



Multiparty entanglement Vs classical information transmission

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IPQI-2014, IOP, Bhubaneswar

OUTLINE

Multiparty entanglement

Dense coding - single to many

Connecting both of them

Analytically

Numerically

Dense coding - many to single

Noiseless

Noisy

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Dense coding - many to single

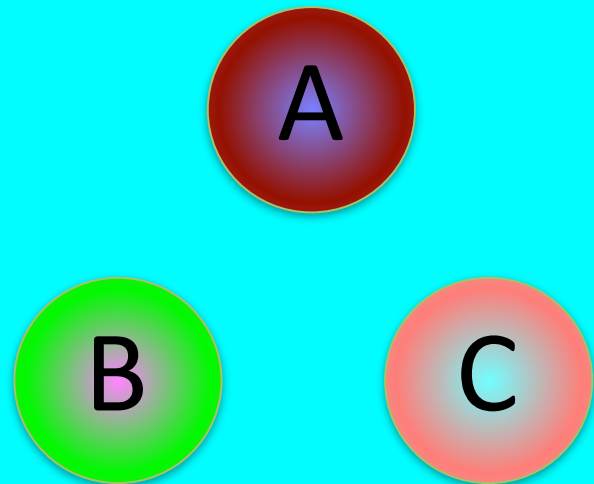
Noiseless

Noisy

Multiparty entanglement

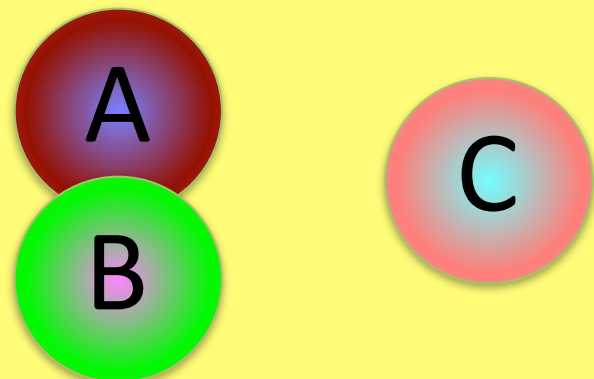
- Relative entropy of entanglement (REE)
 - Genuinely multiparty entangled
- Generalized geometric measure (GGM)
- Monogamy based measure

Types of entangled states



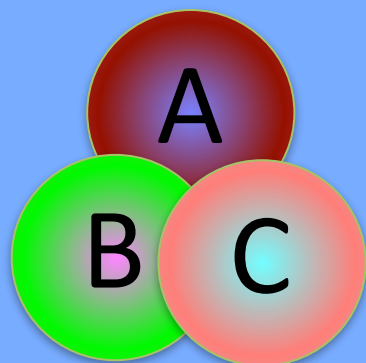
Fully separable:

$$|\phi\rangle \otimes |\xi\rangle \otimes |\chi\rangle$$



Biseparable:

$$|\psi\rangle_{AB} \otimes |\chi\rangle$$



Genuine multiparty entangled:

$$\neq |\phi\rangle \otimes |\xi\rangle \otimes |\chi\rangle$$

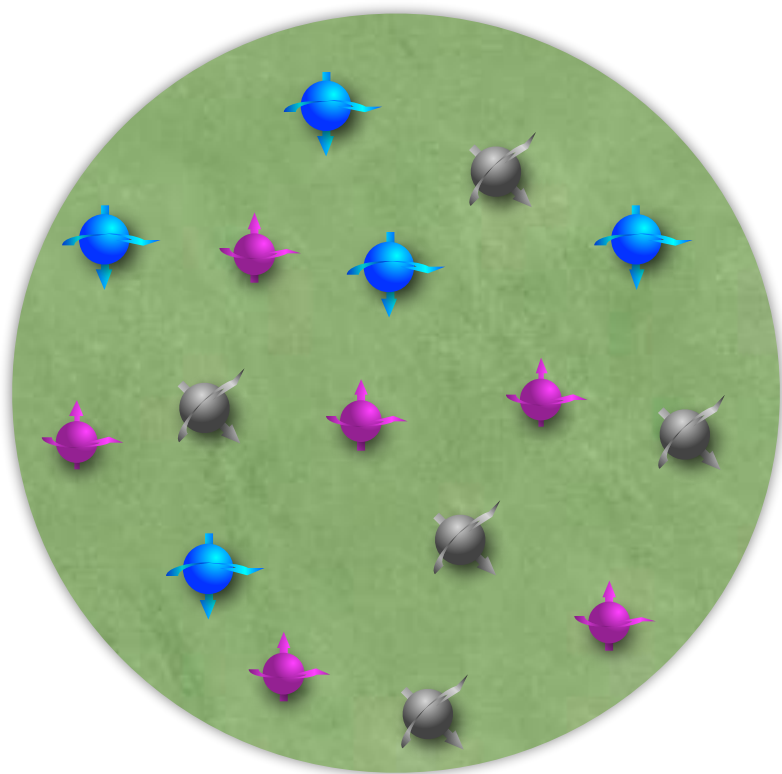
$$\neq |\psi\rangle_{AB} \otimes |\chi\rangle$$

Multiparty entanglement

- Relative entropy of entanglement (REE)
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Multiparty entanglement

Relative entropy of entanglement



$$Q_{A_1 A_2 \dots A_n}$$

↓
Pure or
mixed state

V. Vedral, M. B. Plenio, M. A. Rippin, and P. L. Knight, Phys. Rev. Lett. 78, 2275 (1997)

V. Vedral, M. B. Plenio, K. Jacobs, and P. L. Knight, Phys. Rev. A 56, 4452 (1997)

V. Vedral, Rev. Mod. Phys. 74, 197 (2002)

Multiparty entanglement

Relative entropy of entanglement

$$E_R(\rho_{A_1 A_2 \dots A_n}) = \min_{\sigma \in n\text{-gen}} S(\rho_{A_1 A_2 \dots A_n} || \sigma)$$

{not genuinely multiparty
entangled}

$$S(\rho || \sigma) = \text{tr}(\rho \log_2 \rho - \rho \log_2 \sigma)$$

Relative entropy

V. Vedral, M. B. Plenio, M. A. Rippin, and P. L. Knight, Phys. Rev. Lett. 78, 2275 (1997)

V. Vedral, M. B. Plenio, K. Jacobs, and P. L. Knight, Phys. Rev. A 56, 4452 (1997)

V. Vedral, Rev. Mod. Phys. 74, 197 (2002)

Multiparty entanglement

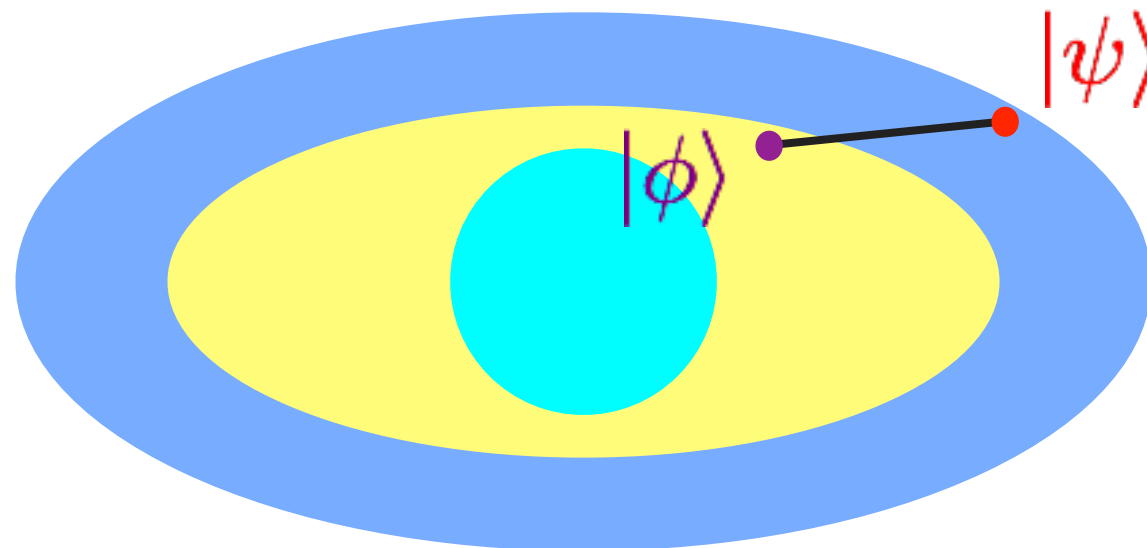
- Relative entropy of entanglement (REE)
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Multiparty entanglement

Generalized geometric measure (GGM)

$$\mathcal{E}(|\psi_{A_1 A_2 \dots A_n}\rangle) = 1 - \max_{|\phi\rangle} |\langle\phi|\psi_{A_1 A_2 \dots A_n}\rangle|^2$$

$|\phi\rangle$ is not genuinely entangled state



Multiparty entanglement

Generalized geometric measure (GGM)

$$\mathcal{E}(|\psi_{A_1 A_2 \dots A_n}\rangle) = 1 - \max\{\lambda_{A:\mathcal{B}}^2 \mid \mathcal{A} \cup \mathcal{B} = \{1, 2, \dots, N\}, \mathcal{A} \cap \mathcal{B} = \emptyset\}$$



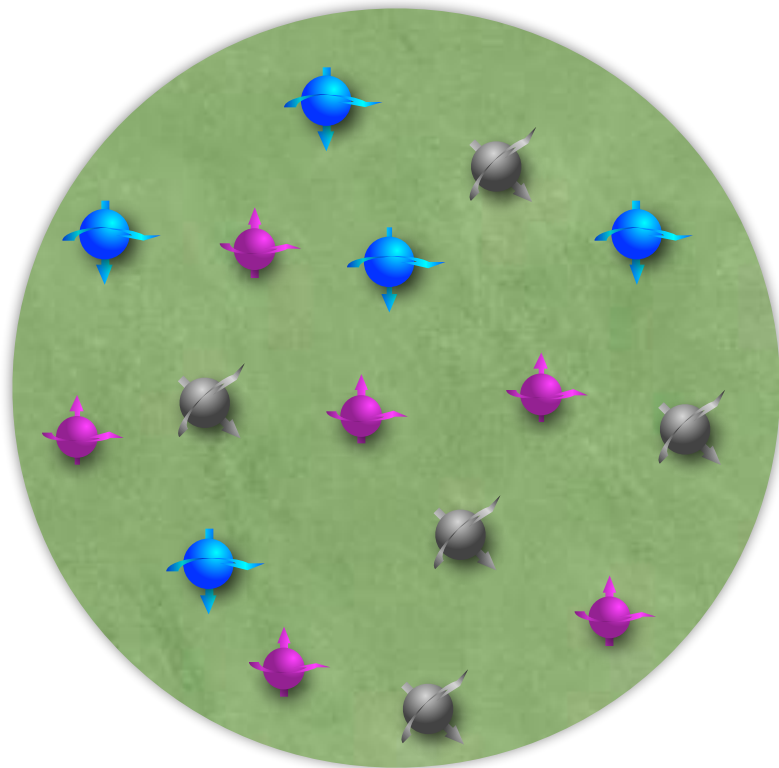
maximal Schmidt coefficients in the
 $A : \mathcal{B}$ bipartite split of $|\phi\rangle$

Multiparty entanglement

- Relative entropy of entanglement (REE)
- Generalized geometric measure (GGM)
- Monogamy based measure

Multiparty entanglement

Monogamy based measures



$$\sum_{i=1}^N Q_{A_i B} \leq Q_{A_1 A_2 \dots A_N : B}$$

Quantum correlation
measure

$$\delta_Q = Q_{A_1 A_2 \dots A_N : B} - \sum_{i=1}^N Q_{A_i B}$$

Positive \longrightarrow Monogamous

Negative \longrightarrow Non-Monogamous

- V. Coffman, J. Kundu, and W. K. Wootters, Phys. Rev. A 61, 052306 (2000)
T. J. Osborne and F. Verstraete, Phys. Rev. Lett. 96, 220503 (2006)
R. Prabhu, A.K. Pati, A. Sen(De), and U. Sen, Phys. Rev. A 86, 052337 (2012)

Multiparty entanglement

Monogamy based measures

$$\delta_Q = Q_{A_1 A_2 \dots A_N : B} - \sum_{i=1}^N Q_{A_i B}$$

$Q = C^2 =$ Concurrence squared

V. Coffman, J. Kundu, and W. K. Wootters, Phys. Rev. A 61, 052306 (2000)

T. J. Osborne and F. Verstraete, Phys. Rev. Lett. 96, 220503 (2006)

$Q = D =$ Quantum discord

R. Prabhu, A.K. Pati, A. Sen(De), and U. Sen, Phys. Rev. A 85, 040102(R) (2012); 86, 052337 (2012)

G. L. Giorgi, Phys. Rev. A 84, 054301 (2011)

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Multiparty entanglement

Dense coding - single to many

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Classical information transfer

Quantum Channel



$$|\psi_{AB}\rangle = \frac{1}{\sqrt{2}}(|0_A 0_B\rangle + |1_A 1_B\rangle)$$



2 bits
requires
2 dimension state

Dense coding capacity

$$\mathcal{C}(\rho_{AB}) = \log_2 d_A + S(\rho_B) - S(\rho_{AB})$$

dimension of
subsystem A

$S(\rho) = -\text{tr}(\rho \log \rho)$
von Neumann entropy

Dense coding capacity

$$\mathcal{C}(\rho_{AB}) = \log_2 d_A + S(\rho_B) - S(\rho_{AB})$$

Classical dense
coding capacity

Dense coding capacity

$$\mathcal{C}(\rho_{AB}) = \log_2 d_A + S(\rho_B) - S(\rho_{AB})$$

Quantum
DC advantage

Quantum advantage

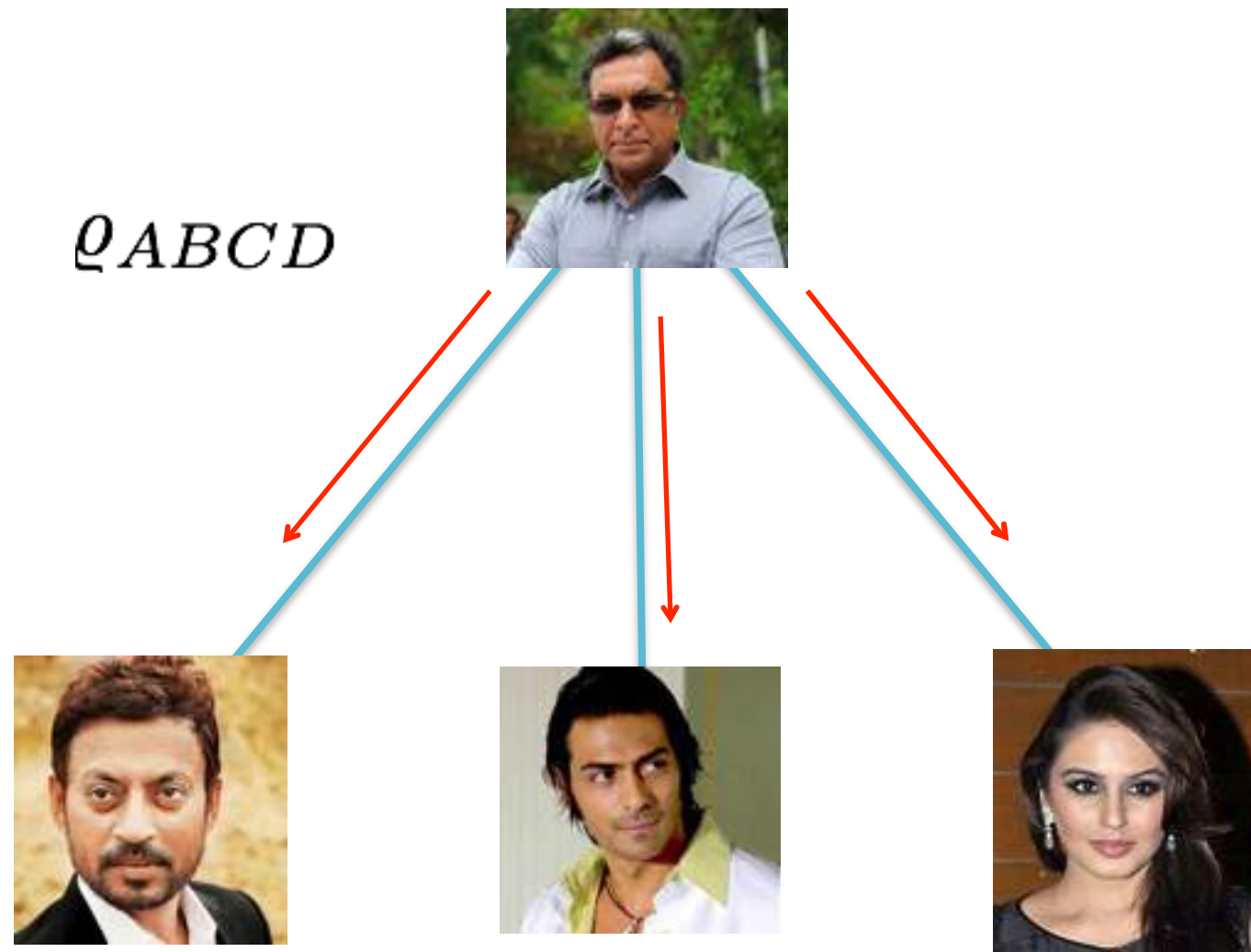
Quantum advantage in dense coding of AB channel:

$$C_{adv}(\rho_{AB}) = \max\{S(\rho_B) - S(\rho_{AB}), 0\}$$

Therefore

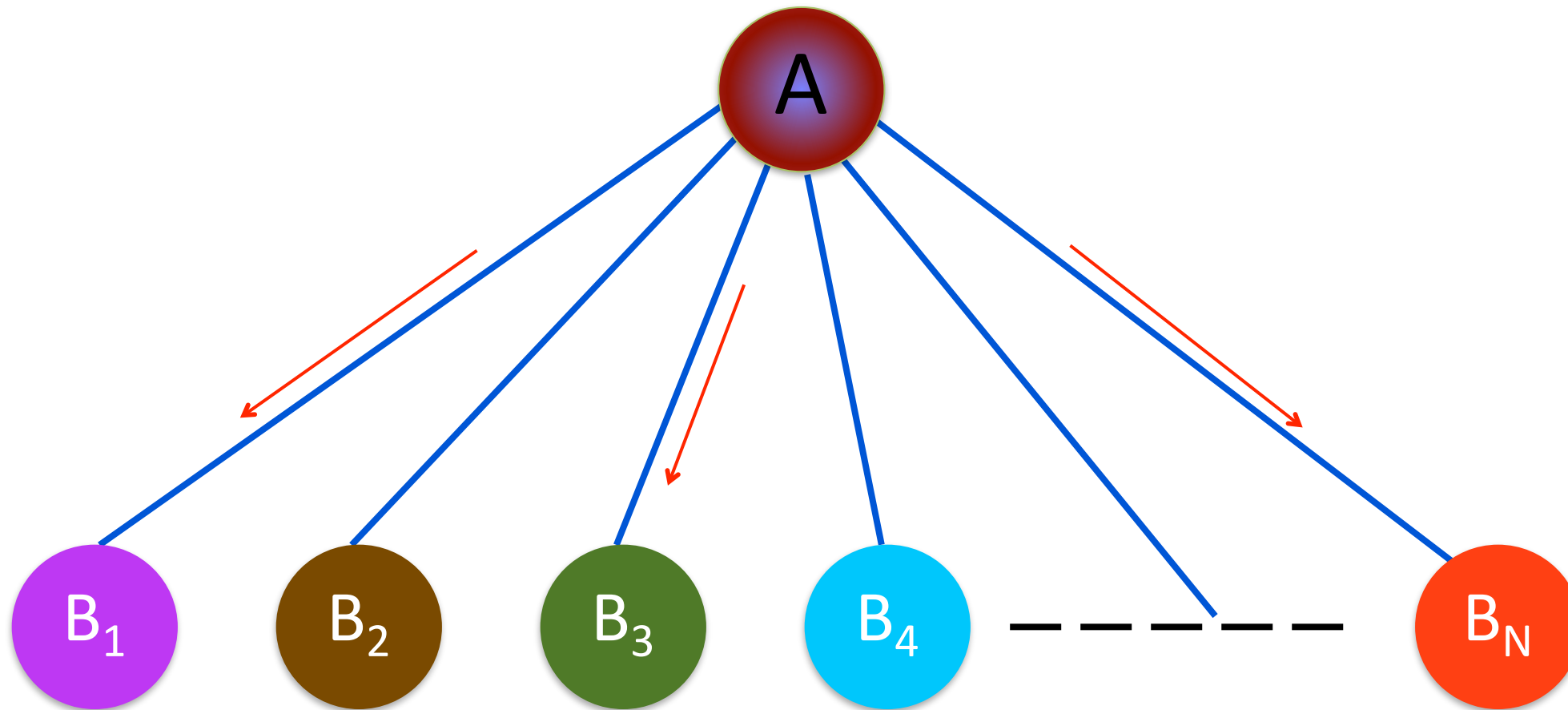
$$C(\rho_{AB}) = \log_2 d_A + C_{adv}(\rho_{AB})$$

Quantum advantage in networks



$$C_{adv}^{max}(\rho_{ABCD}) = \max\{C_{adv}(\rho_{AB}), C_{adv}(\rho_{AC}), C_{adv}(\rho_{AD})\}$$

Quantum advantage in networks



$$C_{\text{adv}}^{\text{max}}(Q_{AB_1B_2\dots B_N}) = \max\{C_{\text{adv}}(Q_{AB_i}) | i = 1, 2, \dots, N\}$$

OUTLINE

Multiparty entanglement

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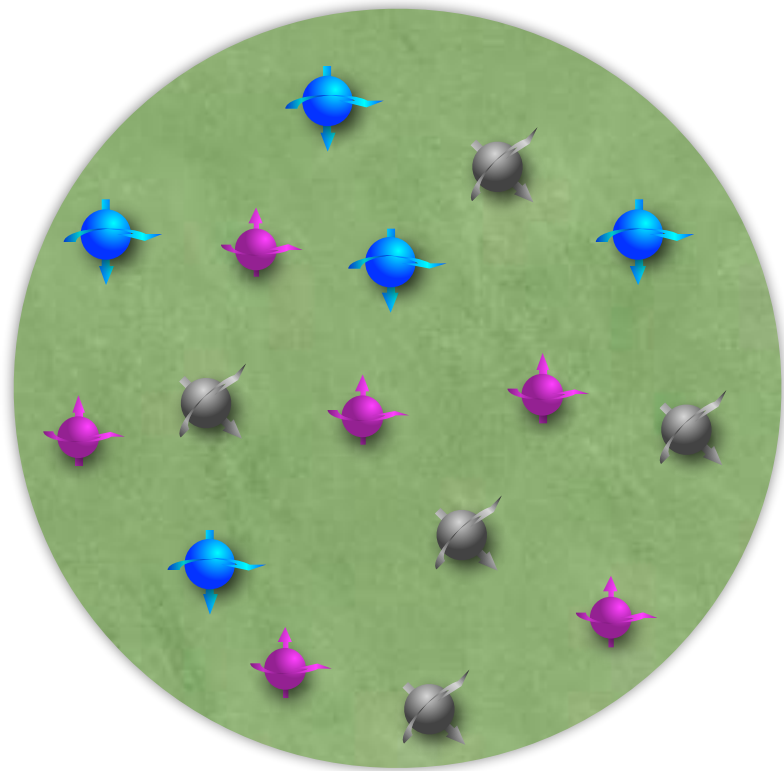
Numerically

Dense coding - many to single

Noiseless

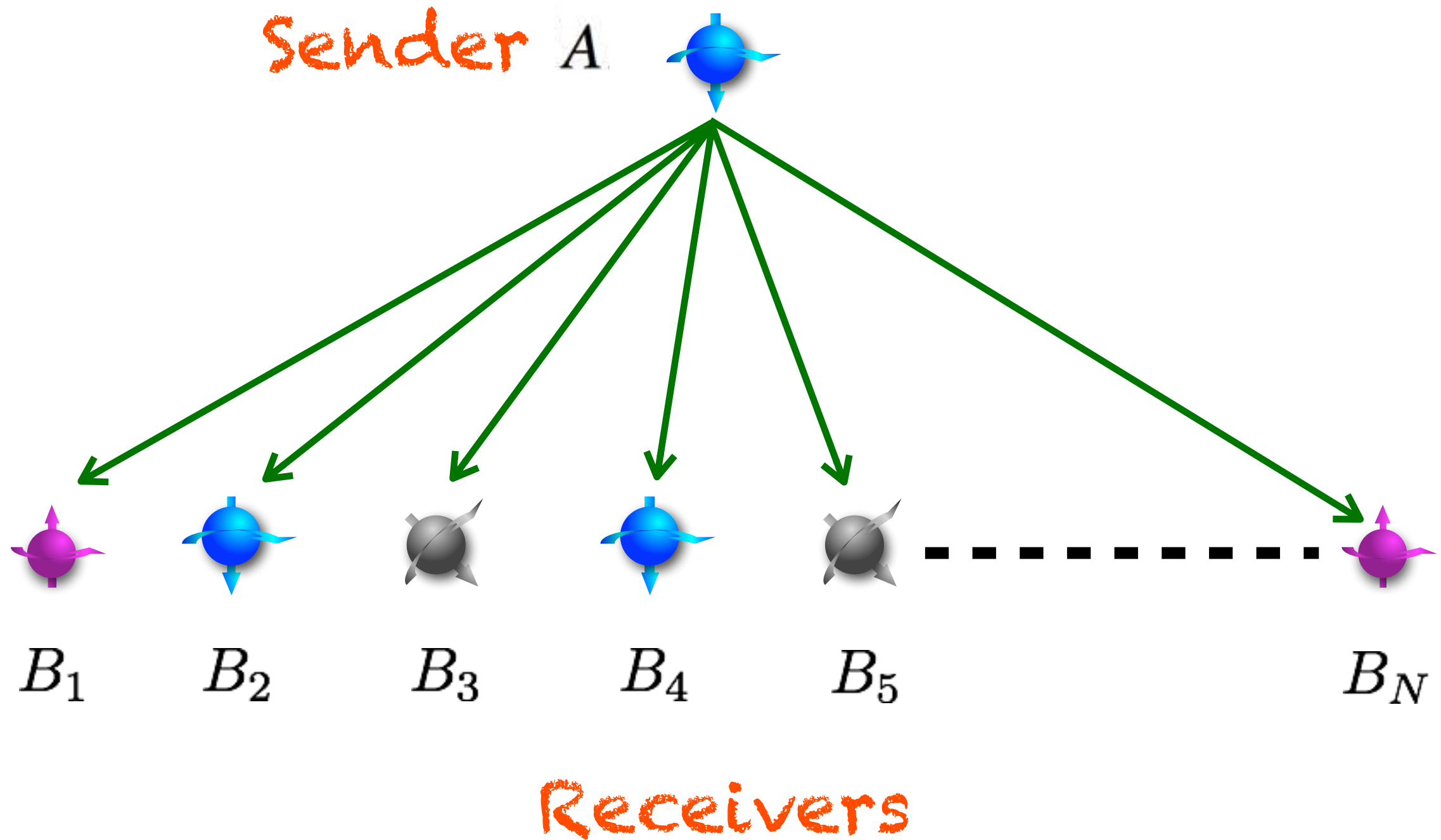
Noisy

Multiparty DC & Entanglement



$$Q_{AB_1 B_2 \dots B_N}$$

Multiparty DC & Entanglement



Multiparty DC & Entanglement

$$C_{\text{adv}}^{\max}(\rho_{AB_1B_2\dots B_N})$$

DC advantage

REE

$$E_R(\rho_{A_1A_2\dots A_n}) = \min_{\sigma \in n\text{-gen}} S(\rho_{A_1A_2\dots A_n} || \sigma)$$

GGM

$$\mathcal{E}(|\psi_{A_1A_2\dots A_n}\rangle) = 1 - \max_{|\phi\rangle} |\langle \phi | \psi_{A_1A_2\dots A_n} \rangle|^2$$

Monogamy score

$$\delta_{\mathcal{Q}} = \mathcal{Q}_{A_1A_2\dots A_N:B} - \sum_{i=1}^N \mathcal{Q}_{A_iB}$$

Multiparty DC & Entanglement

DC advantage & REE

$$\begin{aligned} E_R(\rho_{AB_1 B_2 \dots B_N}) &= \min_{\sigma \in (N+1)\text{-gen}} S(\rho || \sigma) \\ &\leq \min_{\sigma' \in \text{sep}_1} S(\rho || \sigma') \equiv E_R^{AB_1:\text{rest}}(\rho_{AB_1:B_2 \dots B_N}) \\ &\leq E_f^{AB_1:\text{rest}}(\rho_{AB_1:B_2 \dots B_N}) \\ &\leq S(\rho_{AB_1}) \end{aligned}$$

separable states in
 $AB_1 : B_2 \dots B_N$

Entanglement of
formation

C. H. Bennett, D. P. DiVincenzo, J. A. Smolin, and
W. K. Wootters, Phys. Rev. A 54, 3824 (1996)

$$E_R(\rho_{AB_1 B_2 \dots B_N}) \leq S(\rho_{AB_1}) \equiv S_{AB_1}$$

V. Vedral and M. B. Plenio, Phys. Rev. A 57, 1619 (1998)

Multiparty DC & Entanglement

DC advantage & REE

$$C_{\text{adv}}^{\max}(\rho_{AB_1B_2\dots B_N}) = S_{B_1} - S_{AB_1}$$

$$E_R(\rho_{AB_1B_2\dots B_N}) \leq S_{AB_1}$$

$$S_{B_1} \leq \log_2 d_{B_1}$$

$$C_{\text{adv}}^{\max} + E_R \leq \log_2 d$$

Multiparty DC & Entanglement

$$C_{\text{adv}}^{\max}(\rho_{AB_1B_2\dots B_N})$$

DC advantage

REE

$$E_R(\rho_{A_1A_2\dots A_n}) = \min_{\sigma \in n\text{-gen}} S(\rho_{A_1A_2\dots A_n} || \sigma)$$

GGM

$$\mathcal{E}(|\psi_{A_1A_2\dots A_n}\rangle) = 1 - \max_{|\phi\rangle} |\langle \phi | \psi_{A_1A_2\dots A_n} \rangle|^2$$

Monogamy score

$$\delta_{\mathcal{Q}} = \mathcal{Q}_{A_1A_2\dots A_N:B} - \sum_{i=1}^N \mathcal{Q}_{A_iB}$$

Multiparty DC & Entanglement

DC advantage & GGM

$$\frac{1}{\log_2 d} C_{\text{adv}}^{\text{max}} + \frac{d}{d-1} \mathcal{E} \leq 1$$

Normalization terms which makes individual quantities maximum

snippet

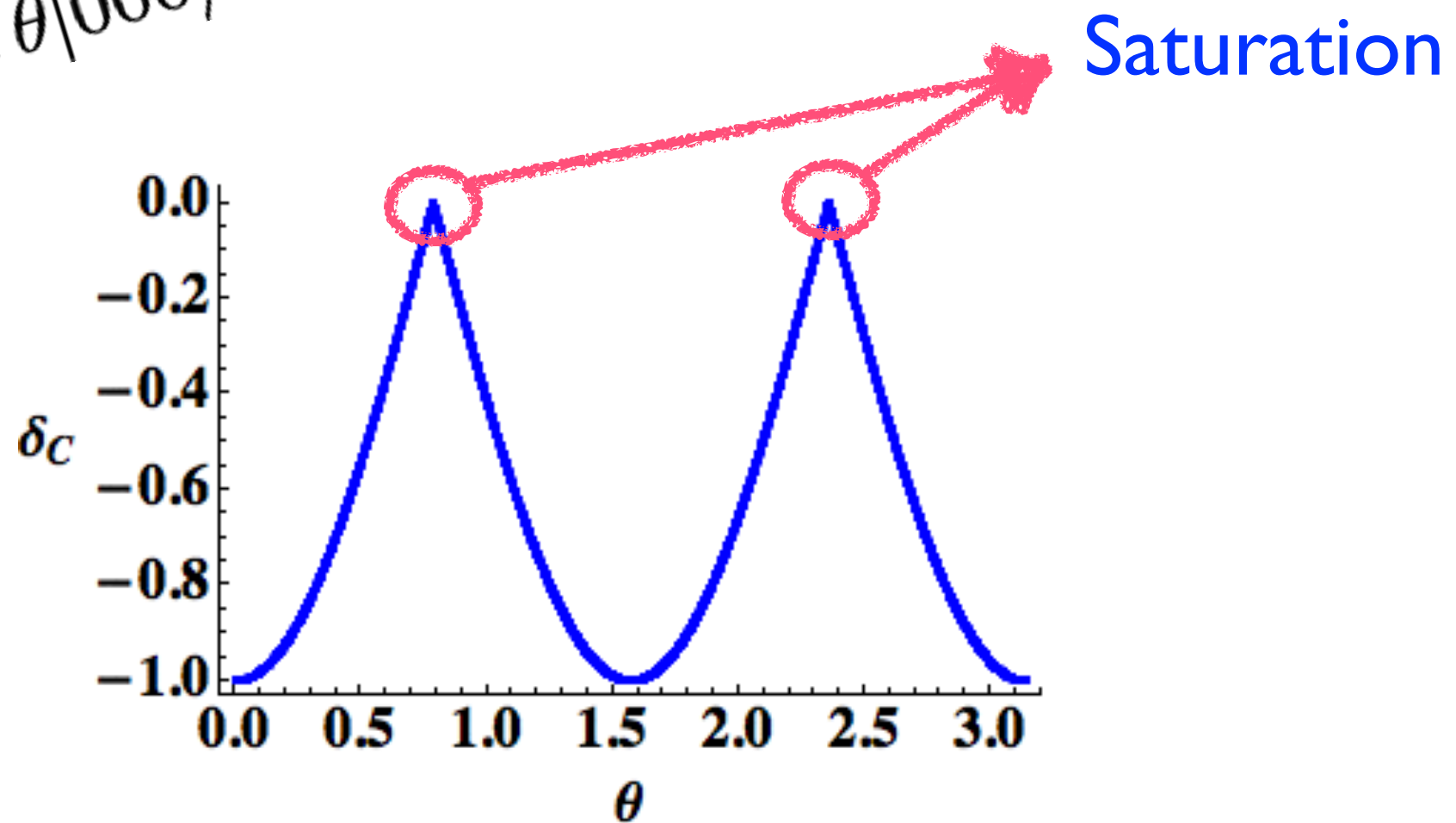
True for arbitrary states

R.P, A. Sen(De), and U. Sen, Phys. Rev. A 88, 042329 (2013)

Multiparty DC & Entanglement

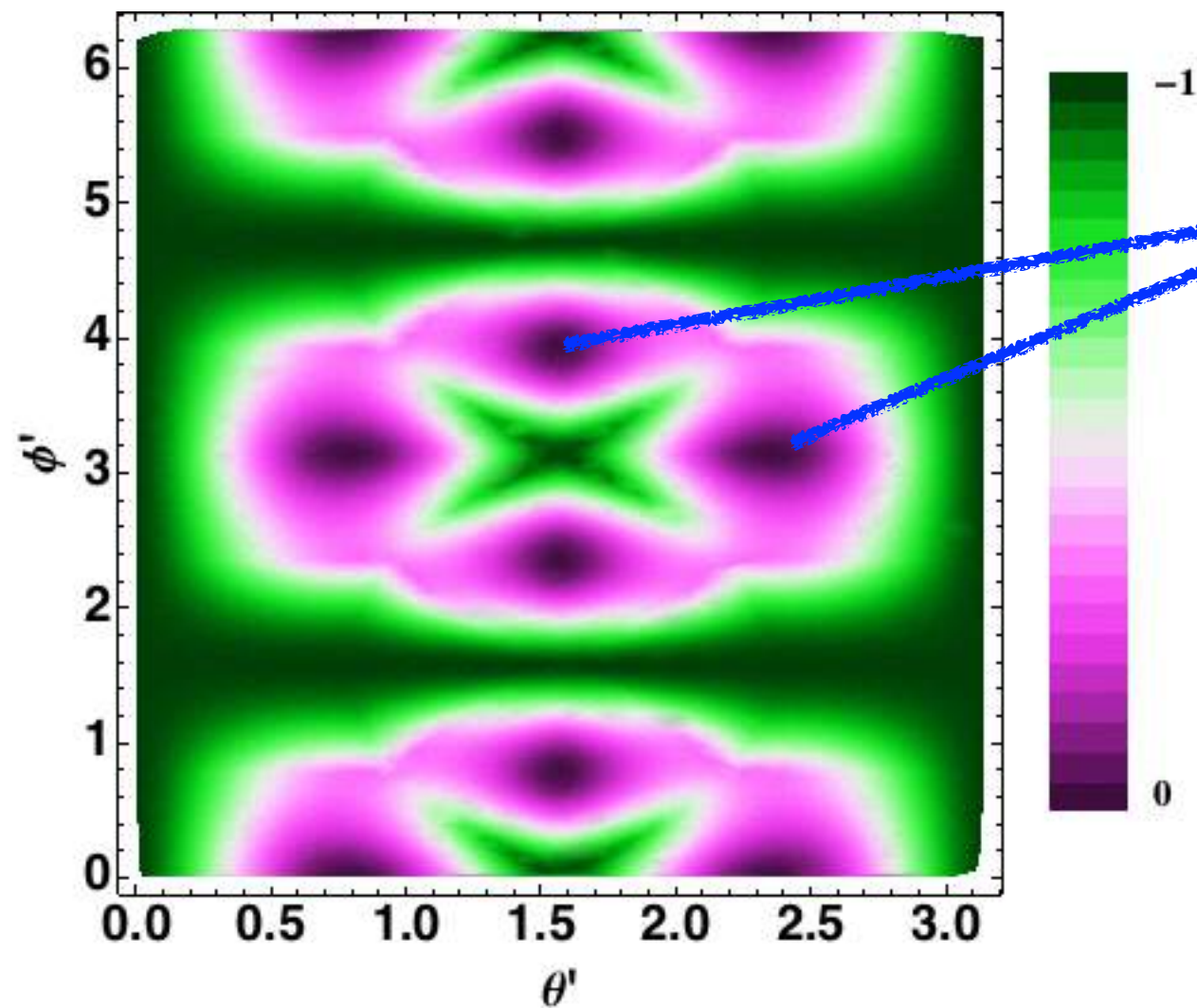
$$\delta_C = \frac{1}{\log_2 d} \mathcal{C}_{\text{adv}}^{\text{max}} + \frac{d}{d-1} \mathcal{E} - 1$$

$$|\psi_{\text{GHZ}}\rangle = \cos \theta |000\rangle + \sin \theta |111\rangle$$



Multiparty DC & Entanglement

$$|\psi_W\rangle = \sin \theta' \cos \phi' |011\rangle + \sin \theta' \sin \phi' |101\rangle + \cos \theta' |110\rangle$$



Saturation

Multiparty DC & Entanglement

$$C_{\text{adv}}^{\max}(\rho_{AB_1B_2\dots B_N})$$

DC advantage

REE

$$E_R(\rho_{A_1A_2\dots A_n}) = \min_{\sigma \in n\text{-gen}} S(\rho_{A_1A_2\dots A_n} || \sigma)$$

GGM

$$\mathcal{E}(|\psi_{A_1A_2\dots A_n}\rangle) = 1 - \max_{|\phi\rangle} |\langle \phi | \psi_{A_1A_2\dots A_n} \rangle|^2$$

Monogamy score

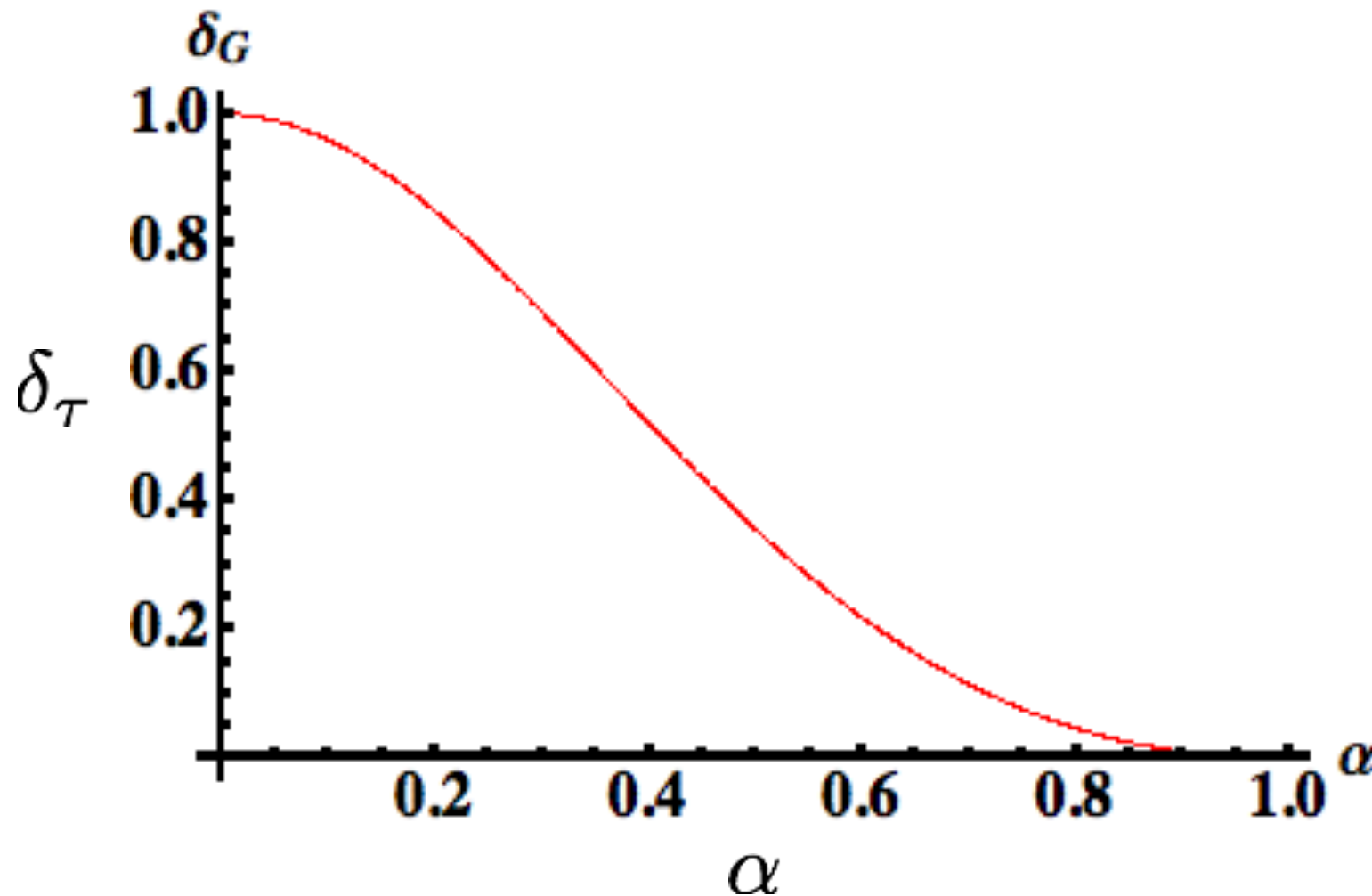
$$\delta_Q = Q_{A_1A_2\dots A_N:B} - \sum_{i=1}^N Q_{A_iB}$$

MDCC states

3-qubit

Maximally-dense-coding-capable states (MDCC)

$$|\psi_\alpha\rangle = |111\rangle + |000\rangle + \alpha(|101\rangle + |010\rangle)$$



MDCC states

3-qubit

Maximally-dense-coding-capable states (MDCC)

$$|\psi_\alpha\rangle = |111\rangle + |000\rangle + \alpha(|101\rangle + |010\rangle)$$



Belongs to GHZ class, except for $\alpha = 1$

Arbitrary state $|\psi\rangle$ and set $\mathcal{E}(|\psi_\alpha\rangle) = \mathcal{E}(|\psi\rangle)$

$$\mathcal{C}_{\text{adv}}(|\psi_\alpha\rangle) \geq \mathcal{C}_{\text{adv}}(|\psi\rangle)$$

MDCC states

3-qubit

Maximally-dense-coding-capable states (MDCC)

$$|\psi_\alpha\rangle = |111\rangle + |000\rangle + \alpha(|101\rangle + |010\rangle)$$

Eigenvalues of marginal density matrices

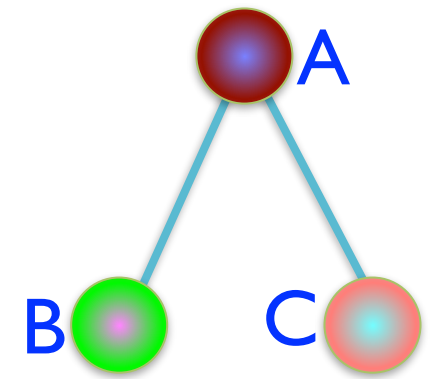
$$\lambda_A = \frac{1}{2} (\text{twice}) = \lambda_C \quad \lambda_B = \frac{(1 \pm \alpha)^2}{2(1 + \alpha^2)}$$

$$\mathcal{E}(|\psi_\alpha\rangle) = 1 - \frac{(1 + \alpha)^2}{2(1 + \alpha^2)} = \frac{1}{2} - \frac{\alpha}{1 + \alpha^2}$$

$$\mathcal{E}(|\psi\rangle) = 1 - \max\{\lambda_A, \lambda_B, \lambda_C\} = 1 - \lambda$$

$$\frac{\alpha}{1 + \alpha^2} = \lambda - \frac{1}{2}$$

MDCC states



$$C_{adv}(|\psi_\alpha\rangle) = S_C^\alpha - S_B^\alpha$$

$$\lambda_C = \frac{1}{2} \text{ (twice)}$$

$$\lambda_B = \frac{(1 \pm \alpha)^2}{2(1 + \alpha^2)}$$

$$\frac{\alpha}{1 + \alpha^2} = \lambda - \frac{1}{2}$$

$$S_{AC}^\alpha = S_B^\alpha$$

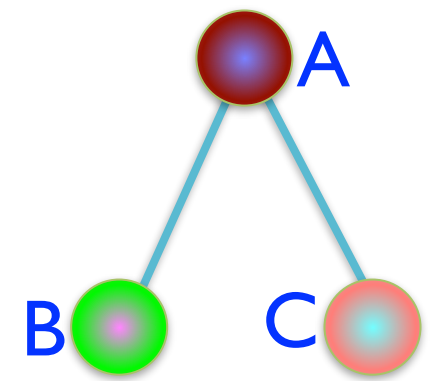
$$C_{adv}(|\psi_\alpha\rangle) = 1 + \lambda \log \lambda + (1 - \lambda) \log(1 - \lambda)$$

$$= 1 - H(\lambda)$$

Shannon entropy

3-qubit

MDCC states



Arbitrary state $|\psi\rangle$

$$\max\{\lambda_A, \lambda_B, \lambda_C\} = \lambda_B$$

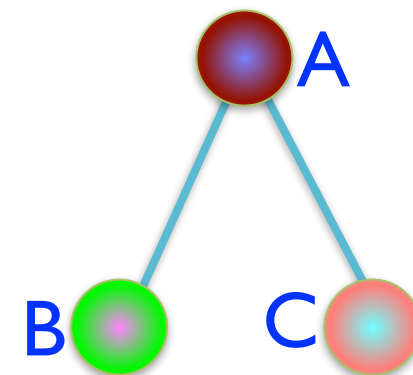
Since $\lambda_B, \lambda_C \geq \frac{1}{2}$ \longrightarrow $S_C - S_B \geq 0$
 ≤ 1

$$\begin{aligned} \mathcal{C}_{adv}(|\psi\rangle) &= \max\{S_B - S_C, S_C - S_B, 0\} \\ &= S_C - S_B \end{aligned}$$

$$\mathcal{C}_{adv}(|\psi_\alpha\rangle) = 1 - H(\lambda_B)$$

$$\mathcal{C}_{adv}(|\psi_\alpha\rangle) \geq \mathcal{C}_{adv}(|\psi\rangle) = \mathcal{C}_{adv}(|\psi\rangle)$$

MDCC states



Arbitrary state $|\psi\rangle$

$$\max\{\lambda_A, \lambda_B, \lambda_C\} = \lambda_A$$

$$\lambda_A \geq \lambda_B \geq \lambda_C \quad \lambda_A, \lambda_B, \text{ and } \lambda_C \geq \frac{1}{2}$$

$$S_C - S_B \geq 0$$

$$S_A \leq S_B$$

$$\mathcal{C}_{adv}(|\psi_\alpha\rangle) = 1 - H(\lambda_A) = 1 - S_A \geq S_C - S_B = \mathcal{C}_{adv}(|\psi\rangle)$$

MDCC states

3-qubit

Maximally-dense-coding-capable states (MDCC)

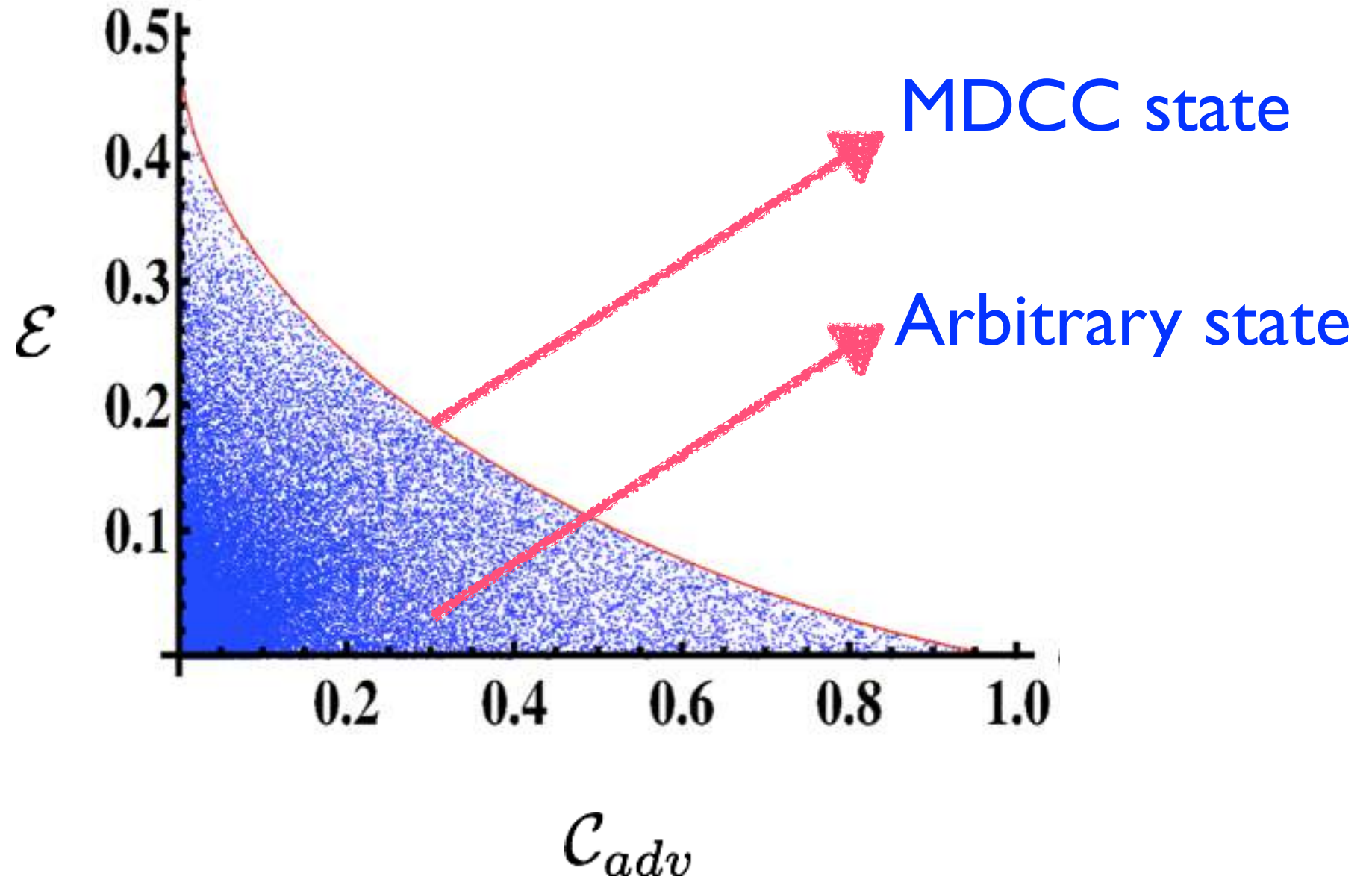
$$|\psi_\alpha\rangle = |111\rangle + |000\rangle + \alpha(|101\rangle + |010\rangle)$$

Arbitrary state $|\psi\rangle$ and set $\mathcal{E}(|\psi_\alpha\rangle) = \mathcal{E}(|\psi\rangle)$

$$\mathcal{C}_{\text{adv}}(|\psi_\alpha\rangle) \geq \mathcal{C}_{\text{adv}}(|\psi\rangle)$$

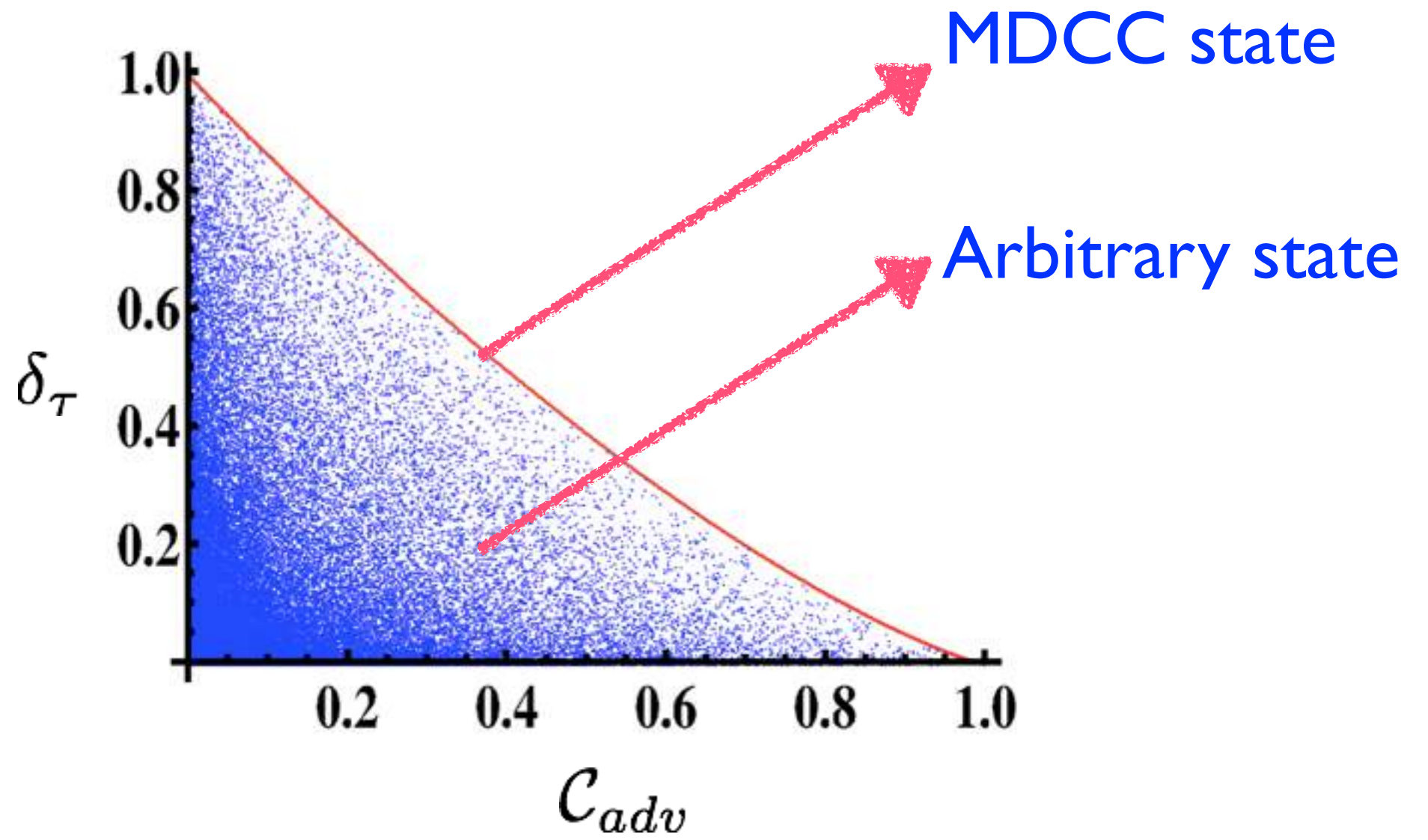
DC advantage & GGM

3-qubit



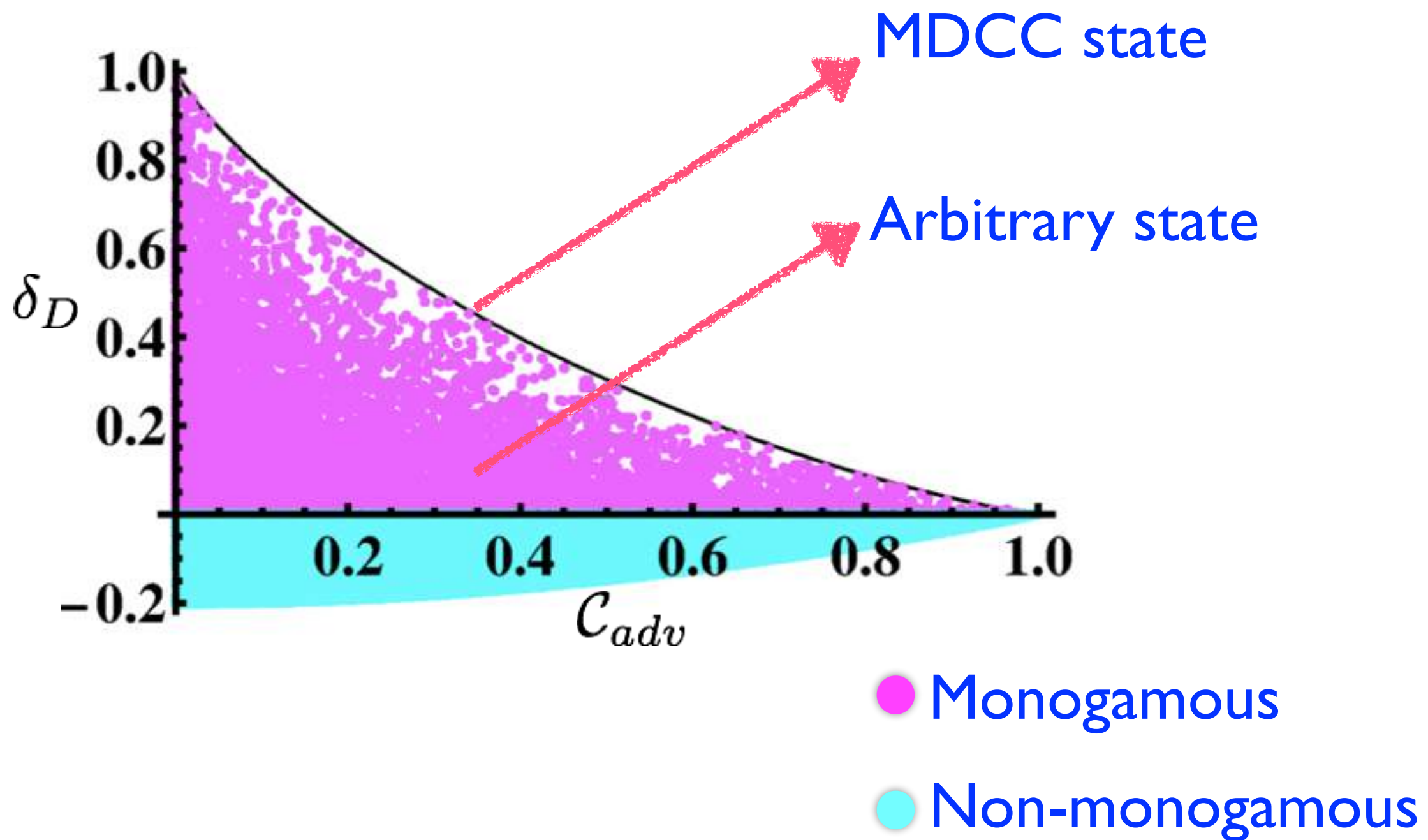
DC advantage & QC Score

3-qubit



DC advantage & QC Score

3-qubit



DC advantage & QC Score

Arbitrary state $|\psi\rangle$ and set $\delta_{\mathcal{Q}}(|\psi_{\alpha}\rangle) = \delta_{\mathcal{Q}}(|\psi\rangle)$

$$\mathcal{C}_{\text{adv}}(|\psi_{\alpha}\rangle) \geq \mathcal{C}_{\text{adv}}(|\psi\rangle)$$

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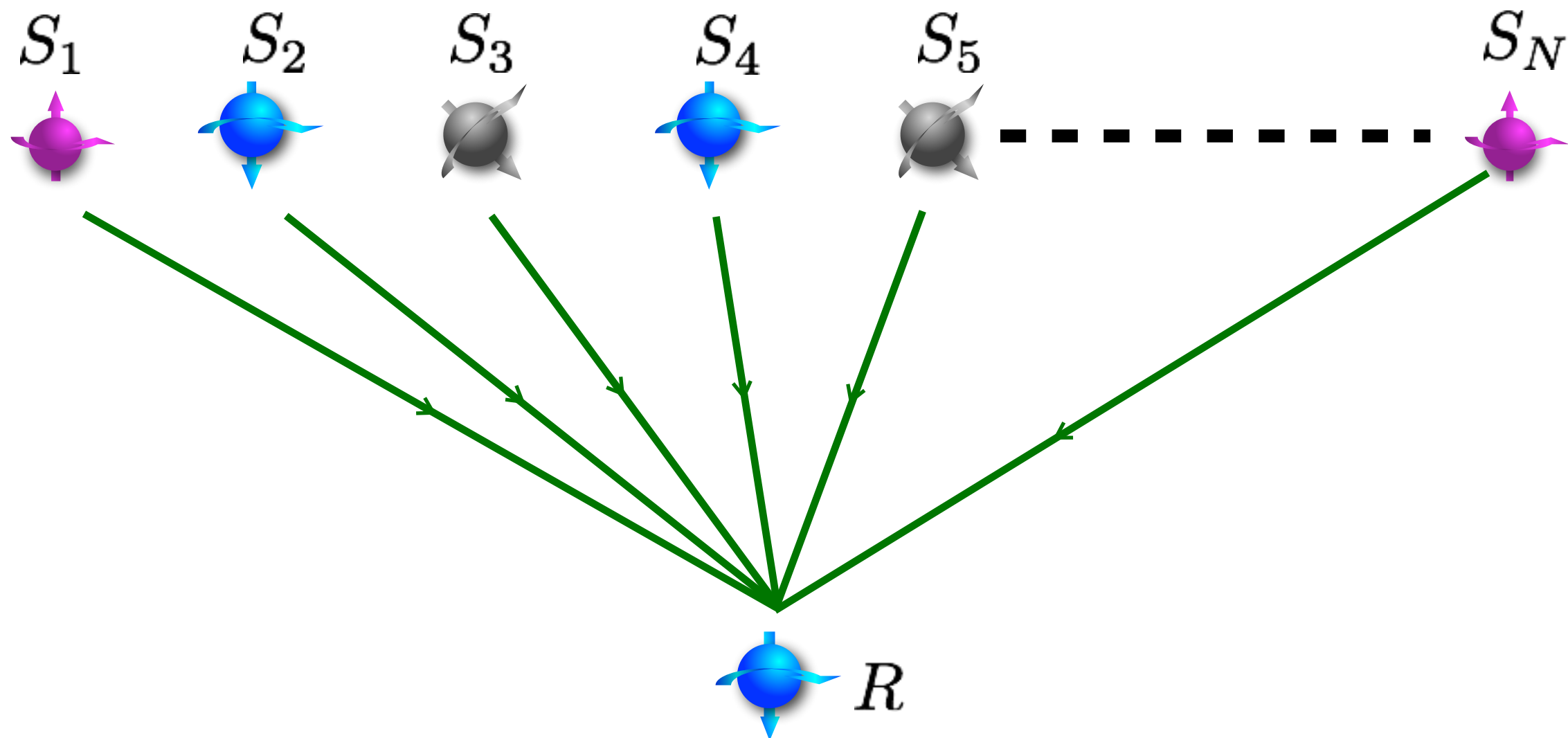
Dense coding - many to single

Noiseless

Noisy

Tamoghna's
poster

Dense coding - many to single



$$C(\rho_{S_1 \dots S_N R}) = \frac{\log_2 d_{S_1 \dots S_N} + S(\rho_R) - S(\rho_{S_1 \dots S_N R})}{\log_2 d_{S_1 \dots S_N R}}$$

$$d_{S_1 \dots S_N} = d_{S_1} \dots d_{S_N}$$

Dense coding - many to single

Generalized GHZ state

$$\begin{aligned} |GGHZ\rangle_{S_1 S_2 \dots S_N R} = & \sqrt{\alpha} |0_{S_1} \dots 0_{S_N}\rangle |0_R\rangle \\ & + \sqrt{1 - \alpha} e^{i\phi} |1_{S_1} \dots 1_{S_N}\rangle |1_R\rangle \end{aligned}$$

Arbitrary state $|\psi\rangle$ and set $C(|GGHZ\rangle) = C(|\psi\rangle)$

$$\mathcal{E}(|\psi\rangle) \leq \mathcal{E}(|GGHZ\rangle)$$

Dense coding - many to single

3-qubit

$$\begin{aligned} C(|GGHZ\rangle) &= \frac{2}{3} + \frac{S(\rho_R)}{3} \\ &= \frac{2}{3} - \frac{\alpha \log_2 \alpha + (1 - \alpha) \log_2 (1 - \alpha)}{3} \end{aligned}$$

$$C(|\psi\rangle) = \frac{2}{3} - \frac{\lambda_R \log_2 \lambda_R + (1 - \lambda_R) \log_2 (1 - \lambda_R)}{3}$$

$$C(|GGHZ\rangle) = C(|\psi\rangle)$$

$$\alpha = \lambda_R$$

Dense coding - many to single

3-qubit

$$\mathcal{E}(|GGHZ\rangle) = 1 - \alpha$$

$$\alpha \geq 1/2$$

$$\mathcal{E}(|\psi\rangle) = 1 - \max[\{l_A\}]$$

Set of all max
eigenvalues
all bipartitions

If max eigenvalue
is from receiver
state (rank 2)

$$\mathcal{E}(|\psi\rangle) = 1 - \lambda_R = 1 - \alpha = \mathcal{E}(|GGHZ\rangle)$$

Dense coding - many to single

3-qubit

If $\lambda_R \neq \max\{\lambda_A\}$

$$\lambda_R \leq \lambda_0 = \max\{\lambda_A\}$$

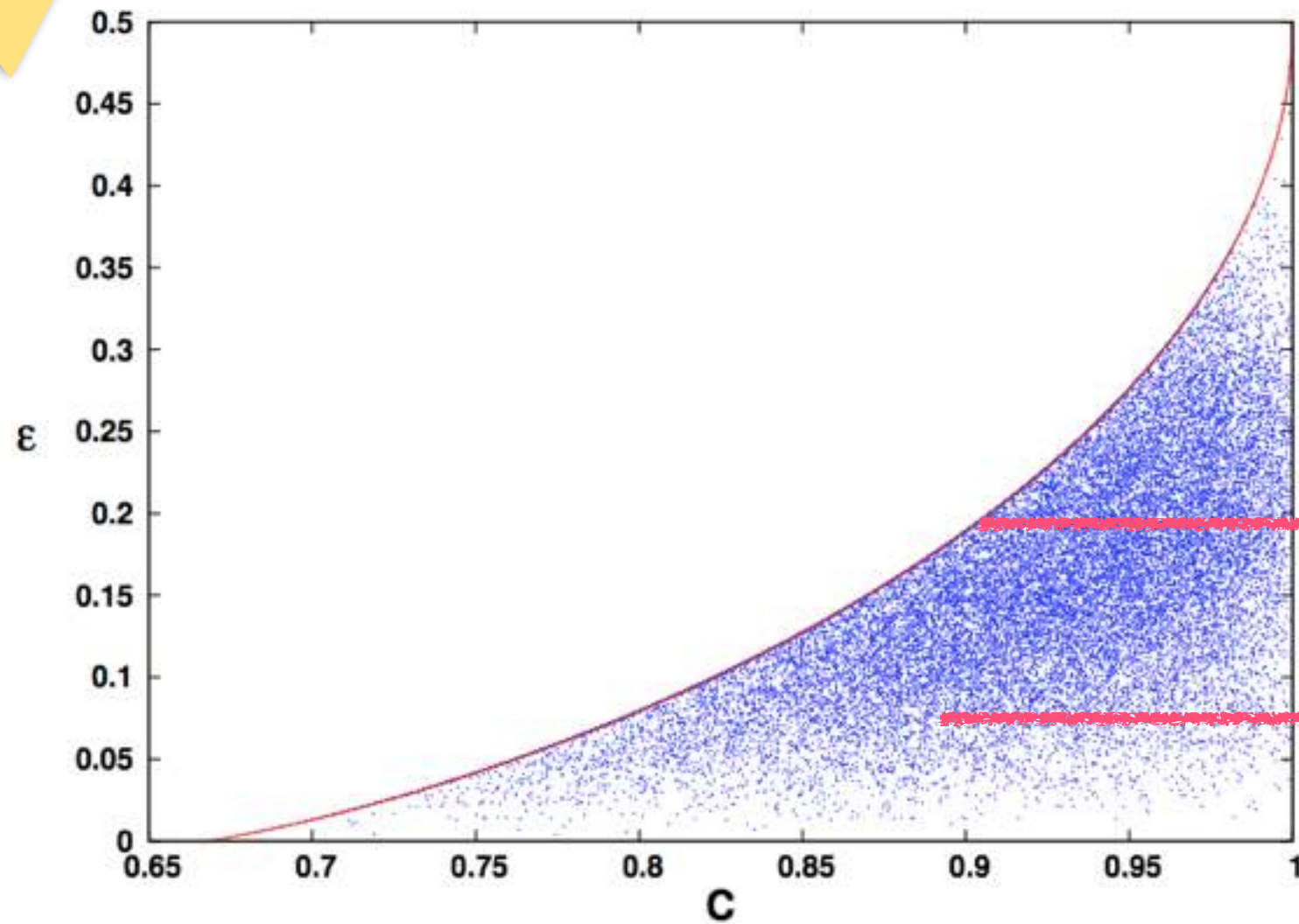
$$\mathcal{E}(|\psi\rangle) = 1 - \lambda_0 \leq 1 - \lambda_R = 1 - \alpha = \mathcal{E}(|GGHZ\rangle)$$

Arbitrary state $|\psi\rangle$ and set $C(|GGHZ\rangle) = C(|\psi\rangle)$

$$\mathcal{E}(|\psi\rangle) \leq \mathcal{E}(|GGHZ\rangle)$$

Dense coding - many to single

3-qubit



GGHZ

Arbitrary state

Dense coding - many to single

Generic ?

If receiver is set as
nodal observer

$$\delta_D(|\psi\rangle) \leq \delta_D(|GGHZ\rangle)$$

$$\delta_\tau(|\psi\rangle) \leq \delta_\tau(|GGHZ\rangle)$$

Conclusions

DC advantage in network

Complimentarity b/w QC and DC adv

$$C_{\text{adv}}^{\text{max}} + E_R \leq \log_2 d$$

$$\frac{1}{\log_2 d} C_{\text{adv}}^{\text{max}} + \frac{d}{d-1} \varepsilon \leq 1$$

Found MDCC states

Related multi DC and multi QC

$$\mathcal{E}(|\psi\rangle) \leq \mathcal{E}(|GGHZ\rangle)$$



<http://www.hri.res.in/~qic/>

Thank you