Controlling quantum state and its coherence

S N Bose National Centre for Basic Sciences



STATE \rightarrow SUPERPOSITION \rightarrow COHERENCE

CONTROLLING STATE



CONTROLLING COHERENCE

Outline of talk

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Quantum correlations

Ducertainty relations

Devantum coherence

Steering of quantum state

Steering of quantum coherence



Example: Pure state and mixed state





It is the weakest form of non-local correlation



It is the strongest form of non-local correlation



Steerable, if Alice can control the state of Bob's system

S. J. Jones, H. M. Wiseman, and A. C. Doherty, Phys. Rev. A 76, 052116 (2007).



Uncertainty relation

Uncertainty relation

Two non-commuting observables can not be measured simultaneously with arbitrary precision.

Coarse grained form : Heisenberg uncertainty relation (HUR) and entropic form of uncertainty relation (EUR).

Fine-grained form : Fine-grained uncertainty relation (FUR).

HUR

$$\Delta \mathcal{A} \ \Delta \mathcal{B} \geq \frac{|\langle [\mathcal{A}, \mathcal{B}] \rangle|}{2}$$

Here, uncertainty is measured by standard deviation which is coarse grained measure of uncertainty.

$$\Delta \mathcal{A} = \sqrt{\sum_{a} a^2 p_a - \sum_{a} (a p_a)^2}$$

W. Heisenberg, Z. Phys. 43, 172 (1927); E. H. Kennard, Z. Phys. 44, 326 (1927).



Here, uncertainty is measured by Shannon entropy.

$$\mathcal{H}(\mathcal{A}) = -\sum_{a} p_{a} \log p_{a}$$

and the complementarity, 1/c, is defined as

$$c = \max_{a,b} |\langle a|b\rangle|^2$$

H. Maassen, and J. B. M. Uffink, Phys. Rev. Lett. 60, 1103 (1988).

Fine-grained Uncertainty relation

Here, uncertainty is measured by probability of a particular measurement outcome or a combination of measurement outcomes, i.e., fine-graining of all possible outcomes.



$$P_{\text{Success}} = \sum_{i=1}^{N} P(\sigma_i) P(b_{\sigma_i} = 0) \leq P_{\text{Success}} = \max_{\rho_B} P_{\text{Success}}$$

J. Oppenheim, and S. Wehner, Science 330, 1072 (2010).



 $P(0_{\sigma_x}) + P(0_{\sigma_y}) + P(0_{\sigma_z}) \le \frac{3}{2} + \frac{3}{2\sqrt{3}}$

Quantum coherence





l₁-norm: $C^{l_1}(\rho_B) = \sum |\rho_B(i,j)|$ $i, j, i \neq j$ Relative entropy of coherence: $C^E(\rho_B) = \mathcal{S}(\rho_B^D) - \mathcal{S}(\rho_B)$ ρ_B^D : diagonal matrix formed with diagonal element of ρ_B .

Measures of Quantum coherence

T. Baumgratz, M. Cramer, and M. B. Plenio, Phys. Rev. Lett. 113, 140401 (2014).



Skew information: the coherence of the state in the basis of eigenvectors of the observable

$$C^{S}_{\mathcal{B}}(\rho_{B}) = -\frac{1}{2} Tr \left[\sqrt{\rho_{B}}, \mathcal{B}\right]^{2}$$

S. Luo, Phys. Rev. Lett. 91, 180403 (2003); Theor. Math. Phys. 143, 681 (2005).

Quantum coherence

Is it possible to measure quantum coherence with arbitrary precision in all possible mutually noncommuting basis, simultaneously?

Coherence complementarity relations

Coherence complementarity relations

$$C_x^{l_1}(\rho) + C_y^{l_1}(\rho) + C_z^{l_1}(\rho) \le \sqrt{6}$$

$$C_x^E(\rho) + C_y^E(\rho) + C_z^E(\rho) \le 2.23$$

 $C_k^{E(l_1)}(\rho)$: is calculated by writing the state in basis σ_k .

$$\rho_{\max}^{C} = \frac{1}{2} \left(I + \frac{1}{\sqrt{3}} \left(\sigma_x + \sigma_y + \sigma_z \right) \right)$$

Coherence complementarity relations
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 $C_k^S(\rho)$: is measured in basis σ_k .

$$\rho_{\max}^{C} = \frac{1}{2} \left(I + \frac{1}{\sqrt{3}} \left(\sigma_x + \sigma_y + \sigma_z \right) \right)$$



Quantum Steering

Quantum steering

EPR paradox

Entanglement is used to put question about incompleteness of quantum physics by Einstein, podolsky and Rosen.

Steering

Schrodinger re-expressed EPR paradox as the power to control of one system by distantly located system.

A. Einstein, D. Podolsky, and N. Rosen, Phys. Rev. 47, 777 (1935).

E. Schrodinger, Proc. Cambridge Philos. Soc. 31, 553 (1935); 32, 446 (1936).



Steerability : Alice's control on the state of Bob's system, i.e., Bob can know his system with higher precision than allowed by uncertainty principle.



Bob's Uncertainty of system B is zero when Alice communicates her results



Steerability : Absence of local hidden state (LHS) model for Bob's system.

$$\rho_{AB} \neq \sum_{\lambda} P(\lambda) \, \rho_{\lambda}^{A} \otimes \rho_{\lambda,Q}^{B}$$

 $P(a_{\mathcal{A}}, b_{\mathcal{B}}) \neq \sum_{\lambda} P(\lambda) P(a_{\mathcal{A}}|\lambda) P_{Q}(b_{\mathcal{B}}|\lambda)$

S. J. Jones, H. M. Wiseman, and A. C. Doherty, Phys. Rev. A 76, 052116 (2007).

Local Hidden State model

LHS model

Local : Alice prepares system B in a state quantumly uncorrelated with other systems possessed by Alice.Hidden : Bob has no information about the state of B.

$$\rho_{AB} = \sum_{\lambda} P(\lambda) \rho_{\lambda}^{A} \otimes \rho_{\lambda,Q}^{B}$$

Result : Once Alice sends the system B to Bob, Alice does not have any control on the state of system B. When Alice and Bob share steerable state, Alice can reduce Bob's uncertainty about his system by controlling its state.

Steering criterion : Intuition

$$P(b_{\mathcal{B}}|a_{\mathcal{A}})$$
It should violate some local
uncertainty relation satisfied
by

$$P(b_{\mathcal{B}})$$



Fine-grained steering criteria

$$P(a_{\mathcal{A}}, b_{\mathcal{B}}) = \sum_{\lambda} P(\lambda)P(a_{\mathcal{A}}|\lambda) P_{Q}(b_{\mathcal{B}}|\lambda)$$

$$+$$

$$q_{min} \sum_{i} p_{i} \leq \sum_{i} p_{i} q_{i} \leq q_{max} \sum_{i} p_{i}$$

$$\alpha$$

$$P(\mathbf{b}_{\sigma_{x}}) + P(\mathbf{b}_{\sigma_{y}}) + P(\mathbf{b}_{\sigma_{y}}) \leq \frac{3}{2} + \frac{3}{2\sqrt{3}}$$

$$P(b_{\sigma_{x}}|a_{\mathcal{A}_{1}}) + P(b_{\sigma_{y}}|a_{\mathcal{A}_{2}}) + P(b_{\sigma_{z}}|a_{\mathcal{A}_{3}}) \leq \frac{3}{2} + \frac{3}{2\sqrt{3}}$$

T. Pramanik, M. Kaplan, and A. S Majumdar, Phys. Rev. A 90, 050305(R) (2014).



All pure entangled states are maximally steerable

 $P(b_{\sigma_x}|a_{\mathcal{A}_1}) + P(b_{\sigma_y}|a_{\mathcal{A}_2}) + P(b_{\sigma_z}|a_{\mathcal{A}_3}) = 3$

T. Pramanik, M. Kaplan, and A. S Majumdar, Phys. Rev. A 90, 050305(R) (2014).



 $\rho_S = \frac{|01\rangle - |10\rangle}{\sqrt{2}}$



Steerability of quantum coherence



Steerability of quantum coherence

$$C_{x}^{l_{1}}(\eta_{B|\Pi_{y(z)}^{a}}) + C_{y}^{l_{1}}(\eta_{B|\Pi_{z(x)}^{a}}) + C_{z}^{l_{1}}(\eta_{B|\Pi_{x(y)}^{a}}) > \sqrt{6}$$

$$C_{x}^{E}(\eta_{B|\Pi_{y(z)}^{a}}) + C_{y}^{E}(\eta_{B|\Pi_{z(x)}^{a}}) + C_{z}^{E}(\eta_{B|\Pi_{x(y)}^{a}}) > 2.23$$

$$C_{x}^{S}(\eta_{B|\Pi_{y(z)}^{a}}) + C_{y}^{S}(\eta_{B|\Pi_{z(x)}^{a}}) + C_{z}^{S}(\eta_{B|\Pi_{x(y)}^{a}}) > 2$$

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All pure entangled states are maximally steerable

 $C_x(\eta_{B|\Pi_y^a}) + C_y(\eta_{B|\Pi_z^a}) + C_z(\eta_{B|\Pi_x^a}) = 3$



State Steerability : p > 0.58



Summary

Steering is a kind of non-local correlation where one of the systems is not trusted as quantum system.

Fine-grained steering criterion overcomes the limitations of coarse grained form of steering criteria.

Coherence complementary relation : No single quantum state is fully coherence under all non-commuting basis.

Vith three measurement settings, Werner state is steerable (state property) for p > 0.58.

Vith three measurement settings, Werner state is steerable (coherence property) for p > 0.82.

Thank You