle plays no role in complementarity.





🐚 Tabish Qureshi, Radhika Vathsan Einstein's Recoiling Slit Experiment, Complementarity and Uncertainty Arxiv: 1210.4248 [quant-ph] Quanta Vol. 2 (April 2013)



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THANK YOU!

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Uncertainty and Complementarity

IPQI 2014, 28/2 38 / 38